The EBLM project – IX. Five fully convective M-dwarfs, precisely measured with CHEOPS and TESS light curves*

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Accepted. Received; in original form

ABSTRACT

Eclipsing binaries are important benchmark objects to test and calibrate stellar structure and evolution models. This is especially true for binaries with a fully convective M-dwarf component for which direct measurements of these stars' masses and radii are difficult using other techniques. Within the potential of M-dwarfs to be exoplanet host stars, the accuracy of theoretical predictions of their radius and effective temperature as a function of their mass is an active topic of discussion. Not only the parameters of transiting exoplanets but also the success of future atmospheric characterisation rely on accurate theoretical predictions. We present the analysis of five eclipsing binaries with low-mass stellar companions out of a sub-sample of 23, for which we obtained ultra high-precision light curves using the CHEOPS satellite. The observation of their primary and secondary eclipses are combined with spectroscopic measurements to precisely model the primary parameters and derive the M-dwarfs mass, radius, surface gravity, and effective temperature estimates using the PYCHEOPS data analysis software. Combining these results to the same set of parameters derived from TESS light curves, we find very good agreement (better than 1% for radius and better than 0.2% for surface gravity). We also analyse the importance of precise orbits from radial velocity measurements and find them to be crucial to derive M-dwarf radii in a regime below 5% accuracy. These results add five valuable data points to the mass-radius diagram of fully-convective M-dwarfs.

Key words: binaries: eclipsing – stars: fundamental parameters – stars: low-mass – techniques: photometric – techniques: spectroscopic

1 INTRODUCTION

Low-mass main-sequence stars of M-type (M-dwarfs) have been in the spotlight of recent exoplanet surveys (Nutzman & Charbonneau

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2008; Delrez et al. 2018; Barclay et al. 2018; Quirrenbach et al. 2019; Donati et al. 2020). This development has two main reasons. First their low masses, and radii, compared to F, G, and K stars make it easier to detect small planets and planetary systems composed of mini Neptunes down to Earth sized planets by means of radial velocity and transit methods (e.g. Gillon et al. 2016; Zechmeister et al. 2019; Günther et al. 2019). Thus, more Earth sized planets have been found in the habitable zone of M-dwarfs than for solar-type stars (e.g. Dressing & Charbonneau 2013). Second, M-dwarfs have low luminosities and, thus offer the first possible window to study transiting rocky planets in their habitable zone and directly analyse their atmospheres with high-precision instruments like the James Webb Space telescope (Kaltenegger & Traub 2009; Morley et al. 2017).

Such studies depend crucially on the knowledge of the parameters of M-dwarf planets which in turn are derived from the mass and radius of the host M-dwarf. Up to now our understanding on the mass and radius distribution of low-mass stars which are fully convective ($M_{\star} < 0.35\,M_{\odot}$, Chabrier & Baraffe 1997) is rather poorly explored compared to more massive stars. This is mainly due to the relative faintness of these stars ¹. Especially the lack of a large sample of M-dwarfs with directly measured mass and radius make it difficult to calibrate stellar evolution models which are typically used to estimate the properties of planet host stars like for example the Exeter/Lyon models (Baraffe et al. 2015) or the Dartmouth models (Dotter et al. 2008).

Studies of M-stars with available radii and masses have revealed that their stellar radii for a given mass are apparently inflated by a few percent, compared to estimates from models (e.g. Casagrande et al. 2008; Torres et al. 2010; Spada et al. 2013; Kesseli et al. 2018).

Several possible explanations have been discussed, like stellar magnetic activity (Mullan & MacDonald 2001; Chabrier et al. 2007), or a bias due to binarity (Ribas 2006; Morales et al. 2009). Also metallicity effects seem to play a role (Berger et al. 2006; von Boetticher et al. 2019). Thus, a representative sample of low-mass M-dwarfs with accurately measured mass, radius, but also metallicity is crucial to understand how the different effects enter into this radius inflation problem.

The eclipsing binaries with low mass (EBLM) project (Triaud et al. 2013) is focusing on a large sample with hundreds of eclipsing binaries of F,G, & K-type stars, orbited by late type M-dwarf companions. These binaries have been detected from the WASP survey (Pollacco et al. 2006). Using a large radial velocity follow-up campaign of these stars, Triaud et al. (2017) derived accurate orbits of many of these systems thus being able to measure fundamental parameters like precise mass and radius of the low-mass M-dwarfs. The binary configuration with a solar-type star allows us to measure accurately the metallicity of the solar-type star. Assuming an equal metallicity of both components, we can constrain the metallicity of the M-dwarf. Thus, EBLM targets are ideal candidates to populate the mass regime of fully convective M-dwarfs with masses below $0.35\,M_{\odot}$ and to establish an empirical mass-radius-metallicity relationship for these stars. Early results from sub samples indicate that models can be matched quite well, when taking accurate measurements of the metallicty of the M-dwarf into account (von Boetticher et al. 2019; Gill et al. 2019). Every low-mass M-dwarf with accurately measured mass, radius and metallicity will help to tighten the constraints on the source of the radius inflation problem and in return

will allow us in future to constrain precise parameters of planet host

CHEOPS (Benz et al. 2021) is a S-class mission of the European Space Agency, which has been launched on the 18th of December 2019. Its primary mission is to perform ultra high-precision photometry of bright exoplanet host-stars. We have started an 'Ancillary Science' programme on a selection of 23 EBLM targets, to obtain precise measurements of primary and secondary eclipses, which allow us to (i) derive the size of both components and (ii) to measure the M-dwarf effective temperature from the surface brightness ratios. Additionally, we use light curves, obtained by the TESS survey (Ricker et al. 2015), which covers the northern and southern hemispheres with observing periods of about one month per pointing (sector). TESS cameras have a three times smaller aperture compared to CHEOPS, leading to a lower accuracy for eclipse eyents in TESS data. Nevertheless, the long coverage of photometric data allows us to gather multiple eclipses of our targets and thus improve and compare orbital parameters, as well as to optimise our analysis of CHEOPS observations.

The three EBLM binaries, analysed in our *CHEOPS* programme EBLM J1741+31, EBLM J1934-42 and EBLM J2046+06 have shown that M-dwarfs with precisely measured radii and metallicities open up the possibility to disentangle the effect of metallicity from different effects on the radius inflation problem for low-mass M-dwarfs (Swayne et al. 2021).

In this paper we present the analysis of five EBLM binaries with fully convective M-dwarfs companions, observed in our *CHEOPS* programme and compare them to the analysis of *TESS* observations.

2 OBSERVATIONS AND METHODS

Primary and secondary eclipses for all our five eclipsing binaries were observed with CHEOPS between November 2020 and January 2021 as part of CHEOPS Guaranteed Time Observation programme ID-037. We obtained one primary eclipse and, depending on the depth of the secondary eclipse, one to three secondary eclipse observations in order to obtain sufficient signal to noise to measure both eclipses. Table 1 gives an overview of the CHEOPS observations and data extraction. All data were reduced by the CHEOPS data reduction pipeline v13.1 (Hoyer et al. 2020), which performs an aperture photometry of the target star, taking contamination in the field as well as instrumental effects like the rotation of the satellite into account. The pipeline offers light curves for different aperture sizes. For our analysis, we selected the aperture size with minimal median absolute deviation of the point-to-point difference in the light curve. The resulting aperture radii are listed as R_{ap} in Table 1. The observations were interrupted due to the low-Earth-orbit of CHEOPS by Earth occultations, as well as crossings of the South Atlantic Anomaly. We derive the time spent on target as the fraction of valid observations compared to the total observation interval.

The *TESS* survey covered all of our targets with 2-min cadence data made available by *TESS* Guest Investigator (GI) programmes. EBLM J0239-20 (TIC64108432) has been observed in sectors 4 and 31 under GI programmes G011278 and G03216. EBLM J0540-17 (TIC46627823) has been observed in sectors 6 and 32 under GI programmes G011278, G03216, & G03251. EBLM J0546-18 (TIC93334206) has been observed in sectors 32 and 33 under GI programme G03216. EBLM J0719+25 (TIC458377744) has recently been observed in sectors 44, 45, & 46 under GI programme G04157 and EBLM J2359+44 (TIC177644756) has been observed in sector 17 under GI programmes G022253 & G022156. Data reduction and

¹ E.g. the planet host star TRAPPIST-1, a M7.5 ultra-cool dwarf in 12 pc distance has a visual magnitude of only 18.8 mag.

Table 1. CHEOPS observations and data extraction for our targets. Effic. is the fraction of the observation that resulted in valid (usable) data and R_{ap} the aperture radius used to extract the light curves.

