



# A SOFIA FORCAST Grism Study of the Mineralogy of Dust in the Winds of Proto-planetary Nebulae: RV Tauri Stars and SRd Variables

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## Abstract

We present a SOFIA FORCAST grism spectroscopic survey to examine the mineralogy of the circumstellar dust in a sample of post-asymptotic giant branch (post-AGB) yellow supergiants that are believed to be the precursors of planetary nebulae. Our mineralogical model of each star indicates the presence of both carbon-rich and oxygen-rich dust species—contrary to simple dredge-up models—with a majority of the dust in the form of amorphous carbon and graphite. The oxygen-rich dust is primarily in the form of amorphous silicates. The spectra do not exhibit any prominent crystalline silicate emission features. For most of the systems, our analysis suggests that the grains are relatively large and have undergone significant processing, supporting the hypothesis that the dust is confined to a Keplerian disk and that we are viewing the heavily processed, central regions of the disk from a nearly face-on orientation. These results help to determine the physical properties of the post-AGB circumstellar environment and to constrain models of post-AGB mass loss and planetary nebula formation.

*Key words:* astrochemistry – binaries: general – circumstellar matter – stars: AGB and post-AGB – stars: evolution

*Supporting material:* extended figures

## 1. Introduction

RV Tauri and yellow semi-regular (SRd) variables are two classes of post-asymptotic giant branch (post-AGB) stars that lie along the horizontal track on the Hertzsprung–Russell (H-R) diagram linking AGB stars to planetary nebulae (PNe). They are thought to be the immediate precursors of PNe and have been termed “proto-planetary nebulae” (PPNe).<sup>12</sup> RV Tauri stars are characterized by semi-regular, bimodal variability (possibly resulting from interaction with a binary companion; Percy 1993; Waelkens & Waters 1993; Fokin 1994), a high mass-loss rate, and often a prominent infrared (IR) excess. SRd variables are similar to RV Tauri stars in many respects but are probably single-star systems, as indicated by the absence of regular pulsations (Percy & Ursprung 2006).

RV Tauri stars are a loosely defined subclass of population II Cepheid variables named after the prototype RV Tau. They are defined as luminous (I-II) mid-F to K supergiants with a typical mass of  $\sim 0.7 M_{\odot}$  (Tuchman et al. 1993; Fokin 1994) that show alternating deep and shallow minima in their light curves (Preston et al. 1963). They have formal periods (defined as the

time between successive deep minima) between 30 and 150 days, but cycle-to-cycle variability is common, and the amplitudes may reach up to 4 mag in *V* (Percy 1993). RV Tauri stars are divided into two photometric classes (“a” and “b”) based on their light curves (Kukarkin 1958). The RVA class contains constant mean magnitude stars, and the RVb class contains stars that have a varying mean magnitude with a period of 600–1500 days. There are several possible explanations for these light variations. One explanation for the alternating minima is that there is a resonance between the fundamental period and the first overtone (Takeuti & Petersen 1983; Shenton et al. 1992; Tuchman et al. 1993; Fokin 1994). Another possibility is that the light variations are due to a geometrical projection effect where the pulsating star is periodically obscured by a circumbinary disk (Van Winckel et al. 1999; Maas et al. 2002).

Preston et al. (1963) classified the RV Tauri stars into three spectroscopic classes (“A,” “B,” “C”). RVA stars are spectral type G–K, and show strong absorption lines and normal CN or CH bands while TiO bands sometimes appear at photometric minima. RVB stars are generally hotter spectral types, weaker lined, and show enhanced CN and CH bands. RVC stars are also weak lined but show normal CN and CH molecular bands. There is no correlation between the photometric and spectroscopic classes.

It has long been known (Gehrz & Woolf 1970; Gehrz 1972; Gehrz & Ney 1972) that some RV Tauri stars (e.g., AC Her, U Mon, R Sct, R Sge) show very strong thermal IR emission from

<sup>11</sup> Michael Jura died on 2016 January 30 while this manuscript was being drafted. He participated in writing the proposals to gather the data, and was aware of the importance of the results at the time of his death.

<sup>12</sup> The phrase “proto-planetary” is also widely used by the exoplanetary and planet formation communities to refer to dusty disks around young stars. Also note that in the literature the terms pre-planetary or proto-planetary nebulae have been used interchangeably with the term post-AGB objects.

**Table 1**  
Properties of the RV Tauri and SRd Variables in This Survey

Name	Type	Spectral Type	Period (day) <sup>a</sup>	[Fe/H] <sub>0</sub> <sup>b</sup>	PC <sup>c</sup>	SC <sup>c</sup>	SED <sup>d</sup>	$T_{\text{eff}}$ (K)	Binarity <sup>e</sup>	Chemical Type <sup>f</sup>	References
TW Cam	RV	F8IbG8Ib	87	-0.40	a	A	Disk	4800			1
UY CMa	RV	G0	114	-0.50	a	B		5500			2
o <sup>1</sup> Cen	SRd	G3Ia0	200								3
RU Cen	RV	A7IbG2pe	65	-1.10	a	B	Disk	6000	Y		4, 5
SX Cen	RV	F5G3/5Vp	33	-0.30	b	B	Disk	6250	Y		4, 5
SU Gem	RV	F5M3	50	0.00	b	A	Disk	5250			6
AC Her	RV	F2pIbK4e	75	-0.90	a	B	Disk	5900	Y	O	7
V441 Her	SRd	F2Ibe	70				Disk		Y	O	8, 9
U Mon	RV	F8IbeK0pIb	91	-0.50	b	A	Disk	5000	Y	O	1, 10
CT Ori	RV	F9	136	-0.60	a	B	Disk	5500			10, 11
TV Per	SRd	K0	358								12
TX Per	RV	Gp(M2)K0e(M2)	78	-0.60	a	A		4250			6
AR Pup	RV	F0IF8I	76	0.40	b	B	Disk	6000		O	10, 13
R Sge	RV	G0IbG8Ib	71	0.10	b	A	Disk	5100			13
AI Sco	RV	G0K2	71	-0.30	b	A	Disk	5300		C?	2, 10
R Sct	RV	G0IaeK2p(M3)Ibe	147	-0.20	a	A	Uncertain	4500			1
RV Tau	RV	G2IaeM2Ia	79	-0.40	b	A	Disk	4500		C	1
V Vul	RV	G4eK3(M2)	76	0.10	a	A	Disk	4500			2, 6

**Notes.**

<sup>a</sup> Pulsation period in days.

<sup>b</sup> The estimated initial metallicity obtained via the Zn or S abundance (Gezer et al. 2015).

<sup>c</sup> Photometric class (PC) and spectroscopic class (SC).

<sup>d</sup> Spectral energy distribution classification from Gezer et al. (2015).