Eclipse Event	Target	Start date (UTC)	Duration (h)	T _{exp} (s)	Effic.	File key	R _{ap} (pixels)
Primary	EBLM J0239-20	2020-11-01T15:43	8.80	60	86.2	CH_PR100037_TG012001_V0200	25
Secondary		2020-11-05T20:30	7.99	60	93.2	CH_PR100037_TG011901_V0200	25
Secondary		2020-11-19T17:24	9.02	60	70.4	CH_PR100037_TG011902_V0200	25
Primary	EBLM J0540-17	2020-12-07T08:39	10.04	60	68.4	CH_PR100037_TG012601_V0200	17.5
Secondary		2021-01-21T09:39	10.75	60	54.1	CH_PR100037_TG012502_V0200	17.5
Secondary		2020-12-04T08:13	10.62	60	66.5	CH_PR100037_TG012501_V0200	17.5
Secondary		2021-01-27T09:20	10.49	60	50.0	CH_PR100037_TG012503_V0200	17.5
Primary	EBLM J0546-18	2020-11-30T22:27	8.67	60	67.5	CH_PR100037_TG012801_V0200	25
Secondary		2020-12-31T05:35	8.77	60	66.3	CH_PR100037_TG012701_V0200	25
Secondary		2021-01-09T19:50	8.05	60	64.0	CH_PR100037_TG012702_V0200	25
Primary	EBLM J0719+25	2020-12-10T07:03	8.80	60	52.8	CH_PR100037_TG013001_V0200	22.5
Secondary		2021-02-03T20:54	8.69	60	57.7	CH_PR100037_TG017301_V0200	22.5
Secondary ¹		2020-12-21T12:03	8.50	60	60.2	CH_PR100037_TG012901_V0200	22.5
Secondary	EBLM J2359+44	2020-11-11T08:59	8.89	60	58.3	CH_PR100037_TG016301_V0200	26.5
Primary		2020-11-28T13:07	15.67	60	51.4	CH_PR100037_TG016401_V0200	26.5

¹ For this observation the secondary eclipse of EBLM J0719+25 has been missed, thus we cannot use this data set for parameter determination of the binary.

light curve extraction were done by the *TESS* Science Processing Operations Center Pipeline (SPOC; Jenkins et al. 2016) and were downloaded via the Mikulski Archive for Space Telescopes² (MAST). For our analysis, we used Pre-search Data Conditioned Simple Aperture Photometry (PDCSAP) flux data and bitmask 175 to exclude data flagged with severe quality issues (Tenenbaum & Jenkins 2018).

For EBLM J2359+44 two radial velocity measurements have been published by Poleski et al. (2010) that confirmed it to be a binary star. Full time series radial velocity observations of EBLM J0719+25 and EBLM J2359+44 have been taken with the SOPHIE high-resolution echelle spectrograph (Perruchot et al. 2008), mounted on the 1.93 m telescope at the Observatoire de Haute-Provence in France as part of the Binaries Escorted By Orbiting Planets (BEBOP) survey to search for circumbinary planets (Martin et al. 2019). For EBLM J0719+25, 8 SOPHIE spectra have been obtained between November 2018 and October 2019 in High-Resolution mode (R = 75,000). For EBLM J2359+44, 15 SOPHIE spectra have been obtained between November 2018 and September 2020 in High-Resolution mode (R = $75\ 000$) as well as in High-Efficiency (HE) mode (R = $40\ 000$). The HE mode allows an about 2.5 times higher throughput compared to the High-Resolution mode. The spectra have an average signal to noise of about 30 with a typical exposure time of 1800 s. To allow the removal of the background contamination from the Moon, all observations were taken with one fibre on target and one on the sky. The spectra were reduced using the SOPHIE Data Reduction Software (Baranne et al. 1996) and radial velocities were measured by cross correlation with a G2 mask (Courcol et al. 2015) for which we achieved a typical precision of 10 m s⁻¹ for our spectra. All radial velocity measurements are listed in the Appendix Tables B1 & B2. We submitted a target list of 40 EBLM systems from Triaud et al. (2017) as a priority 4 proposal to be observed with high resolution spectrograph (Crause et al. 2014) of the Southern African Large Telescope (SALT) in medium resolution ($R \approx 37,000$). In total, 30 of them were observed between the 19th of May and 7th August 2017,

including EBLM J0239-20. These observations were made in long slit mode with an exposure time scaling as a function of magnitude to ensure a SNR ≥ 100 . Data was reduced and processed using standard pipelines (Crawford 2015; Craig et al. 2015) to produce two spectra for each observation (370–550 nm & 550–890 nm) as a result of the dual-beam nature of the spectrograph.

3 ANALYSIS

For data analysis, we followed the methods, described in Swayne et al. (2021), hereafter SW21. Both TESS and CHEOPS light curves were modelled using the qpower2 transit model, which applies a power-2 limb darkening law (Maxted & Gill 2019). We use it as binary star model including primary and secondary eclipses which is implemented in the python software PYCHEOPS³ (Maxted et al. 2021). The parameters of the binary star model are the orbital period P, the mid-time of the primary eclipse T_0 ; the primary and secondary eclipse depths D and L, the impact parameter b, the parameters $f_c =$ $\sqrt{e}\cos(\omega)$ and $f_s = \sqrt{e}\sin(\omega)$, which parameterise the eccentricity e and the longitude of periastron ω , the limb darkening parameters h₁ and h₂ (Maxted 2018), and W, which becomes the width of the eclipse for e = 0 and is defined by the stellar radii, impact parameter, and the semi mayor axis a (see Maxted et al. (2021) for details). We used gaussian priors for f_c, f_s. These priors were derived from radial velocity measurements of the systems. Orbital parameters from radial velocity measurements for EBLM J0239-20, EBLM J0540-17, and EBLM J0546-18 have been published in Triaud et al. (2017). Their eccentricities are reported to be consistent to zero, thus we set those priors to zero for all three systems. For EBLM J0719+25 and EBLM J2359+44, we used the binary star python code ellc (Maxted 2016), to model the radial velocity from SOPHIE measurements as well as the two measurements from Poleski et al. (2010) for EBLM J2359+44. We sampled the posterior probability distribution (PPD) of our model parameters fc, fs, and the semi amplitude K, using the

² https://archive.stsci.edu/

https://github.com/pmaxted/pycheops

4 D. Sebastian et al.

Table 2. Stellar and orbital parameters of the primary stars. Coordinates are in J2000.

	EBLM J0239-20	EBLM J0540-17	EBLM J0546-18	EBLM J0719+25	EBLM J2359+44
Name	TYC 5862-1683-1	TYC 5921-745-1	TIC 93334206	TYC1913-0843-1	TYC3245-0077-1
RA	02 39 29.29	05 40 43.58	05 46 04.81	07 19 14.26	23 59 29.74
Dec.	$-20\ 02\ 24.0$	$-17\ 32\ 44.8$	-18 17 54.6	+25 25 30.8	+44 40 31.2
G (mag)	10.57	11.42	12.01	11.15	10.46
Sp. Type	G0	G0	G0	G0	F8
$T_{eff,1}(K)^a$	5758 ± 100	6290 ± 77	6180 ± 80	6026 ± 67	6799 ± 83
$\log g_1(cgs)^c$	4.053 ± 0.016	4.058 ± 0.017	4.100 ± 0.034	4.239 ± 0.022	4.068 ± 0.010
[Fe/H] ^a	0.27 ± 0.12	-0.04 ± 0.05	-0.45 ± 0.08	0.04 ± 0.05	0.12 ± 0.05
$R_1(R_{\odot})^c$	1.587 ± 0.039	1.636 ± 0.040	1.509 ± 0.064	1.305 ± 0.038	1.711 ± 0.033
$M_1(M_{\odot})^c$	1.037 ± 0.060	1.120 ± 0.062	1.051 ± 0.059	1.078 ± 0.059	1.253 ± 0.070
Orbital parameters:					
$K(km s^{-1})$	21.316±0.036d	16.199 ± 0.010^{d}	26.15 ± 0.10^{d}	15.02 ± 0.04^{b}	23.62 ± 0.08^{b}
e	< 0.0032 ^d	0.00029 ± 0.00057^{d}	< 0.015 ^d	0.0730 ± 0.0045^{b}	0.4773 ± 0.0010^{b}
$\omega(\deg)$	_	-164 ± 10^{d}	_	-155.8 ± 5.4^{b}	-94.290±0.060 ^b
$f(m) (10^{-3} M_{\odot})$	2.788 ± 0.014^{d}	2.6444 ± 0.0096^{d}	2.1332 ± 0.0023^d	2.597±0.021 ^b	10.53±0.11 ^b

References: a From spectral analysis, from radial velocity analysis, from light curve modelling, from Triaud et al. (2017)

Table 3. Priors on $f_C = \sqrt{e} \cos \omega$ and $f_S = \sqrt{e} \sin \omega$ used in the analysis of the *CHEOPS* and *TESS* light curves based on the spectroscopic orbits for each binary system.