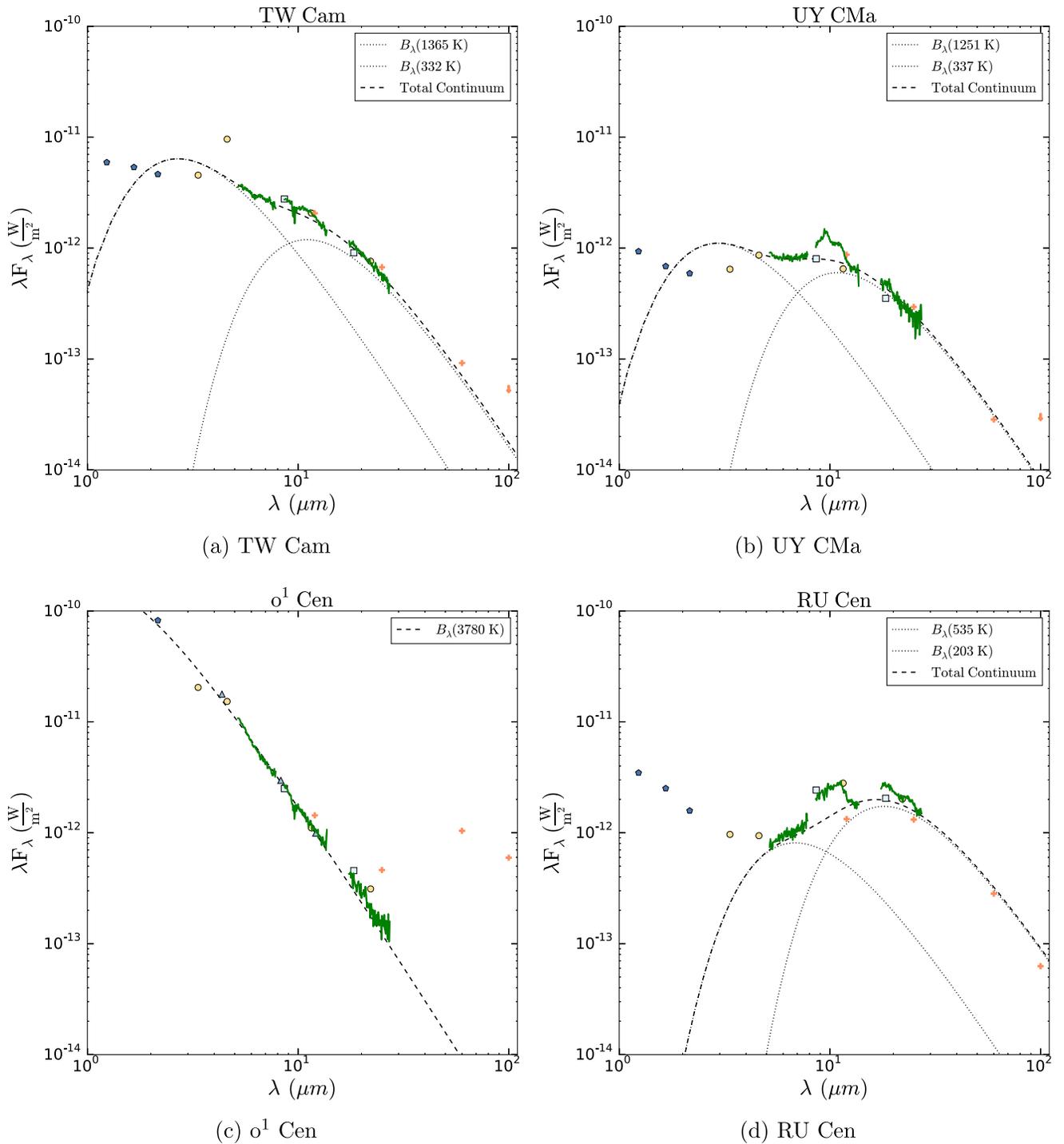
<sup>e</sup> Y indicates confirmed binarity based on radial velocity measurements. Confirming binarity using this method is difficult because the photospheres of these variables have large amplitude radial pulsations.

<sup>f</sup> Stellar chemical type from He et al. (2014) and references therein.

**References.** (1) Giridhar et al. (2000), (2) Giridhar et al. (2005), (3) O’Connell (1961), (4) Maas et al. (2002), (5) Maas et al. (2005), (6) Rao & Giridhar (2014), (7) Giridhar et al. (1998), (8) Waters et al. (1993), (9) de Ruyter et al. (2006), (10) Kiss et al. (2007), (11) Gonzalez et al. (1997a), (12) Payne-Gaposchkin (1952), (13) Gonzalez et al. (1997b).

circumstellar dust. Observations by the *Infrared Astronomical Satellite (IRAS)* confirmed these previous detections and increased the sample size. *IRAS* detected considerable cool, circumstellar dust around many of the RV Tauri stars, which has been interpreted as being due to strong, dusty mass loss during AGB evolution (Jura 1986). From CO observations, Alcolea & Bujarrabal (1991) suggest a mass loss rate of  $4 \times 10^{-7}$  to  $7 \times 10^{-6} M_{\odot} \text{ yr}^{-1}$  within the last  $10^2$  to  $10^3$  years for most of the RV Tauri stars. Because of their position on the H-R diagram, variability, high mass-loss rate and rarity (about 110 are known), it is generally believed that RV Tauri stars represent a relatively short-term, unstable transitional phase between the AGB and PNe phases of solar-mass stars. If RV Tauri stars are assumed to be evolving on the post-AGB track, models predict that they spend about 200 years as variables and take about 1000 years to go from AGB to PNe (Schoenberner 1983; Percy et al. 1991; Fokin 1994), however some (e.g., R Sct) are “lazy,” spend more time in the RV Tauri stage as evidenced by their mass-loss history, have higher CO emission, and show a relatively low IR excess in the 5–20  $\mu\text{m}$  range but a larger IR excess for  $\lambda \gtrsim 50 \mu\text{m}$  (Bujarrabal et al. 1988; Alcolea & Bujarrabal 1991). The lack of a near-IR excess and the presence of a large mid- and far-IR excess is evidence for a thick and extended dust envelope that is relatively cool. The detection of SiO around R Sct could be an indication of ongoing, weak mass loss (Bujarrabal et al. 1989). Alcolea & Bujarrabal (1991) estimate that  $\sim 1/5$  RV Tauri stars exhibit this “lazy” evolutionary behavior.

Interestingly, RV Tauri stars do not have the same high C and s-process overabundances that are characteristic of other post-AGB stars. Instead, their photospheres are “depleted” of refractory elements (Gonzalez et al. 1997a, 1997b; Giridhar et al. 1998, 2000, 2005; Van Winckel et al. 1998; Maas et al. 2002, 2005; Molster et al. 2002a, 2002b, 2002c; Deroo et al. 2005; Gielen et al. 2007, 2009; Rao & Giridhar 2014). This phenomenon is not fully understood, but is apparently due to a chemical rather than a nucleosynthetic process (van Winckel 2003). The refractory elements, which have a high dust condensation temperature, are separated from the volatiles as the condensed grains are driven away by radiation pressure in the circumstellar environment. The depleted gas is then re-accreted onto the stellar photosphere, leaving it with a peculiar composition similar to that of the depleted gas in the interstellar medium (ISM) (Hinkle et al. 2007). Waters et al. (1992) suggested that these abundance anomalies are more likely to occur when the dust is trapped in a circumstellar disk. These same depletion patterns are also observed in binary post-AGB stars with circumbinary disks (Van Winckel et al. 1995). This has led to the suggestion that the depleted RV Tauri stars are also binaries with a circumstellar disk (Van Winckel et al. 1999). In their study, de Ruyter et al. (2006) included many RV Tauri stars in the class of post-AGB binaries with disks and suggested that the thermal IR spectrum originates in the Keplerian circumstellar disk. The narrow velocity width lines of the  $^{12}\text{CO}$  and  $^{13}\text{CO}$   $J = 2 \rightarrow 1$  and  $J = 1 \rightarrow 0$  transitions of AC Her, V441 Her, and R Sct and the  $^{12}\text{CO}$   $J = 3 \rightarrow 2$  and  $J = 2 \rightarrow 1$  transitions of AR Pup are



**Figure 1.** Observed SOFIA FORCAST spectrum (green curve) of our sample of stars plotted together with the archival photometry from 2MASS (blue pentagons), MSX (blue triangles), AKARI (light blue squares), WISE (yellow circles), IRAS (orange crosses), Herschel (red diamonds), and the best-fitting Planck functions (black dotted and dashed curves). WISE photometry upper limits are depicted as orange downward arrows.

indicative of a Keplerian disk (Bujarrabal et al. 2013, 2015). The trapping of grains in a disk could provide the environment needed to enable grains to grow to sizes significantly larger than those in the ISM or in the stellar outflows of single stars (Shenton et al. 1995; de Ruyter et al. 2005). The similar class of SRd stars, however, does not show evidence of disks and may evolve from single stars.

Depending on the oxygen and carbon abundances in the circumstellar gas, either oxygen- or carbon-rich dust is formed.