Target	f_c	f_s
EBLM J0239-20	0.0	0.0
EBLM J0540-17	0.0	0.0
EBLM J0546-18	0.0	0.0
EBLM J0719+25	-0.247 ± 0.013	-0.111 ± 0.023
EBLM J2359+44	-0.0517 ± 0.0007	-0.6889 ± 0.0007

Markov chain Monte Carlo (MCMC) code EMCEE (Foreman-Mackey et al. 2013) to take the RV-jitter of the data into account by weighting the fit by the log-likelihood function. For this we used the period from our TESS fit (see Sec. 3.1) as fixed prior and did not need to fit any additional trend to the data. The resulting orbital parameters, as well as the mass function f(m) (see equation 6 in Triaud et al. 2017) are listed in Table 2. The resulting priors for f_c , f_s , are listed in Table 3. The errors represent the one sigma error of the resulting PPD.

3.1 TESS light curve analysis

Only segments of the *TESS* light curve within one eclipse duration of the time of mid-eclipse were used in this analysis. To remove trends in the light curve, we divided these segments by a linear polynomial model fitted to the data either side of the eclipse. Unlike SW21, we preferred this method over the use of a Gaussian process in order to securely preserve the transit shape of the faint secondary eclipses.

To model the light curve, we first determined the initial orbital parameters using a least-squares fit and then sampled the PPD of our transit model using EMCEE. We placed normal priors on the orbital parameters f_c , f_s , as listed in Table 3 as well as on the white noise, using the residual rms of the least-squares fit. The resulting parameters from the *TESS* light curves are detailed in Tables 4, 5, & 6. These represent the median of the PPD as well as the standard errors from the 15.9% and 84.1% percentile-points of the PPD. We show the resulting fits of all targets in the Appendix, Fig. D3 and Fig. D2.

3.2 CHEOPS light curve analysis

CHEOPS light curves were analysed in two steps. First we analysed every visit separately to derive initial model parameters (see Table 1 for an overview of all visits). As described in detail in SW21, instrumental effects like roll angle, contamination, and background level can be represented using linear correlation parameters or for roll angle ϕ , $\sin(\phi)$, $\cos(\phi)$, $\sin(2\phi)$, etc., which were iteratively selected⁴. The PPD of all model and decorellation parameters were sampled simultaneously using EMCEE. We used the same Gaussian priors for f_c, and f_s as for the TESS data and since we obtained single eclipse events, we fixed our transit model to accurately measured orbital period P, from the TESS light curve fit. For secondary eclipses, we used priors on the parameters D, W and b, as derived from the primary eclipse of each target.

In a second step, we were using a single MCMC to perform a 'multivisit' analysis including all visits for a specific target. We used the same priors as for the individual analysis as well as the results as input parameters and used the function multivisit of PYCHEOPS to sample the joint PPD with EMCEE. Hereby we used the implicit decorrelation method for instrumental trends as described in Maxted et al. (2021), keeping the number of harmonic terms to its default ($N_{\rm roll}=3$). The resulting parameters from the *CHEOPS* light curves are detailed in Tables 4, 5, & 6. These represent the median of the PPD as well as the standard errors from the 15.9% and 84.1% percentile-points of the PPD. We show the resulting fits of all targets in the Appendix, Fig. E1, Fig. E2, and Fig. E3 and in Table A1 the resulting decorrelation parameters from the multivisit analysis.

3.3 Stellar parameters

We used co-added high-resolution spectra to derive the stellar parameters of the primary components (T_{eff} and [Fe/H]). For EBLM J0540-17, we used co-added CORALIE spectra, obtained by Triaud et al. (2017) and available from the ESO science archive facility and co-added SOPHIE spectra for EBLM J0719+25 and EBLM J2359+44. The stellar parameters for these three targets were derived

⁴ See Table A1 for the decorrelation parameters selected for each visit

⁵ http://archive.eso.org/

using the equivalent width method following the same methodology, model atmospheres, and line list as described in Sousa (2014) and Santos et al. (2013). In here we applied the ARES code (Sousa et al. 2015), as well as the MOOG radiative transfer code (Sneden et al. 2012), assuming ionisation and excitation equilibrium of iron lines. For EBLM J0546-18 we used co-added CORALIE spectra and applied a wavelet decomposition method where we compare the coefficients from a wavelet decomposition to those from a grid of model spectra. Those model spectra were synthesised using the code SPECTRUM (Gray & Corbally 1994), MARCS model atmospheres (Gustafsson et al. 2008) as well as the atomic line list version 5 of the Gaia ESO survey (Heiter et al. 2015). The method is detailed in Gill et al. (2018) and has been found to deliver robust measurements for effective temperature and metallicity for spectra with relatively low SNR (SNR ≥ 40). For EBLM J0239-20 we used the SALT spectra and modeled the stellar fundamental parameters using the software SME⁶ (Spectroscopy Made Easy; Valenti & Piskunov 1996; Piskunov & Valenti 2017) that computes synthetic spectra with atomic and molecular line data from VALD⁷ (Ryabchikova et al. 2015) which is compared to the observations. We chose the stellar atmosphere grid Atlas12 (Kurucz 2013) and modelled Teff, log g1, abundances and v sin i one parameter at a time. Due to the high rotational velocity (v sin i = 31 ± 4 km s⁻¹), the uncertainties in log g_1 derived from the line wings of the Ca₁ triplet around 6200 is with 0.2 dex relatively high. We thus rely on the lightcurve modelling to derive the surface gravity of our targets.

Similarly to SW21, we derived the system parameters using the function massradius in PYCHEOPS. As explained in Maxted et al. (2021), this function applies a Monte Carlo approach to derive basic system parameters like the primaries mean stellar density, the mass and radius of the M-dwarf, using the PPD of our *CHEOPS* light curve fit. It additionally uses the primaries mass and radius, as well as the orbital parameters which were not sampled in the PPD like period, and eccentricity as input and derives the surface gravity $\log g_2$ of the M-dwarf using the radial-velocity semi-amplitudes. We used this function to optimise the global system parameters in a two stage iterative process.

In the first step, we used the primaries mass and radius estimates available from the *TESS* input catalogue v8 (Stassun et al. 2019) as initial parameters. The derivation of these estimates is based on an empirical relation including photometric effective temperature estimates for stars with well measured Gaia distances. We used the same priors for period and eccentricity that we used for our *CHEOPS* fit, as well as the semi-amplitudes from radial velocity measurements. For EBLM J0239-20, EBLM J0540-17, and EBLM J0546-18 we have used the published semi-amplitudes (Triaud et al. 2017), For For EBLM J0719+25 and EBLM J2359+44, we use the results from our orbital fit (see Table 2).

In a second iteration, we made use of the massradius function again in order to find the best fitting parameters of the primary mass and radius from our light curve fit. We used the relation of Enoch et al. (2010) (equation 4), to derive a mass sample for the primary star. This sample is based on the stellar density samples obtained from the first iteration and created similar sized samples for $T_{\rm eff}$ and [Fe/H] based on our spectroscopic stellar parameters. We then added a normal distributed scatter of 0.023 to account for the resulting scatter for this relation found by Enoch et al. (2010). We derived a radius sample using this mass sample as well as the density sample.

We used the mass and radius samples to re-run the massradius function to derive the final stellar parameters of the primary and M-dwarf components. We finally derived the surface gravity $\log g_1$ from the stellar density, directly measured from the light curve fit of our *CHEOPS* data, as well as the primaries mass derived from the previous step.

We derived the effective temperature $T_{eff,2}$ of the M-dwarf companion using the surface brightness ratio L/D, derived from the light curve fit of primary and secondary eclipses. Similar to SW21, we derived the integrated surface brightness in the *CHEOPS* and *TESS* bandbands of the primary star, using the spectral parameters $T_{eff,2}$, $\log g_1$, and [Fe/H] using PHOENIX model atmospheres with no alpha-element enhancement (Husser et al. 2013) and sampled a large set of surface temperatures over the known parameters, L/D, $\log g$, and [Fe/H] (assuming similar metallicity for both companions) to derive the effective temperature.

The light contribution from the primary star reflected to the M-dwarf can be expressed by $A_g(R_2/a)^2,$ where A_g is the geometric albedo and R_2/a is the radius of the M-dwarfs in units of the semi mayor axis, which we directly measure from our model. With a typical albedo of $A_g\sim 0.1$ (Marley et al. 1999), the light contribution for our targets is very small and thus negligible. Nevertheless, for the two shortest period binaries in our sample, EBLM J0239-20 and EBLM J0546-18 the light contribution might cause an underestimation of the secondary eclipse depth on the one sigma level and thus an underestimation of $T_{\rm eff,2}$ in the order of $1\,\%$ for both CHEOPS and TESS passbands. Thus, we increased the relative uncertainties for $T_{\rm eff,2}$ for EBLM J0239-20 and EBLM J0546-18 by 1% in order to account for the unknown uncertainty of A_g .

All parameters of the primary stars are listed in Table 2, all parameters for the M-dwarf companions are listed in Tables 4, 5, & 6.

4 DISCUSSION

We have derived the stellar parameters for both companions for all of our targets thanks to high precision *CHEOPS* light curves. For the M-dwarfs we derive accurate radii with an average uncertainty of $3.2 \pm 1.3\%$ and the surface gravity with an average uncertainty of $0.4 \pm 0.3\%$. This precision for the surface gravity of M-dwarfs is better then, or hardly reached with state of the art high-resolution spectroscopic measurements of field M-dwarfs (e.g. Olander et al. 2021; Marfil et al. 2021).

4.1 Radial velocity priors

We used priors obtained from the radial velocity (RV) orbital parameters eccentricity (e) and longitude of periastron (ω) to fit our CHEOPS and TESS light curves. Only EBLM J0719+25 and EBLM J2359+44 have eccentricities significantly larger than zero, the others we have fixed to zero eccentricity. We analysed the effect of imposing RV priors on the CHEOPS parameter fit by repeating it with f_c and f_s kept as free parameters. Two of our binaries with previously fixed eccentricities, resulted in eccentricties consistent to zero with EBLM J0239-20 (e = 0.028 ± 0.058) and EBLM J0546-18 (e = 0.0005 ± 0.008 0.0007). For EBLM J0540-17 and EBLM J0719+25 this fit resulted in a longer MCMC chain, which finally ended with a less uniformly defined PPD for W, which was strongly correlated to fc and fs. This led to up to 5% overestimated radii for the M-dwarfs. Except for these two stars, the derived model parameters did not deviate more than 1σ from the parameters listed in Tables 4, 5, and 6. Nevertheless, we found that for the orbital parameters all resulting uncertainties

⁶ http://www.stsci.edu/~valenti/sme.html

⁷ http://vald.astro.uu.se

Table 4. The derived parameters for EBLM J0239-20 and EBLM J0540-17 using *CHEOPS* and *TESS* light curve fits with eclipse depths being in the relevant instrumental bandpass.

	EBLM J	10239-20	EBLM J0540-17		
	CHEOPS	TESS	CHEOPS	TESS	
Model parameters					
$T_0(BJD)$	2163.70805 ± 0.00015	1413.46145 ± 0.00012	2209.12086 ± 0.00021	1470.51285 ± 0.00030	
P (days)	2.778691 (fixed)	2.778691 ± 0.000001	6.004940 (fixed)	6.004940 ± 0.000003	
D	0.01679 ± 0.00019	0.016716 ± 0.000092	0.01404 ± 0.00021	0.01381 ± 0.00018	
W	0.05268 ± 0.00037	0.05286 ± 0.00015	0.03818 ± 0.00019	0.03827 ± 0.00018	
b	0.654 ± 0.014	0.6428 ± 0.0092	0.167 ± 0.105	0.253 ± 0.089	
f_c	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	
f_s	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	
L	$(3.68 \pm 0.45) \times 10^{-4}$	$(7.30 \pm 0.42) \times 10^{-4}$	$(3.66 \pm 0.53) \times 10^{-4}$	$(6.61 \pm 0.78) \times 10^{-4}$	
h_1	0.766 ± 0.020	0.836 ± 0.011	0.767 ± 0.015	0.811 ± 0.013	
h_2	0.47 ± 0.22	0.59 ± 0.20	0.54 ± 0.18	0.47 ± 0.21	
Derived parameters					
R_2/R_1	0.12957 ± 0.00073	0.12929 ± 0.00035	0.11850 ± 0.00087	0.11752 ± 0.00075	
R_1/a	0.1797 ± 0.0027	0.1788 ± 0.0015	0.1084 ± 0.0018	0.1105 ± 0.0023	
R_2/a	0.02288 ± 0.00042	0.02289 ± 0.00024	0.01265 ± 0.00028	0.01264 ± 0.00034	
i(°)	83.25 ± 0.24	83.40 ± 0.15	88.96 ± 0.67	88.40 ± 0.59	
e	0.0	0.0	0.0	0.0	
$\omega(^{\circ})$	_	_	_	\mathcal{L}	
Absolute parameters					
a(AU)	0.04106 ± 0.00076	0.04107 ± 0.00076	0.0703 ± 0.0012	0.0703 ± 0.0012	
$R_2(R_{\odot})$	0.2056 ± 0.0052	0.2041 ± 0.0044	0.1939 ± 0.0050	0.1959 ± 0.0056	
$M_2(M_{\odot})$	0.1597 ± 0.0059	0.1597 ± 0.0059	0.1633 ± 0.0058	0.1634 ± 0.0058	
$\log g_2(cgs)$	5.015 ± 0.014	5.0214 ± 0.0076	5.075 ± 0.015	5.066 ± 0.019	
$T_{eff,2}(K)$	3027 ± 88	2982 ± 71	3220 ± 70	3143 ± 66	

Table 5. The derived parameters for EBLM J0546-18 and EBLM J0719+25 using *CHEOPS* and *TESS* light curve fits with eclipse depths being in the relevant instrumental bandpass.

	EBLM J	0546-18	EBLM J	EBLM J0719+25		
	CHEOPS	TESS	CHEOPS	TESS		
Model parameters						
$T_0(BJD)$	2203.71457 ± 0.00027	2174.98660 ± 0.00032	2216.39007 ± 0.00024	2559.38262 ± 0.00019		
P (days)	3.191919 (fixed)	3.191919 ± 0.000034	7.456295 (fixed)	7.456295 ± 0.000045		
D	0.0239 ± 0.0018	0.02328 ± 0.00081	0.02145 ± 0.00051	0.02092 ± 0.00017		
W	0.0415 ± 0.0016	0.04020 ± 0.00047	0.02491 ± 0.00029	0.02456 ± 0.00018		
b	0.777 ± 0.040	0.824 ± 0.013	0.498 ± 0.033	0.520 ± 0.016		
f_c	0.0 (fixed)	0.0 (fixed)	-0.2589 ± 0.0069	-0.2588 ± 0.0053		
f_s	0.0 (fixed)	0.0(fixed)	-0.116 ± 0.023	-0.139 ± 0.022		
L	$(11.0 \pm 1.3) \times 10^{-4}$	$(17.6 \pm 1.2) \times 10^{-4}$	$(6.4 \pm 1.2) \times 10^{-4}$	$(9.32 \pm 0.65) \times 10^{-4}$		
h_1	$0.44 \pm 0.14^*$	0.719 ± 0.100	0.731 ± 0.020	0.813 ± 0.013		
h_2	0.31 ± 0.14	0.37 ± 0.24	0.24 ± 0.24	0.56 ± 0.19		
Derived parameters	A \ Y					
R_2/R_1	0.1546 ± 0.0059	0.1526 ± 0.0027	0.1465 ± 0.0018	0.144625 ± 0.000593		
R_1/a	0.1533 ± 0.0057	0.1569 ± 0.0026	0.0757 ± 0.0017	0.076857 ± 0.001019		
R_2/a	0.0223 ± 0.0014	0.02361 ± 0.00034	0.01076 ± 0.00033	0.010941 ± 0.000176		
i(°)	83.17 ± 0.54	82.58 ± 0.22	87.84 ± 0.19	87.711 ± 0.100		
e	0.0	0.0	0.0807 ± 0.0041	0.086242 ± 0.003542		
$\omega(^{\circ})$	_	_	-155.9 ± 4.6	-151.8 ± 4.3		
Absolute parameters	7					
a(AU)	0.04587 ± 0.00080	0.04586 ± 0.00080	0.0802 ± 0.0014	0.0801 ± 0.0014		
$R_2(R_{\odot})$	0.233 ± 0.013	0.2356 ± 0.0072	0.1912 ± 0.0060	0.1915 ± 0.0044		
$M_2(M_{\odot})$	0.2129 ± 0.0075	0.2131 ± 0.0075	0.1584 ± 0.0056	0.1583 ± 0.0056		
$\log g_2(cgs)$	5.029 ± 0.047	5.020 ± 0.021	5.075 ± 0.023	5.073 ± 0.012		
T _{eff,2} (K)	3409 ± 111	3332 ± 90	3208 ± 89	3063 ± 40		

*The limb darkening parameters are not well constrained from CHEOPS data for EBLM J0546-18 (see discussion in Sec. 4.2.1.)

Table 6. The derived parameters for EBLM J2359+44 using *CHEOPS* and *TESS* light curve fits with eclipse depths being in the relevant instrumental bandpass.