It has been found that many of the RV Tauri stars have very weak CO  $J = 2 \rightarrow 1$  lines, unlike the stronger CO  $J = 2 \rightarrow 1$  lines characteristic of other, more massive, post-AGB stars (Alcolea & Bujarrabal 1991; He et al. 2014; Bujarrabal et al. 2015). It is possible that the unique circumstellar environment of RV Tauri stars either suppresses the formation of CO molecules (possibly through ultraviolet photodissociation) or underexcites them (Alcolea & Bujarrabal 1991; McDonald et al. 2015). However, the weakness may be due

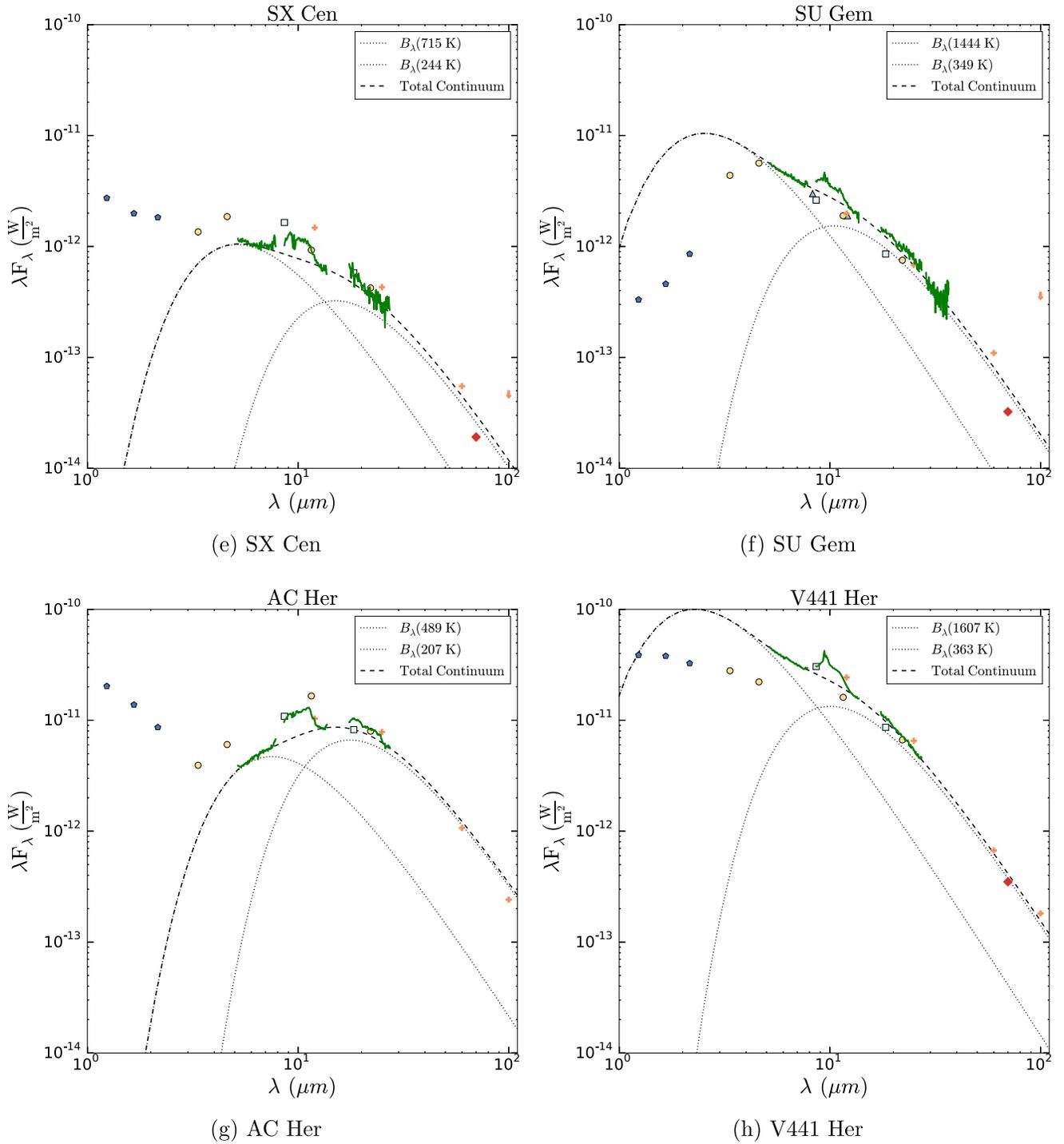


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to the compactness of the disk, as has been suggested for the case of AC Her (Bujarrabal et al. 2015). In oxygen-rich gas, dust such as olivine ( $\text{Mg}_{2(1-x)}\text{Fe}_{2x}\text{SiO}_4$ ) and pyroxene ( $\text{Mg}_{1-x}\text{Fe}_x\text{SiO}_3$ ) is formed. In carbon-rich gas, carbon-rich dust particles such as SiC, amorphous carbon, and possible polycyclic aromatic hydrocarbons (PAHs) are formed.

The amount of crystalline grain material compared to amorphous grain material is generally low,  $\sim 10\%$ – $15\%$ , and is dominant only in rare cases. Crystalline material is also

generally only detected in stars that have experienced high mass-loss rates (above  $10^{-5} M_{\odot} \text{ yr}^{-1}$ ) (Cami et al. 1998; Sogawa & Kozasa 1999; Sylvester et al. 1999; Suh 2002). However, both crystalline and amorphous grains have been detected simultaneously in stellar outflows, both in the present study as well as in others. Gielen et al. (2008, 2009) showed that dust processing in circumstellar disk environments is conducive to creating large, crystalline grains. As the disk is subjected to the hard radiation and stellar wind from the central

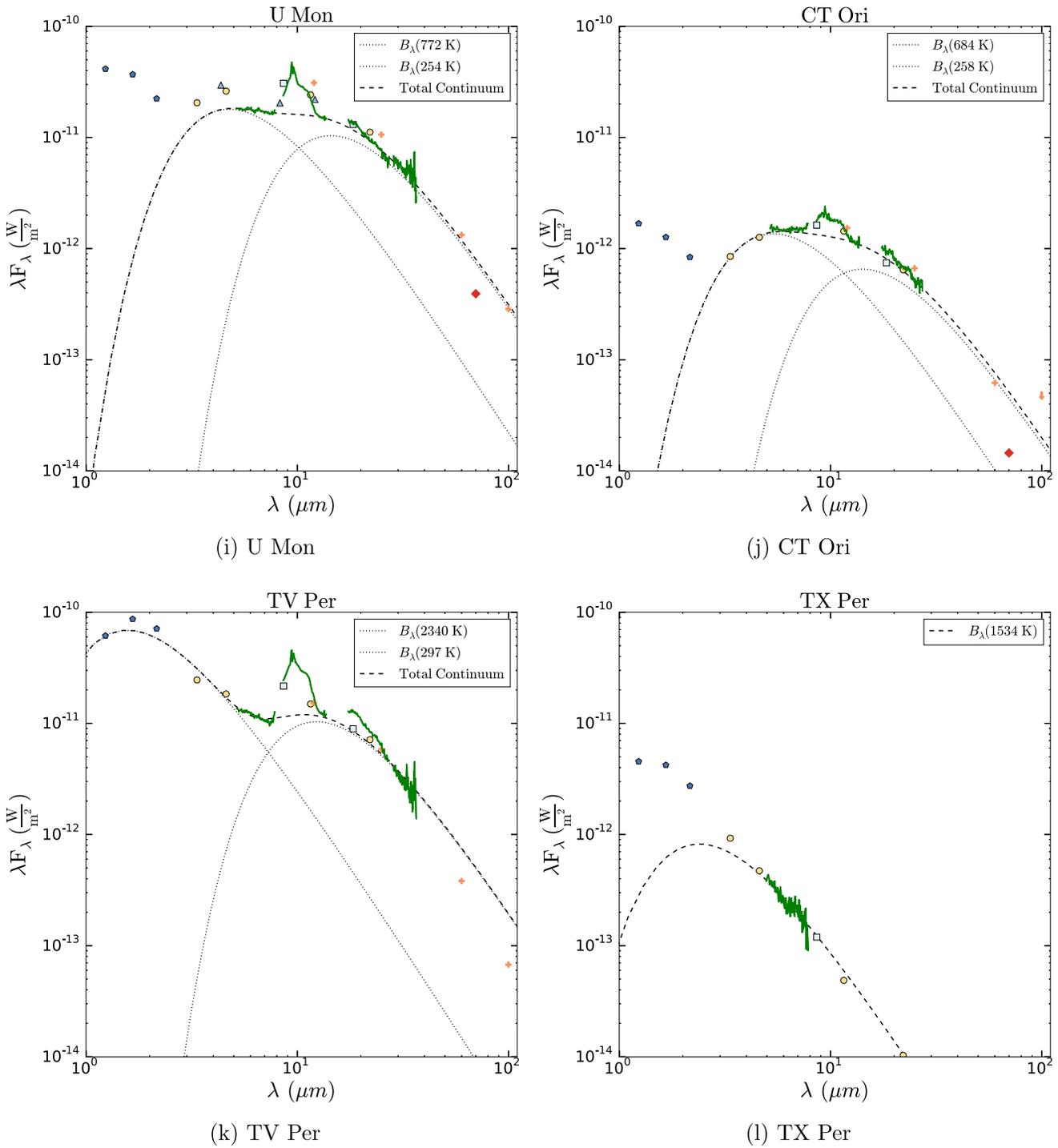


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source, the dust crystallization fraction increases (Gielen et al. 2011) and the disk dissipates (Kastner et al. 2004, 2016; Gezer et al. 2015; Lisse et al. 2017). Thus, the IR excess associated with the dusty disk diminishes as the system ages and transitions to a PN.

In this work we present 5–40  $\mu\text{m}$  IR spectra on a diverse sample of RV Tauri and SRd variables from a grism spectroscopic study of suspected proto-planetary nebula precursors with the Faint Object infraRed CAmera for the

SOFIA Telescope (FORCAST; Herter et al. 2012) instrument on board the NASA Stratospheric Observatory for Infrared Astronomy (SOFIA; Becklin et al. 2007; Gehrz et al. 2009; Young et al. 2012). With this rich data set, we produce spectroscopic sampling of these objects in the mid-IR. By modeling the emission we can determine the source of the IR-excess, identify the dust species present and quantify fundamental dust properties, such as the grain size distribution and dust temperature. These parameters help to determine the

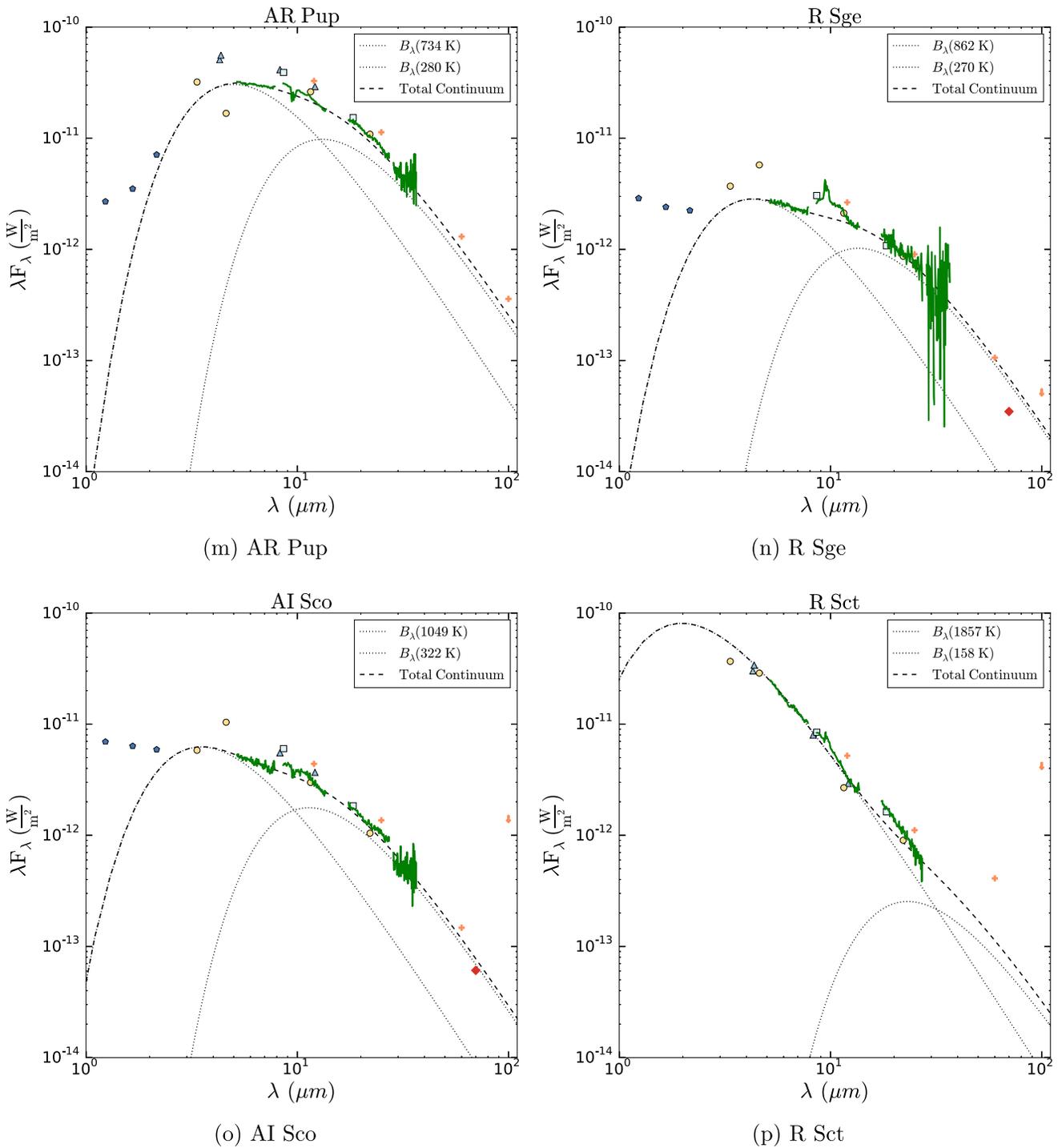


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physical properties of the post-AGB circumstellar environment and to constrain models of post-AGB mass loss and PN formation.

In Section 2 we summarize the stars observed by our program. An overview of the observations and data reduction strategies is given in Section 3. Section 4 contains the construction of the spectral energy distributions (SEDs) and the spectral decomposition model we used to measure the mineralogy of the program stars. The results of our model, a

discussion of our results, and our conclusions are presented in Sections 5–7.

## 2. Program Stars

We have selected a sample of RV Tauri and related SRd stars based upon: (1) their availability for SOFIA flights from Palmdale, CA and Christchurch, NZ, (2) diversity of their IR SEDs, and (3) our ability to obtain a signal-to-noise ratio