	EBLM J	2359+44
	CHEOPS	TESS
Model parameters		
$T_0(BJD)$	1977.85239 ± 0.00015	1773.4230 ± 0.0027
P (days)	11.3627 (fixed)	11.3627 ± 0.0027
D	0.02997 ± 0.00016	0.03015 ± 0.00023
W	0.025946 ± 0.000091	0.02611 ± 0.00017
b	0.096 ± 0.024	0.141 ± 0.033
f_c	-0.05175 ± 0.00032	-0.05242 ± 0.00053
f_s	-0.68888 ± 0.00071	-0.68906 ± 0.00072
L	$(8.91 \pm 0.63) \times 10^{-4}$	$(20.21 \pm 0.98) \times 10^{-4}$
h_1	0.7754 ± 0.0043	0.8393 ± 0.0093
h_2	0.61 ± 0.13	0.60 ± 0.19
Derived parameters		
R_2/R_1	0.17311 ± 0.00045	0.17363 ± 0.00067
R_1/a	0.06971 ± 0.00033	0.07040 ± 0.00066
R_2/a	0.011990 ± 0.000077	0.01207 ± 0.00015
i(°)	89.619 ± 0.098	89.43 ± 0.14
e	0.47724 ± 0.00098	0.47755 ± 0.00099
$\omega(^{\circ})$	-94.30 ± 0.027	-94.350 ± 0.044
Absolute parameters		
a(AU)	0.1144 ± 0.0020	0.1144 ± 0.0020
$R_2(R_{\odot})$	0.2963 ± 0.0058	0.3001 ± 0.0064
$M_2(M_{\odot})$	0.293 ± 0.010	0.293 ± 0.010
$\log g_2(cgs)$	4.9602 ± 0.0049	4.9490 ± 0.0089
$T_{\rm eff,2}(K)$	3465 ± 46	3513 ± 41

were about one order of magnitude larger then obtained from the RV fitting alone. We conclude that even for high precision *CHEOPS* light curves, (i) radial velocity measurements are essential to derive precise radii for low mass eclipsing binaries and (ii) our analysis method does not allow to constrain the orbital eccentricity from the light curves better then from radial velocity measurements.

4.2 Comparison to TESS

For all targets, we compared our results from TESS light curve fitting with the CHEOPS results. Both instruments comprise different passbands with the TESS having an redder effective wavelength of 745.6 nm compared to CHEOPS with 581.1 nm⁸. In this, we do not compare the limb darkening parameters and absolute eclipse depths, since these depend on the instrumental passband. The secondary eclipses are thus 1.5 to 2.5 times deeper in TESS, compared to CHEOPS. We find a good agreement on the derived radius ratio, inclination and relative primary radii R_1/a (<1%). As discussed in the previous section, using radial velocity priors is essential to derive precise radii for the M-dwarfs. We find that keeping fc and fs as free parameters results in 3-6% smaller radii for TESS light curves (for EBLM J0540-17 and EBLM J0719+25), compared to CHEOPS. Using similar radial velocity priors (see chapter 3.1), we find that the derived radii and surface gravity for the M-dwarfs agree well for all targets (on average within 0.9% and 0.15% respectively) between TESS and CHEOPS. We find that the uncertainties of the derived parameters from TESS light curves are of a similar order, compared to CHEOPS results. TESS is in favour, for relatively bright secondary companions with deep secondary eclipse and for targets with short orbital periods and thus, many eclipses covered during the monitoring. We find that the effective temperature of the M-dwarfs, derived from TESS light curves is in agreement with our CHEOPS value for EBLM J2359+44, but about 2-4% cooler for our other targets. We included the result from SW21 for EBLM J1934-42 to analyse for any systematic difference between the effective temperature of the M-dwarf, derived with TESS relative to CHEOPS. We modelled a constant difference between two instruments using EMCEE to take the RV-jitter of the effective temperatures of both TESS and CHEOPS into account by weighting the fit by the log-likelihood function. The offset from our sample of six stars results in a slightly lower (1.11 \pm 0.99%) temperature for TESS light curves with a remaining jitter of 0.0076%.

The small discrepancy in $T_{\rm eff,2}$ might be caused by an underestimation of the secondary eclipse depth (L). In Sec 3.3 we have discussed that reflected light might lead to an underestimated depth of the secondary eclipse. Nevertheless, this effect affects both passbands of *CHEOPS* and *TESS* in a comparable level and only for the shortest period binaries in our sample. Thus, reflection can not explain this discrepancy. Possible explanations might be uncertainties introduced by the stellar model we used to derive the temperature from the surface brightness, or stellar activity of the primary star, linked to stellar spots which are not accounted for in the eclipse model, we have used.

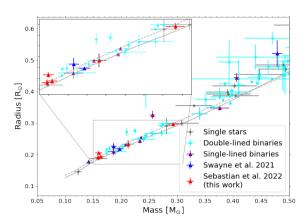
4.2.1 Limb darkening parameters

For our CHEOPS and TESS fits, we kept the limb darkening parameters h₁ and h₂ free. To compare our results, we derived expected limb darkening parameters for EBLM J0239-20, EBLM J0540-17, EBLM J0546-18, and EBLM J0719+25 by interpolating the tables for the TESS bandpass and Kepler passband (for CHEOPS data respectively) published in Maxted (2018) using the stellar parameters T_{eff.1}, log g₁, and [Fe/H] as listed in Table 2, and applying an offset $(h_1 + 0.01)$ and $h_2 - 0.045$; (Maxted 2018)). This method did not converge for the hottest star in our sample EBLM J2359+44 since its effective temperature exceeds the tabulated temperature range. Thus, we used the other four targets for this comparison. The expected limb darkening parameters are listed in Table C1. We find that h₁ agrees on average well with differences of a few percent, while we find larger discrepancies for h2 in the order of several 10 percent similarly in the CHEOPS and TESS data sets. This finding, as well as the derived uncertainties follow the trend from Maxted (2018), (Fig 4) for h₂ to be about one order of magnitude less constrained than h₁. We find some cases of larger uncertainties in CHEOPS light curve fits. EBLM J0546-18 we derive about 31% uncertainty for h₁ and the derived parameter, differs more than 70% from the expectations. This is not surprising, given the large impact parameter which does not allow to constrain the limb darkening parameters for this star. We have repeated the CHEOPS and TESS fits for these four targets, using the expected limb darkening parameters as priors, but found that introducing these priors will neither improve the fit, nor has it any significant impact on the derived M-dwarf parameters. We, thus, present in Table 5 the derived parameters without priors for h₁ and h₂, noting that the corresponding values are less well constrained with CHEOPS compared to TESS.

4.3 Mass-radius diagram

The main goal of the CHEOPS programme is to build a well defined mass-radius diagram for stars below the fully convective bound-

⁸ Filter profiles and effective wavelengths can be accessed using the VSO Filter Profile Service.



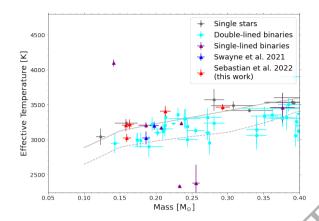


Figure 1. Left: Mass-radius diagram for low mass stars. Triangles: Single lined eclipsing binaries, with *CHEOPS* programme targets highlighted in red and blue. Gray, and Cyan squares: single stars and double lined binaries from literature with measured mass, radius, and effective temperature. The zoom in section highlights the MIST model tracks for [Fe/H]=0, grey line, and [Fe/H]=0.25, grey dotted line. Right: Mass-effective temperature diagram of the same data set compared to same MIST models.

ary. In Fig. 1 we show our five targets together with the theoretical mass relation from MIST stellar models for 1 Gyr stars of solar metallicity ([Fe/H]=0.0) as well as for slightly more metal rich stars ([Fe/H]=0.25) (Dotter 2016; Choi et al. 2016; Paxton et al. 2011). Similarly to SW21, we compiled a comparison sample of precisely measured low mass stars from literature, classified in single stars, double lined binaries, and single lined binaries. (Carter et al. 2011; Nefs et al. 2013; Gillen et al. 2017; Parsons et al. 2018; Smith et al. 2021; Swayne et al. 2021). We compared the radii with both the MIST and the Exeter/Lyon (Baraffe et al. 2015) models for solar metallicity. The M-dwarf radius for EBLM J0239-20 is $11.0\pm2.6\%$ ($12.5\pm2.6\%$) larger for the MIST (and Exeter/Lyon) model, the others are on average $2.6 \pm 1.3\%$ ($3.5 \pm 1.3\%$) larger compared to both models. Despite most of our targets being within the uncertainties in agreement with the theoretical radii, we observe that they follow the trend of very low, mass stars to be slightly larger than predicted by models. In Fig. 1 we also show the effective temperature of our five M-dwarfs, the result from SW21, as well as the same literature sample. Our targets effective temperatures follow the overall trend of low mass stars. We note that EBLM J0239-20, similarly to EBLM J1934-42 (blue triangle from SW21) have a slightly higher metallicity ([M/H] > 0.2). Both stars are slightly larger and cooler, compared to models for stars with solar metallicity. As shown in Fig. 1 this trend is predicted by the MIST models for more metal rich stars. But also in this case, both stars are slightly larger than predicted by models for higher metallicity stars. Fig. 1 shows three single lined stars from literature with measured M-dwarf effective temperatures being outliers of more than 500 K compared to model predictions. These are KIC 1571511B (Ofir et al. 2012) as well as SAO 106989 and HD 24465 (Chaturvedi et al. 2018). Populating the low-mass main-sequence with M-dwarfs having precise effective temperature measurements will help us to constrain possible trends for low-mass dwarfs. This is one of the main goals of our CHEOPS programme.