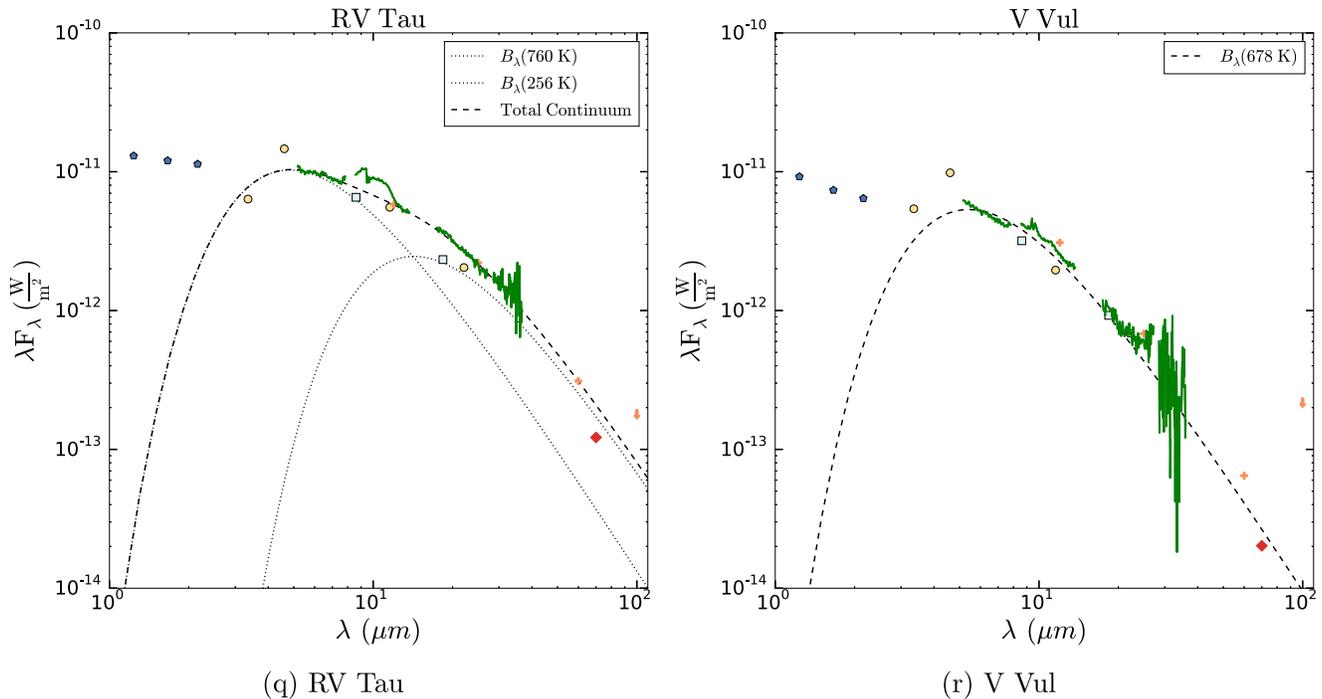


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compatible with our science objectives in a reasonable integration time.

The properties of the 18 RV Tauri and SRd variables presented in this work are summarized in Table 1. TX Per is sometimes categorized as an RV Tauri star and sometimes as an SRd variable. We concur with Percy & Coffey (2005), which refers to TX Per as being a “mild” RV Tauri as the consecutive minima are very similar in depth, and categorize TX Per as an RV Tauri variable.

The mineralogy of many of these systems has been studied previously (Molster et al. 2002c; Deroo et al. 2006; Gielen et al. 2007, 2011; Blommaert et al. 2014; Hillen et al. 2015). These studies have mostly focused on crystalline silicates. Most of the studies found evidence for large, crystalline dust grains indicative of highly processed material. Some of the studies suggest that the crystalline and amorphous silicates are at different temperatures, suggesting that the two species are spatially separated and have different formation histories.

### 3. Observations and Data Reduction

The targets were observed with SOFIA during Guest Investigator (GI) Cycles 2, 3, and 4. Descriptions of the SOFIA Observatory and its science instrument (SI) suite have been given by Becklin et al. (2007), Gehrz et al. (2009), and Young et al. (2012). All of the targets in our survey were observed using FORCAST.

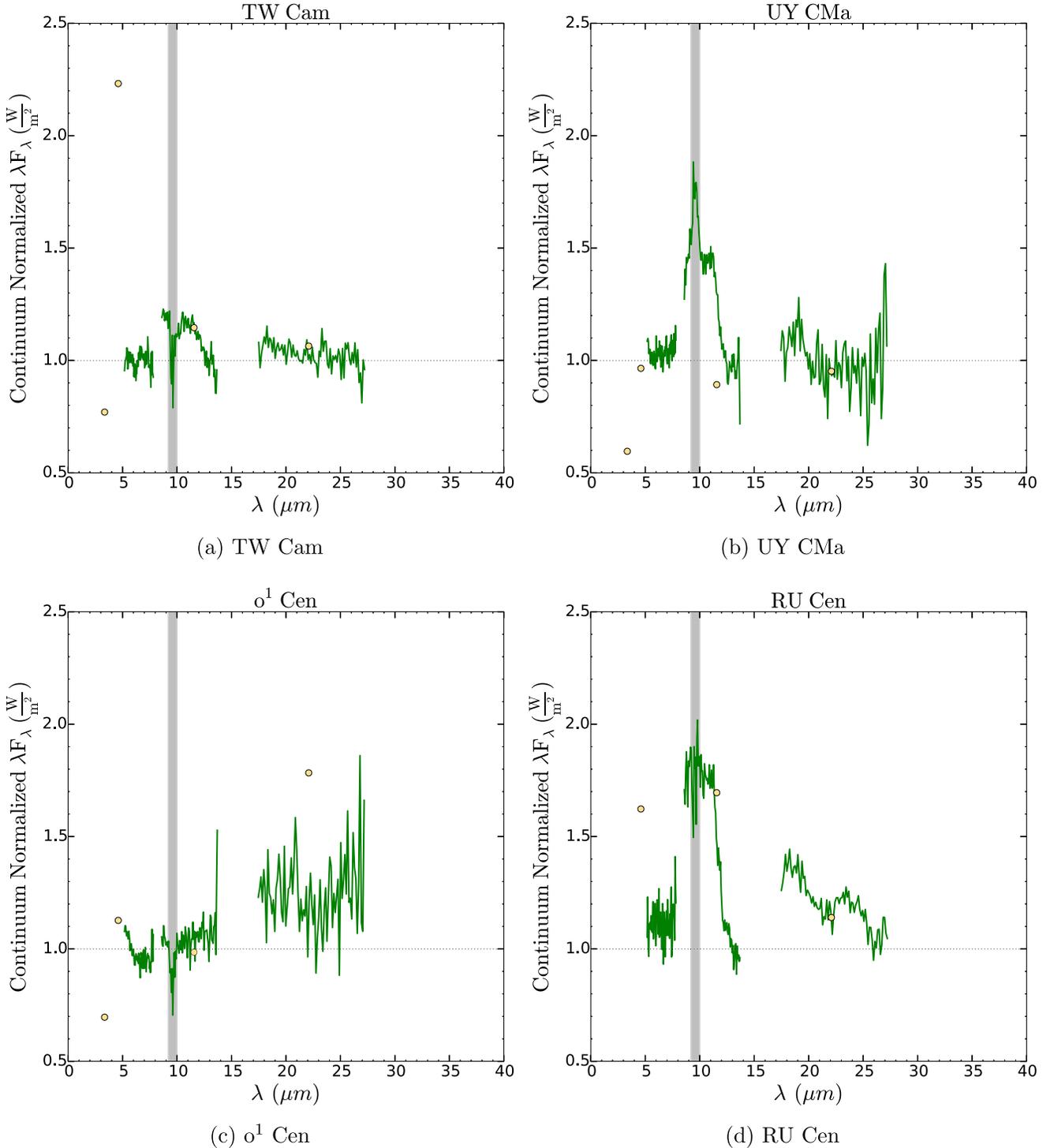
FORCAST is a dual-channel mid-IR camera and spectrograph operating between 5 and 40  $\mu\text{m}$ . Each channel consists of a  $256 \times 256$  pixel array that yields a  $3'.4 \times 3'.2$  field of view with a square plate scale of  $0''.768$ , after distortion correction. The Short Wave Camera (SWC) uses a Si:As blocked-impurity band (BIB) array optimized for  $\lambda < 25 \mu\text{m}$ , while the Long Wave Camera’s (LWC) Si:Sb BIB array is optimized for  $\lambda > 25 \mu\text{m}$ . Observations can be made through either of the two channels individually or, by use of a dichroic

mirror, with both channels simultaneously across the entire range. All of the observations presented in this work were taken in the single-channel, long-slit mode. We utilized FORCAST’s suite of gratings, which provided low spectral resolution ( $R \approx 200$ ) over the 5–40  $\mu\text{m}$  range. The following gratings were used for our observations: G1 covering 4.9–8.0  $\mu\text{m}$ , G3 covering 8.4–13.7  $\mu\text{m}$ , G5 covering 17.6–27.7  $\mu\text{m}$ , and G6 covering 28.7–37.1  $\mu\text{m}$ . All of the observations were taken in the “nod match chop” mode (C2N) which used a chop throw of 30”, a chop angle of either 0° or 30°, and no dithering.

The data were reduced by the SOFIA Science Center using FORCAST Redux v1.5.0 and v1.2.0 pipeline versions (Clarke et al. 2015) and released to the authors as level 3 results. We stacked the spectra when there were multiple observations of a given target. We did not use any of the data points between 9.19–10.0  $\mu\text{m}$  as these are strongly affected by telluric ozone absorption. The spectra were smoothed with a three-point unweighted boxcar to emphasize spectral features.

### 4. SEDs of the Survey Objects

We present the IR SEDs of the survey objects in Figure 1. As can be seen in Figure 2, all of the program stars except for TX Per were observed with the G1, G3, and G5 gratings, and only 7 of the 18 stars were also observed with the G6 grism. TX Per was fainter than expected and we were only able to obtain the G1 grism spectrum (4.9–8.0  $\mu\text{m}$ ). We were unable to model the SED of TX Per with such a limited wavelength range. The continuum-normalized SOFIA grism spectra are shown in Figure 2. We have also gathered archival broadband IR photometry for comparison with the SOFIA spectra. Photometry from the Two Micron All-Sky Survey (2MASS; Skrutskie et al. 2006) at 1.25, 1.65, and 2.17  $\mu\text{m}$ , the AKARI satellite (Murakami et al. 2007) at 9 and 18  $\mu\text{m}$ , the *Wide-field Infrared Survey Explorer* (WISE; Wright et al. 2010) at 3.4, 4.6, 12 and 22  $\mu\text{m}$ , and IRAS (Neugebauer et al. 1984) at 12, 25, 60,



**Figure 2.** Continuum-normalized SOFIA FORCAST spectrum (green curve) after dividing by the best-fitting continuum of our sample of stars showing the *WISE* archival photometry (yellow points) and telluric ozone region (gray band). (An extended version of this figure is available.)

and  $100\ \mu\text{m}$  were all taken from the NASA/IPAC Infrared Science Archive (IRSA; Berriman 2008) database. Photometry from the Herschel (Pilbratt 2003) Photoconductor Array Camera and Spectrometer (PACS; Poglitsch et al. 2010) at  $70$ ,  $100$ , and  $160\ \mu\text{m}$  and Spectral and Photometric Imaging Receiver (SPIRE; Griffin et al. 2010) at  $250$ ,  $350$ , and  $500\ \mu\text{m}$  were obtained using aperture photometry after sky background subtraction. The archival photometry data points were *not* used

in any of the least-squares fitting routines and are only plotted to visualize the SED of each star.

#### 4.1. Dust Species

Previous studies have shown that the most common dust species present in circumstellar environments are amorphous and crystalline silicates with olivine and pyroxene stoichiometries (Molster et al. 2002a, 2002b, 2002c; Gielen et al. 2011).

**Table 2**  
Dust Species and Properties Used in This Work

Dust Species	Composition	Structure	Density ( $\text{g cm}^{-3}$ )	Grain Size ( $\mu\text{m}$ )	References
Forsterite	$\text{Mg}_2\text{SiO}_4$	C	3.27	0.1	Koike et al. (2003)
Olivine	$\text{Mg}_2\text{SiO}_4$	A	3.71	0.1, 2.0	Dorschner et al. (1995)
Pyroxene	$\text{MgSiO}_3$	A	3.20	0.1, 2.0	Dorschner et al. (1995)
Carbon	Pyrolyzed at 400°C	A	1.435	0.1, 2.0	Jaeger et al. (1998b)
Silicon carbide	$\alpha\text{-SiC}$	C	3.26	0.1, 2.0	Pegourie (1988)
Graphite	...	C	2.24	0.1, 2.0	Draine & Lee (1984)
Metallic iron	Fe	C	7.87	0.1, 2.0	Pollack et al. (1994)

**Note.** The mineral structure is denoted as either amorphous (A) or crystalline (C).

Amorphous olivine has very prominent broad features around 9.8 and 18  $\mu\text{m}$ . These features arise from the Si–O stretching and O–Si–O bending modes. Amorphous pyroxene shows a 10  $\mu\text{m}$  feature similar to that of amorphous olivine, but shifted toward shorter wavelengths. Crystalline forsterite has prominent emission features at 11.2, 23.7, and 33.7  $\mu\text{m}$ . Many of the stars in our sample are known to be oxygen rich (see Table 1); however, an increased abundance of carbon dredged up as the star evolves is possible (Iben 1981; Chan & Kwok 1990). Models suggest that stars with main sequence masses less than 1.5  $M_\odot$  do not experience third dredge-up. Therefore, the surface composition of these stars is fixed by the first dredge-up and red giant branch extra-mixing, and they are oxygen-rich AGB stars. For main sequence stars with masses over 6  $M_\odot$  the stars undergo shallow third dredge-up episodes. Stars in the intermediate main sequence mass range undergo repeated third dredge-up episodes that bring carbon to the surface to become carbon stars (for a review see Straniero et al. 2006; Karakas & Lattanzio 2014). The interpulse time for the third dredge-up is a few times  $10^4$  years with carbon stars undergoing several pulses after the carbon abundance first exceeds the oxygen abundance (Straniero et al. 1997). The interpulse time exceeds the expansion time for circumstellar shells and, as expected, nearly all normal AGB stars have shells of the same composition as the star. Carbon-rich AGB stars are expected to have circumstellar shells dominated by amorphous carbon or graphite grains with some silicon carbide (SiC) possibly present (Suh 2000; Speck et al. 2005, 2009). Amorphous carbon does not have prominent IR features but contributes to the dust continuum emission; however, graphitic carbon and silicon carbide have emission features at 11.53  $\mu\text{m}$  and in the 10–13  $\mu\text{m}$  region, respectively. The strong depletion of iron in the photospheres of RV Tauri stars suggests that metallic iron may be present in the circumstellar environment. Iron grains can form at temperatures 50–100 K lower than silicates and are stable in O-rich environments above 700 K (Kemper et al. 2002). Like amorphous carbon, metallic iron lacks prominent IR features and contributes to the overall dust continuum. Therefore, the dust species we included in our model are crystalline forsterite, amorphous olivine, amorphous pyroxene, amorphous carbon, silicon carbide, graphite, and metallic iron. We only considered the magnesium-rich crystalline species of olivine (forsterite) as it has been found that the iron content of circumstellar crystalline olivine around evolved stars is lower than 10% (Tielens et al. 1998; Molster et al. 2002c). In addition, we only included amorphous enstatite and pyroxene with an iron content of 0% as these

species were found to be better fits to the spectra than the same species with an iron content of 50%. We tried including Mg-rich crystalline olivine, Mg-rich crystalline enstatite, crystalline bronzite, crystalline fayalite, iron oxide, amorphous alumina, and amorphous silica in our model; however none of these dust species were significantly present. We discuss the addition of some of these species further in Section 6.4.

Mass absorption coefficients for the different dust species are calculated from optical constants using a homogeneous sphere approximation (Min et al. 2005). Although the continuous distribution of ellipsoids (CDE) approximation (Bohren & Huffman 1983) is widely used, it is only valid in the Rayleigh limit for small grain sizes. Because we are interested in the grain size distribution we did not use the CDE approximation. The details of the different optical constants that we used can be found in Table 2. In cases where the refractive index was reported for the three crystallographic directions, we assumed even distributions of each orientation. There are many laboratory measurements of IR optical constants available, corresponding to different material compositions, crystal structures, annealing temperatures, measurement environments, grain sizes, and grain orientations. These different measurements produce spectra with similar global features but with unique differences. While the minerals we have chosen to use in our model may result in different relative abundances compared to those derived in previous studies, we are less concerned with making comparisons with those works and more concerned with drawing comparisons between the program stars in the present study.