Magnetic activity of the primary star, like spot crossing is not accounted for in our eclipse model, thus, can affect the size determinations of the M-dwarfs. We used the *TESS* light curves to search for variability linked to magnetic activity, like rotational pattern and flares. No flares have been found in the *TESS* data set. EBLM J0239-20 shows a variable modulation of 2-3% close to the orbital period,

most probably linked to stellar activity aligned with the rotational period of the G-dwarf. All our other targets show no or small variability of less than 1 %. Since we found a good agreement between the M-dwarf radii in the different passbands of TESS and CHEOPS, we conclude that stellar activity can only have a minor (< 1%) effect on the derived M-dwarf radius for the five stars, analysed in this work. Depending on the actual contrast between the primary star and the M-dwarf the contribution of the M-dwarf is between 300 and 1200 ppm in CHEOPS data. From this we can exclude large flares with exceed relative intensities of 25 to 100 % compared to the Mdwarfs average brightness. M-dwarfs with such flaring activity exist but account only for about 10% of the flaring M-dwarfs found in TESS (Günther et al. 2020). We can assume that the M-dwarf rotation period is synchronised with the orbital period, since the tidal synchronisation timescale for EBLM systems is about 1 Gyr or less (Barker 2020). Thus the M-dwarfs are expected to be fast rotators $(P \le 10 \,\mathrm{d})$, which are expected to show enhanced activity levels (e.g. Morales et al. 2010; Wright et al. 2018). Activity induced photometric variations, observed for field M-dwarfs is typically in the order of 1% of the M-dwarfs average brightness (Medina et al. 2020). This results in an expected photometric variability in the order of 10 ppm for active M-dwarfs which is below the detection efficiency in our data.

Reflected light from the primary star (See discussion in Sec 3.3) can cause an underestimated radius of the M-dwarfs. We note that this effect is negligible for the five binaries analysed in this work, as it would result in a relative underestimation of about 100 ppm of the M-dwarfs radius for the shortest period binaries in our sample.

5 SUMMARY

Within the framework of our EBLM project, we initiated a *CHEOPS* observing programme of 23 low-mass stars to measure precise stellar parameters as well as effective temperatures. In this paper, we have analysed high precision *CHEOPS* light curves of primary and secondary eclipses for five eclipsing binaries with low mass companions. Using the qpower2 transit model, of PYCHEOPS, we find an average uncertainty of $3.2 \pm 1.3\%$ for the M-dwarfs radius and $0.4 \pm 0.3\%$ for the M-dwarfs surface gravity. Thus, using precision light curves

allowed us to overcame the larger uncertainties to derive stellar parameters typically involved with high-resolution spectroscopy. We have derived the M-dwarfs effective temperature from the contrast between primary and secondary eclipses and the metallicity from spectroscopic analysis of the primary star, assuming equal metallicities of both components.

This allows us to compare the M-dwarfs parameters to theoretical structural models, like the MIST models. We find that all our Mdwarfs are on average larger, but agree within the uncertainty with the model predictions. This is also true for low-mass M-dwarfs with enhanced metallicity, which follow the predicted trend of having a larger radius as well as a cooler effective temperature. Up to now, the stellar models, as well as our transit model do not include stellar activity. We have analysed TESS light curves for all our five targets and find a good (better than 1%) agreement on the M-dwarf radius in the different passband of both instruments. Given the absence of strong activity indicated variability and flare activity as well as this good agreement, we conclude that stellar activity does not play a strong role in the derived uncertainties for our five stars. This result is of particular importance for more active stars on our CHEOPS programme, where activity induced changes in parameters between the TESS and CHEOPS passbands might need to be accounted for. We have analysed the dependence of derived M-dwarf parameters with priors used in the fit. We find that limb darkening parameters as well as orbital parameters like the eccentricity and the argument of periastron are not well constrained from our model fit. Nevertheless, we find that, other than the limb darkening coefficients, precise orbital parameters, obtained from radial velocity observations are crucial to derive M-dwarf radii better than 5%.

Together with SW21, we increased the sample to eight low-mass stars, with precise measured radii from *CHEOPS* data. Due to the fact that the F,G,K-type primary companions are single lined binaries, that allow high-precision orbital characterisation as well as the determination of precise stellar parameters like metallicity, this survey, once completed, will allow us to empirically shed light on the radius inflation problem for very low mass stars.

ACKNOWLEDGEMENTS

CHEOPS is an ESA mission in partnership with Switzerland with important contributions to the payload and the ground segment from Austria, Belgium, France, Germany, Hungary, Italy, Portugal, Spain, Sweden, and the United Kingdom. The CHEOPS Consortium would like to gratefully acknowledge the support received by all the agencies, offices, universities, and industries involved. Their flexibility and willingness to explore new approaches were essential to the success of this mission. Spectroscopic data were obtained via observing time allocations at OHP awarded by the French PNP (18B.PNP.SAN1, 19A.PNP.SANT). Some of the observations reported in this paper were obtained with the Southern African Large Telescope (SALT). This research has made use of the services of the ESO Science Archive Facility. This research is also supported work funded from the European Research Council (ERC) the European Union's Horizon 2020 research and innovation programme (grant agreement n°803193/BEBOP). PM acknowledges support from STFC research grant number ST/M001040/1. MIS acknowledges support from STFC grant number ST/T506175/1 S.G.S. acknowledge support from FCT through FCT contract nr. CEECIND/00826/2018 and POPH/FSE (EC). SH gratefully acknowledges CNES funding through the grant 837319. YA and MJH acknowledge the support of the Swiss National Fund under grant 200020_172746.

We acknowledge support from the Spanish Ministry of Science and Innovation and the European Regional Development Fund through grants ESP2016-80435-C2-1-R, ESP2016-80435-C2-2-R, PGC2018-098153-B-C33, PGC2018-098153-B-C31, ESP2017-87676-C5-1-R, MDM-2017-0737 Unidad de Excelencia Maria de Maeztu-Centro de Astrobiología (INTA-CSIC), as well as the support of the Generalitat de Catalunya/CERCA programme. The MOC activities have been supported by the ESA contract No. 4000124370. S.C.C.B. acknowledges support from FCT through FCT contracts nr. IF/01312/2014/CP1215/CT0004. XB, SC, DG, MF and JL acknowledge their role as ESA-appointed CHEOPS science team members. ABr was supported by the SNSA. ACC acknowledges support from STFC consolidated grant numbers ST/R000824/1 and ST/V000861/1, and UKSA grant number ST/R003203/1. This project was supported by the CNES. The Belgian participation to CHEOPS has been supported by the Belgian Federal Science Policy Office (BELSPO) in the framework of the PRODEX Program, and by the University of Liège through an ARC grant for Concerted Research Actions financed by the Wallonia-Brussels Federation. L.D. is an F.R.S.-FNRS Postdoctoral Researcher. This work was supported by FCT - Fundação para a Ciência e a Tecnologia through national funds and by FEDER through COMPETE2020 Programa Operacional Competitividade e Internacionalizacão by these grants: UID/FIS/04434/2019, UIDB/04434/2020, UIDP/04434/2020, PTDC/FIS-AST/32113/2017 & POCI-01-0145-FEDER- 032113, PTDC/FIS-AST/28953/2017 & POCI-01-0145-FEDER-028953, PTDC/FIS-AST/28987/2017 & POCI-01-0145-FEDER-028987, O.D.S.D. is supported in the form of work contract (DL 57/2016/CP1364/CT0004) funded by national funds through FCT. B.-O.D. acknowledges support from the Swiss National Science Foundation (PP00P2-190080). This project has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (project Four Aces. grant agreement No 724427). It has also been carried out in the frame of the National Centre for Competence in Research PlanetS supported by the Swiss National Science Foundation (SNSF). DE acknowledges financial support from the Swiss National Science Foundation for project 200021_200726. MF and CMP gratefully acknowledge the support of the Swedish National Space Agency (DNR 65/19, 174/18). DG gratefully acknowledges financial support from the CRT foundation under Grant No. 2018.2323 "Gaseousor rocky? Unveiling the nature of small worlds". M.G. is an F.R.S.-FNRS Senior Research Associate. KGI is the ESA CHEOPS Project Scientist and is responsible for the ESA CHEOPS Guest Observers Programme. She does not participate in, or contribute to, the definition of the Guaranteed Time Programme of the CHEOPS mission through which observations described in this paper have been taken, nor to any aspect of target selection for the programme. This work was granted access to the HPC resources of MesoPSL financed by the Region IIe de France and the project Equip@Meso (reference ANR-10-EQPX-29-01) of the programme Investissements d'Avenir supervised by the Agence Nationale pour la Recherche. ML acknowledges support of the Swiss National Science Foundation under grant number PCEFP2_194576. LBo, GBr, VNa, IPa, GPi, RRa, GSc, VSi, and TZi acknowledge support from CHEOPS ASI-INAF agreement n. 2019-29-HH.0. This work was also partially supported by a grant from the Simons Foundation (PI Queloz, grant number 327127). IRI acknowledges support from the Spanish Ministry of Science and Innovation and the European Regional Development Fund through grant PGC2018-098153-B- C33, as well as the support of the Generalitat de Catalunya/CERCA programme. GyMSz acknowledges the support of the Hungarian National Research, Development and Innovation Office (NKFIH) grant K-125015, a PRODEX Institute Agreement between the ELTE Eötvös Loránd University and the European Space Agency (ESA-D/SCI-LE-2021-0025), the Lendület LP2018-7/2021 grant of the Hungarian Academy of Science and the support of the city of Szombathely. V.V.G. is an F.R.S-FNRS Research Associate. NAW acknowledges UKSA grant ST/R004838/1. V.K. acknowledges support from NSF award AST2009501. This work was also supported by the STFC PATT Travel grant ST/S001301/1. We thank the reviewer for valuable comments and suggestions.