#### 4.2. Spectral Decomposition Model

To identify the minerals present and to quantify the grain size distributions, we fit the observed SOFIA spectra with synthetic spectra of various mineral species. The synthetic spectra were calculated from the optical constants of each mineral. The conversion from laboratory measured optical constants of dust to mass absorption coefficients is not straightforward. Several factors affect the observed emission features, including the chemical composition of the dust, the grain size, and the grain shape (Min et al. 2003, 2005). We constructed a basic model to fit the full FORCAST wavelength range. Given that we see the silicate features in emission, we assume that the dust features are in an optically thin part of the disk and, therefore, we approximate the spectrum as a linear combination of dust absorption profiles. The emission model is

**Table 3**  
RV Tauri and SRd Star Mineralogy

Star	$\chi^2_{\text{red}}$	Mineral	$\bar{c}_i$	$\sigma_{\bar{c}_i}$	$V_f$	$\sigma_{V_f}$	$T_1$ (K)	$F_{T_1}$	$T_2$ (K)	$F_{T_2}$
TW Cam	0.26	Graphite-small	3.75	0.15	0.58	2.6E-2	1365 ± 105	0.019 ± 3.6E-3	332 ± 8	0.981 ± 3.6E-3
		Graphite-large	1.14	0.06	0.18	1.1E-2				
		Carbon-large	1.08	0.10	0.17	1.7E-2				
		Pyroxene-small	0.26	0.03	0.04	4.1E-3				
		SiC-small	0.20	0.04	0.03	5.7E-3				
UY CMa	0.08	Graphite-small	1.76	0.49	0.36	1.0E-1	1251 ± 201	0.010 ± 4.1E-3	337 ± 7	0.990 ± 4.1E-3
		Graphite-large	1.29	0.11	0.26	3.6E-2				
		Carbon-large	1.20	0.21	0.24	4.9E-2				
		Pyroxene-small	0.47	0.06	0.10	1.7E-2				
		Forsterite-small	0.21	0.04	0.04	8.6E-3				
o <sup>1</sup> Cen	0.32	Iron-small	7.01	0.48	0.49	3.8E-2	3780 ± 193	1.0	...	...
		Graphite-small	5.96	0.49	0.42	3.7E-2				
		Graphite-large	1.17	0.06	0.08	5.5E-3				
		Pyroxene-small	0.17	0.03	0.01	2.1E-3				
RU Cen	0.23	Graphite-small	3.53	0.20	0.45	2.8E-2	535 ± 16	0.010 ± 1.4E-3	203 ± 2	0.990 ± 1.4E-3
		Carbon-large	2.01	0.18	0.25	2.4E-2				
		Pyroxene-small	0.89	0.05	0.11	8.0E-3				
		Graphite-large	0.86	0.08	0.11	1.1E-2				
		SiC-small	0.50	0.06	0.06	8.3E-3				
		Forsterite-small	0.11	0.02	0.01	3.0E-3				
SX Cen	0.24	Graphite-small	2.65	0.26	0.41	4.5E-2	715 ± 14	0.042 ± 4.8E-3	244 ± 7	0.958 ± 4.8E-3
		Carbon-large	1.77	0.14	0.27	2.7E-2				
		Graphite-large	0.95	0.08	0.15	1.5E-2				
		Pyroxene-large	0.72	0.19	0.11	3.0E-2				
		Olivine-large	0.32	0.18	0.05	2.8E-2				
		Pyroxene-small	0.08	0.06	0.01	9.8E-3				
SU Gem	0.40	Graphite-small	4.06	0.12	0.59	1.9E-2	1444 ± 83	0.023 ± 3.1E-3	349 ± 5	0.977 ± 3.1E-3
		Carbon-large	1.16	0.07	0.17	1.1E-2				
		Graphite-large	1.07	0.05	0.15	7.5E-3				
		Pyroxene-small	0.38	0.02	0.06	2.9E-3				
		SiC-small	0.27	0.03	0.04	3.9E-3				
AC Her	1.43	Graphite-small	3.43	0.06	0.48	9.4E-3	489 ± 4	0.022 ± 9.4E-4	207 ± 1	0.978 ± 9.4E-4
		Carbon-large	1.51	0.05	0.21	6.6E-3				
		Graphite-large	0.92	0.02	0.13	3.6E-3				
		Pyroxene-small	0.76	0.01	0.11	2.3E-3				
		SiC-small	0.50	0.02	0.07	2.6E-3				
		Forsterite-small	0.09	0.01	0.01	9.8E-4				
V441 Her	0.15	Graphite-small	3.14	0.14	0.40	2.2E-2	1607 ± 68	0.019 ± 1.8E-3	363 ± 3	0.981 ± 1.8E-3
		Carbon-small	2.76	0.20	0.35	2.6E-2				
		Graphite-large	1.07	0.05	0.14	8.3E-3				
		Olivine-small	0.44	0.08	0.06	1.0E-2				
		Pyroxene-small	0.23	0.04	0.03	5.2E-3				
		SiC-small	0.21	0.04	0.03	5.3E-3				
U Mon	0.50	Carbon-small	3.55	0.50	0.30	6.3E-2	772 ± 10	0.020 ± 1.0E-3	254 ± 1	0.980 ± 1.0E-3
		Graphite-small	3.17	0.13	0.27	4.5E-2				
		Iron-large	2.94	1.88	0.25	1.6E-1				
		Pyroxene-small	0.87	0.06	0.07	1.3E-2				
		SiC-small	0.56	0.05	0.05	9.0E-3				
		Graphite-large	0.42	0.15	0.04	1.4E-2				
		Olivine-small	0.38	0.10	0.03	9.8E-3				
CT Ori	0.25	Graphite-small	3.91	0.17	0.53	2.5E-2	684 ± 13	0.041 ± 4.0E-3	258 ± 5	0.959 ± 4.0E-3
		Carbon-large	1.58	0.12	0.21	1.7E-2				
		Graphite-large	0.97	0.07	0.13	1.1E-2				
		Pyroxene-small	0.57	0.03	0.08	5.0E-3				
		SiC-small	0.38	0.04	0.05	5.7E-3				
TV Per	0.58	Carbon-small	3.75	0.58	0.29	6.5E-2	2340 ± 12	0.002 ± 4.0E-5	297 ± 1	0.998 ± 4.0E-5
		Iron-large	3.24	2.20	0.25	1.7E-1				

**Table 3**  
(Continued)

Star	$\chi_{\text{red}}^2$	Mineral	$\bar{c}_i$	$\sigma_{\bar{c}_i}$	$V_j$	$\sigma_{V_j}$	$T_1$ (K)	$F_{T_1}$	$T_2$ (K)	$F_{T_2}$
		Graphite-small	2.47	0.16	0.19	3.5E-2				
		Pyroxene-small	1.47	0.07	0.11	2.0E-2				
		Olivine-small	1.04	0.14	0.08	1.7E-2				
		SiC-small	0.87	0.07	0.07	1.3E-2				
		Graphite-large	0.30	0.17	0.02	1.4E-2				
AR Pup	1.93	Graphite-small	4.07	0.11	0.48	1.4E-2	734 ± 7	0.062 ± 3.0E-3	280 ± 3	0.938 ± 3.0E-3
		Carbon-small	2.88	0.06	0.34	8.3E-3				
		Graphite-large	1.02	0.02	0.12	2.8E-3				
		Olivine-large	0.23	0.01	0.03	1.5E-3				
		Pyroxene-small	0.19	0.01	0.02	6.9E-4				
R Sge	0.17	Graphite-small	3.66	0.54	0.43	6.8E-2	862 ± 29	0.026 ± 3.6E-3	270 ± 7	0.974 ± 3.6E-3
		Carbon-small	3.02	0.42	0.36	5.5E-2				
		Graphite-large	0.97	0.13	0.11	1.8E-2				
		Pyroxene-small	0.51	0.03	0.06	6.2E-3				
		SiC-small	0.34	0.05	0.04	6.8E-3				
AI Sco	0.32	Graphite-small	3.94	0.17	0.59	3.0E-2	1049 ± 74	0.030 ± 6.1E-3	322 ± 7	0.970 ± 6.1E-3
		Carbon-large	1.15	0.14	0.17	2.1E-2				
		Graphite-large	1.10	0.08	0.16	1.3E-2				
		Pyroxene-small	0.30	0.03	0.04	4.2E-3				
		SiC-small	0.19	0.04	0.03	6.3E-3				
R Sct	1.18	Graphite-small	5.42	0.23	0.70	3.1E-2	1857 ± 44	0.017 ± 7.0E-3	158 ± 18	0.983 ± 7.0E-3
		Graphite-large	1.10	0.05	0.14	8.0E-3				
		Carbon-large	1.04	0.07	0.13	1.0E-2				
		Pyroxene-small	0.23	0.02	0.03	2.2E-3				
RV Tau	0.26	Graphite-small	3.70	0.13	0.43	1.9E-2	760 ± 10	0.052 ± 2.9E-3	256 ± 3	0.948 ± 2.9E-3
		Carbon-small	3.21	0.21	0.37	2.5E-2				
		Graphite-large	0.92	0.06	0.11	7.1E-3				
		Pyroxene-small	0.47	0.03	0.05	3.3E-3				
		SiC-small	0.38	0.03	0.04	3.9E-3				
V Vul	0.36	Iron-small	8.06	0.55	0.55	4.1E-2	678 ± 4	1.0	...	...
		Graphite-small	4.92	0.41	0.34	3.1E-2				
		Graphite-large	1.05	0.05	0.07	5.0E-3				
		Pyroxene-large	0.31	0.07	0.02	4.9E-3				
		SiC-large	0.16	0.06	0.01	4.4E-3				
		Pyroxene-small	0.10	0.05	0.01	3.3E-3				