DATA AVAILABILITY

All CHEOPS data and data products are publicly available via the Data Analysis Center for Exoplanets web platform. This paper includes data collected by the TESS mission, which is publicly available from the Mikulski Archive for Space Telescopes (MAST) at the Space Telescope Science Institute (STScI) (https://mast.stsci.edu). Funding for the TESS mission is provided by the NASA Explorer Program directorate. STScI is operated by the Association of Universities for Research in Astronomy, Inc., under NASA contract NAS 5-26555. We acknowledge the use of public TESS Alert data from pipelines at the TESS Science Office and at the TESS Science Processing Operations Center. SOPHIE high-resolution spectra are available trough the data archives of the Observatoire de Haute-Provence via http://atlas.obs-hp.fr/. Programme ID were 18B.PNP.SAN1, and 19A.PNP.SANT.

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APPENDIX D: TESS FITS

APPENDIX E: CHEOPS FITS

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APPENDIX A: DECORRELATION PARAMETERS FITTED FROM CHEOPS FITS

APPENDIX B: RADIAL VELOCITY MEASUREMENTS

COEFFICIENTS

APPENDIX C: EXPECTED LIMB DARKENING



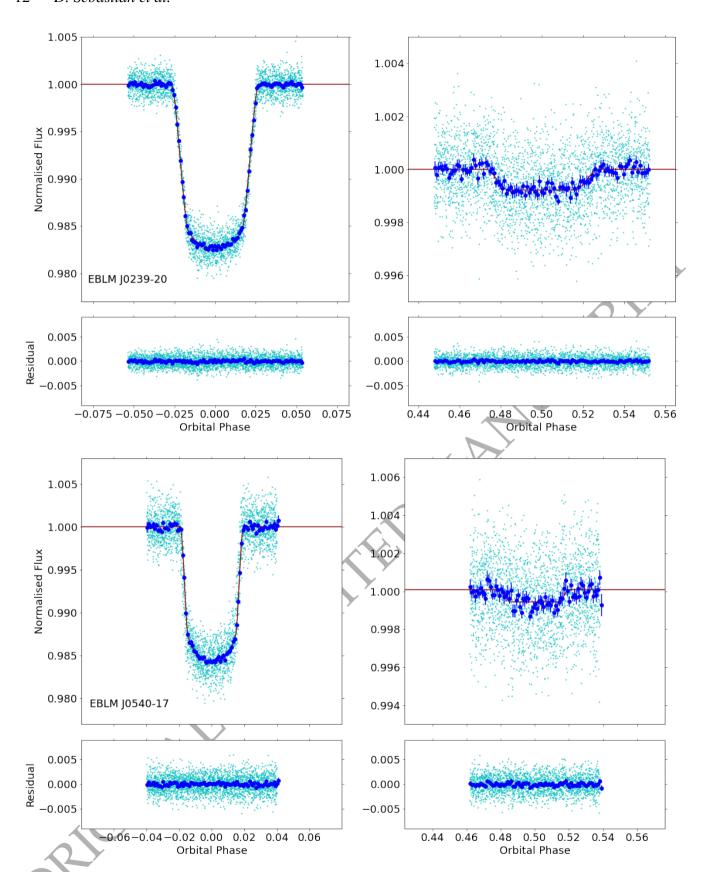


Figure D1. Fitted TESS light curves of all targets in phase intervals around the primary and secondary eclipse events. The observed data points are shown in cyan. The fitted light curve is shown in red. The residual of the fit is displayed the fitted curves.

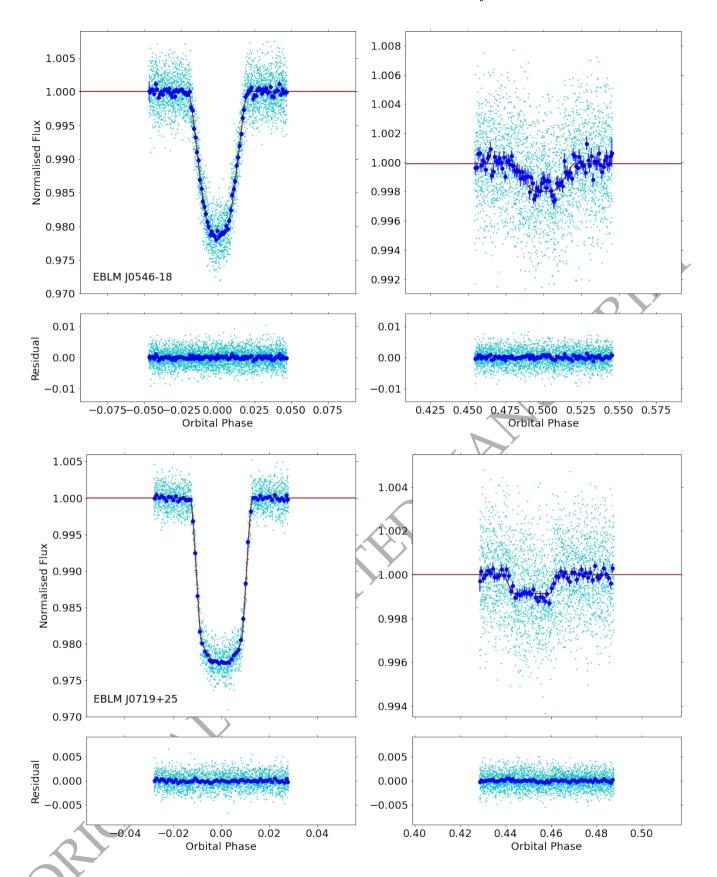


Figure D2. Fitted TESS light curves of all targets in phase intervals around the primary and secondary eclipse events. The observed data points are shown in cyan. The fitted light curve is shown in red. The residual of the fit is displayed the fitted curves.

Table A1. Decorrelation parameters fitted from *CHEOPS* multivisit analysis for each visit (in the same order as in Table 1). The parameters are: Image background level (dfdbg), PSF centroid position (dfdx and dfdy), time (dfdt), and aperture contamination (dfdcontam).

Target	Eclipse	$ dfdbg (10^{-3}) $	dfdx (10 ⁻⁴)	$ dfdy (10^{-3}) $		dfdcontam (10 ⁻³)
EBLM J0239-20	primary	-	_	_	_	_
	secondary	1.57 ± 0.90	_	0.311 ± 0.085	2.924 ± 0.029	_
	secondary	1.21 ± 0.23	_	_	1.680 ± 0.029	_
EBLM J0540-17	primary	1.20 ± 0.82	7.33 ± 1.79	_	-0.31 ± 0.43	_
	secondary	0.71 ± 0.77	_	_	_	_
	secondary	_	_	-0.51 ± 0.14	0.163 ± 0.036	_
	secondary	_	5.95 ± 1.71	-0.87 ± 0.17	_	_
EBLM J0546-18	primary	4.80 ± 0.87	_	0.78 ± 0.23	_	-1.73 ± 0.56
	secondary	_	_	_	_	-1.59 +/- 0.83
	secondary	2.85 ± 0.66	11.32 ± 2.51	_	1.367 ± 0.079	_
EBLM J0719+25	primary	_	_	_	-0.496 ± 0.060	_
	secondary	1.22 ± 0.93	_	_	0.291 ± 0.061	_
	secondary	_	_	_	_	-
EBLM J2359+44	secondary	0.83 ± 0.40	_	0.208 ± 0.088	_	-0.48 ± 0.27
	primary	0.83 ± 0.26	_	_	_	- 🔎

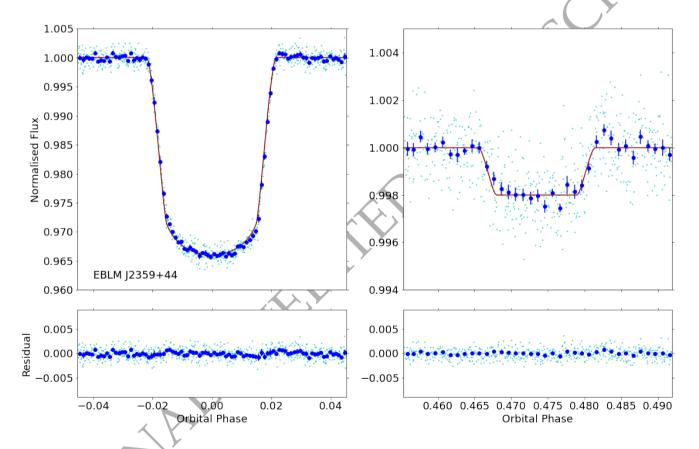


Figure D3. Fitted *TESS* light curves of all targets in phase intervals around the primary and secondary eclipse events. The observed data points are shown in cyan. The fitted light curve is shown in red. The residual of the fit is displayed the fitted curves.

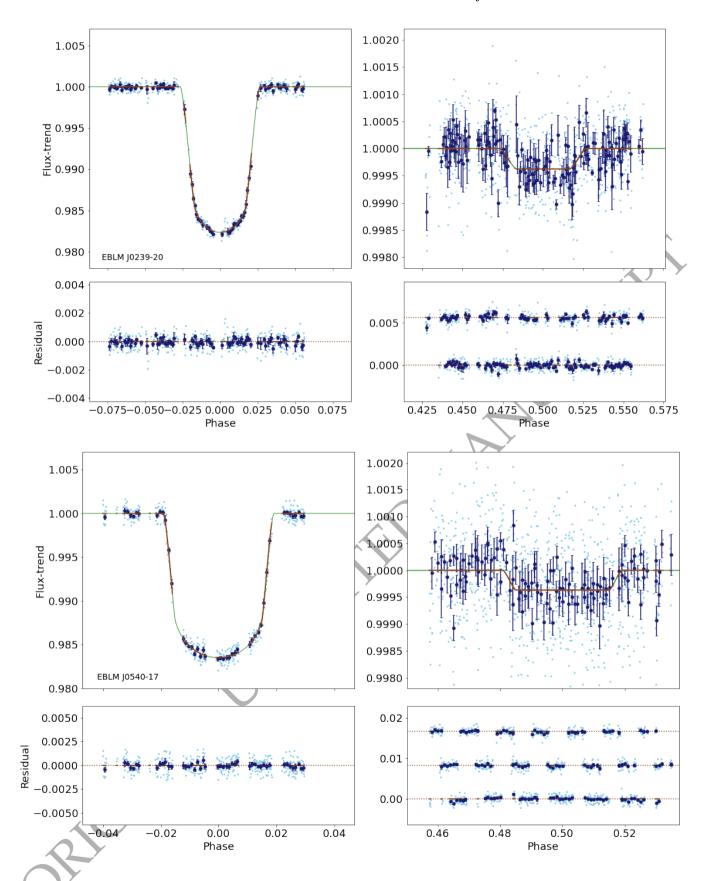


Figure E1. Fitted *CHEOPS* light curves of all targets in phase intervals around the primary and secondary eclipse events. The observed data points are shown in cyan. The fitted light curve is shown in red. The residual of the fit is displayed in blue below the fitted curves.

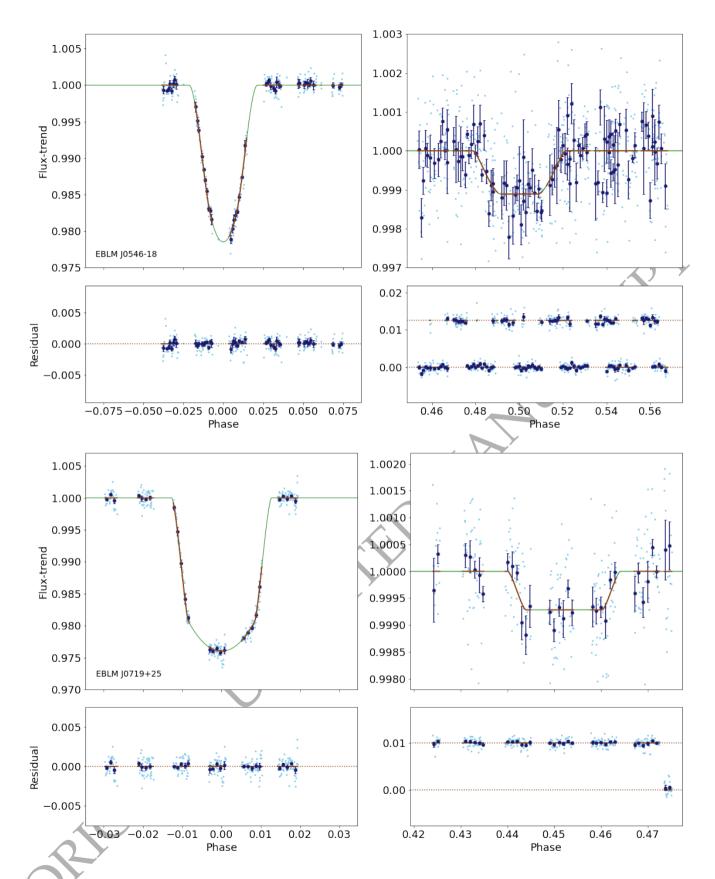


Figure E2. Fitted CHEOPS light curves of all targets in phase intervals around the primary and secondary eclipse events. The observed data points are shown in cyan. The fitted light curve is shown in red. The residual of the fit is displayed in blue below the fitted curves.

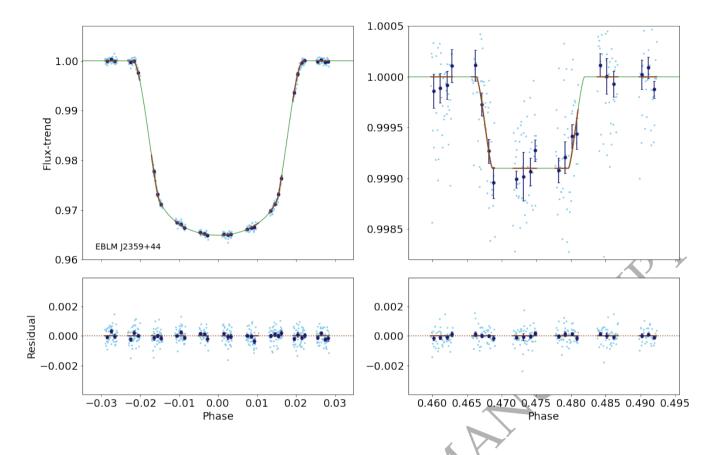


Figure E3. Fitted *CHEOPS* light curves of all targets in phase intervals around the primary and secondary eclipse events. The observed data points are shown in cyan. The fitted light curve is shown in red. The residual of the fit is displayed in blue below the fitted curves.

Table B1. Radial velocity measurements for EBLM J0719+25

BJD - 2400000	RV [km s ⁻¹]	RV error $[km s^{-1}]$	Source
58436.57258	-5.9492	0.0079	SOPHIE
58438.59676	12.5703	0.0057	SOPHIE
58536.40291	11.1258	0.0058	SOPHIE
58538.42658	-9.091	0.012	SOPHIE
58542.39085	10.1391	0.0047	SOPHIE
58562.39379	-15.9404	0.0073	SOPHIE
58566.37826	10.2797	0.0053	SOPHIE
58761.63689	-3.306	0.011	SOPHIE

Table B2. Radial velocity measurements for EBLM J2359+44

BJD - 2400000	$RV [km s^{-1}]$	RV error $[km s^{-1}]$	Source
53310.6391	-19.07	0.42	Poleski et al.
53311.7990	-26.36	0.50	Poleski et al.
58436.31776	-33.537	0.011	SOPHIE
58438.40839	2.8147	0.0086	SOPHIE
58685.56693	-29.4759	0.012	SOPHIE
58704.54724	-8.063	0.014	SOPHIE
58729.61888	-20.846	0.013	SOPHIE
58734.5406	11.81	0.015	SOPHIE
58754.47118	-33.987	0.015	SOPHIE
58765.46162	-31.893	0.011	SOPHIE
59030.57795	10.110	0.011	SOPHIE
59043.50347	1.726	0.014	SOPHIE
59045.53151	-9.040	0.012	SOPHIE
59071.56389	-27.920	0.012	SOPHIE
59077.5554	1.898	0.012	SOPHIE
59094.51791	-29.440	0.011	SOPHIE
59100.57485	0.226	0.012	SOPHIE

Table C1. Expected limb darkening coefficients derived for *TESS* and *CHEOPS* passbands.

Target	СНЕС	OPS .	TES	SS
	h_1	h_2	h_1	h ₂
EBLM J0239-20	0.743±0.012	0.40±0.05	0.798±0.012	0.39±0.05
EBLM J0540-17	0.773 ± 0.011	0.41 ± 0.05	0.826±0.011	0.38 ± 0.05
EBLM J0546-18	0.771 ± 0.011	0.41 ± 0.05	0.822±0.011	0.37±0.05
EBLM J0719+25	0.754±0.011	0.41 ± 0.05	0.808±0.011	0.39±0.05

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