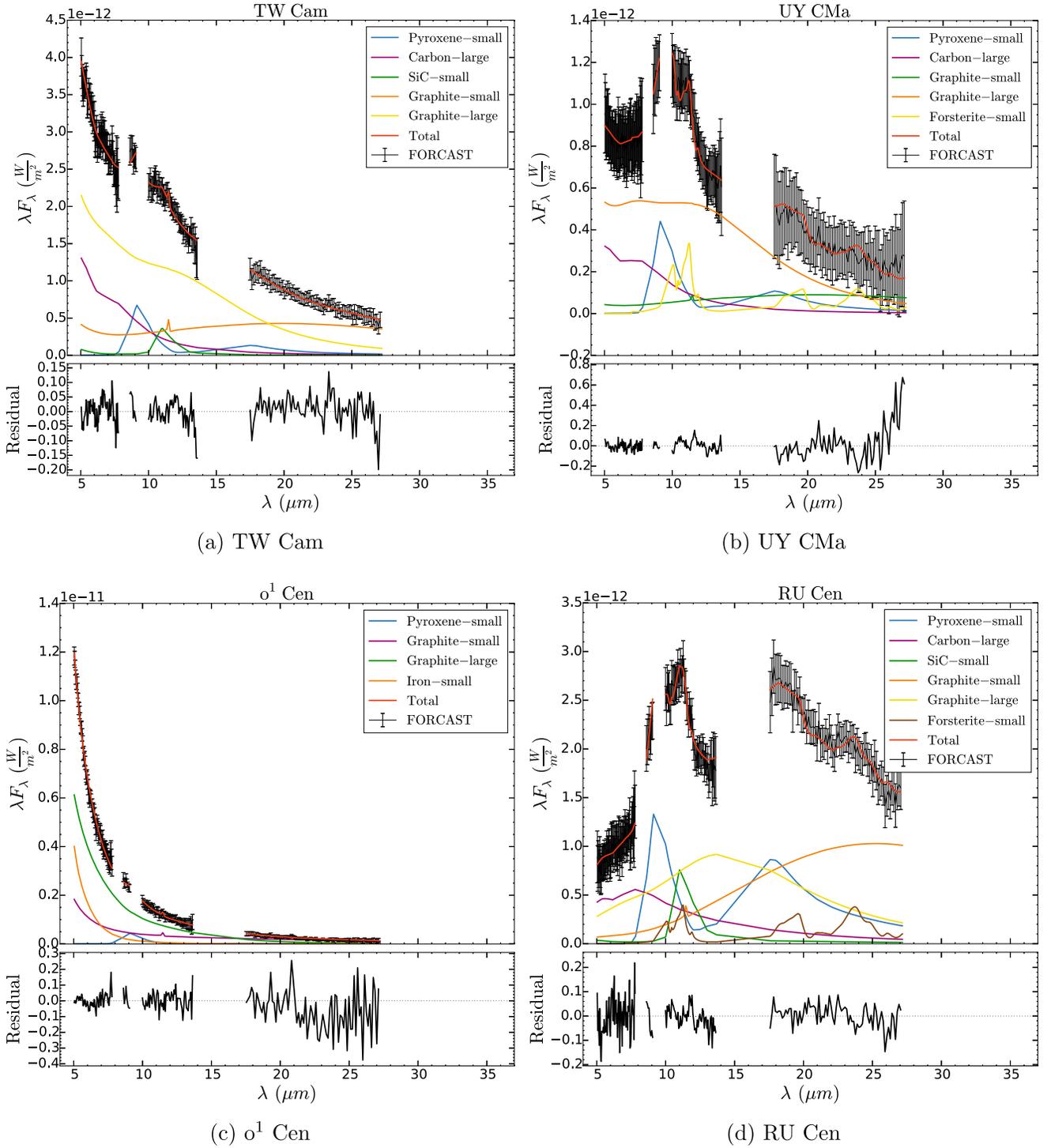
**Note.**  $\bar{c}_i$  is the average best-fit coefficient of a given mineral species.  $\sigma_{\bar{c}_i}$  is the standard deviation of the average best-fit coefficient,  $V_j$  is the volume fraction,  $\sigma(V_j)$  is the error in the volume fraction,  $T_{1,2}$  are the blackbody dust temperatures, and  $F_{T_{1,2}}$  is the fraction of dust at those temperatures, respectively; “small” refers to 0.1  $\mu\text{m}$  spherical grains, and “large” designates 2.0  $\mu\text{m}$  spherical grains.

given by

$$\lambda F_{\lambda} \propto \sum_i c_i \mu_i(\lambda) \times \sum_j a_j \lambda B_{\lambda}(T_j), \quad (1)$$

where  $\mu_i(\lambda)(\text{cm}^{-1})$  is the absorption coefficient of dust component  $i$  and  $c_i$  gives the volume fraction of that dust component,  $B_{\lambda}(T_j)$  ( $\text{W sr}^{-1} \text{m}^{-3}$ ) denotes the Planck function at temperature  $T_j$  and  $a_j$  is the scaling factor for the  $j$ th Planck function. The absorption coefficient is related to the mass absorption coefficient (opacity) by  $\mu = \kappa\rho$ , where  $\kappa(\text{cm}^2 \text{g}^{-1})$  is the mass absorption coefficient (opacity) and  $\rho(\text{g cm}^{-3})$  is the density of the dust component. The formulation of our model assumes that the stellar contribution to the SED in this range is negligible. We further assume that all of the dust in a population is in thermal equilibrium with all of the other dust species, regardless of particle size or the ability to absorb and

re-emit starlight. We first used a least-squares minimization to fit the Planck functions to the FORCAST continuum. We fit the functions to the entire FORCAST wavelength range available except for the 8–12.5  $\mu\text{m}$  range, which is dominated by silicate emission. Two Planck functions were used for all of the spectral models except for o<sup>1</sup> Cen and V Vul where only a single Planck function was needed to fit the underlying continuum. The best-fitting Planck function parameters are shown in Figure 1 and summarized in Table 3; however, the sum of the Planck scaling factors,  $a_j$ , has been normalized to unity in order to represent the fraction of dust at each temperature  $T_j$ . The best-fitting Planck functions are plotted with the FORCAST spectra in Figure 1. It is probably not realistic to model the dust as a single-temperature component or even as a two-temperature component—a temperature gradient is probably more realistic—but in order to keep the



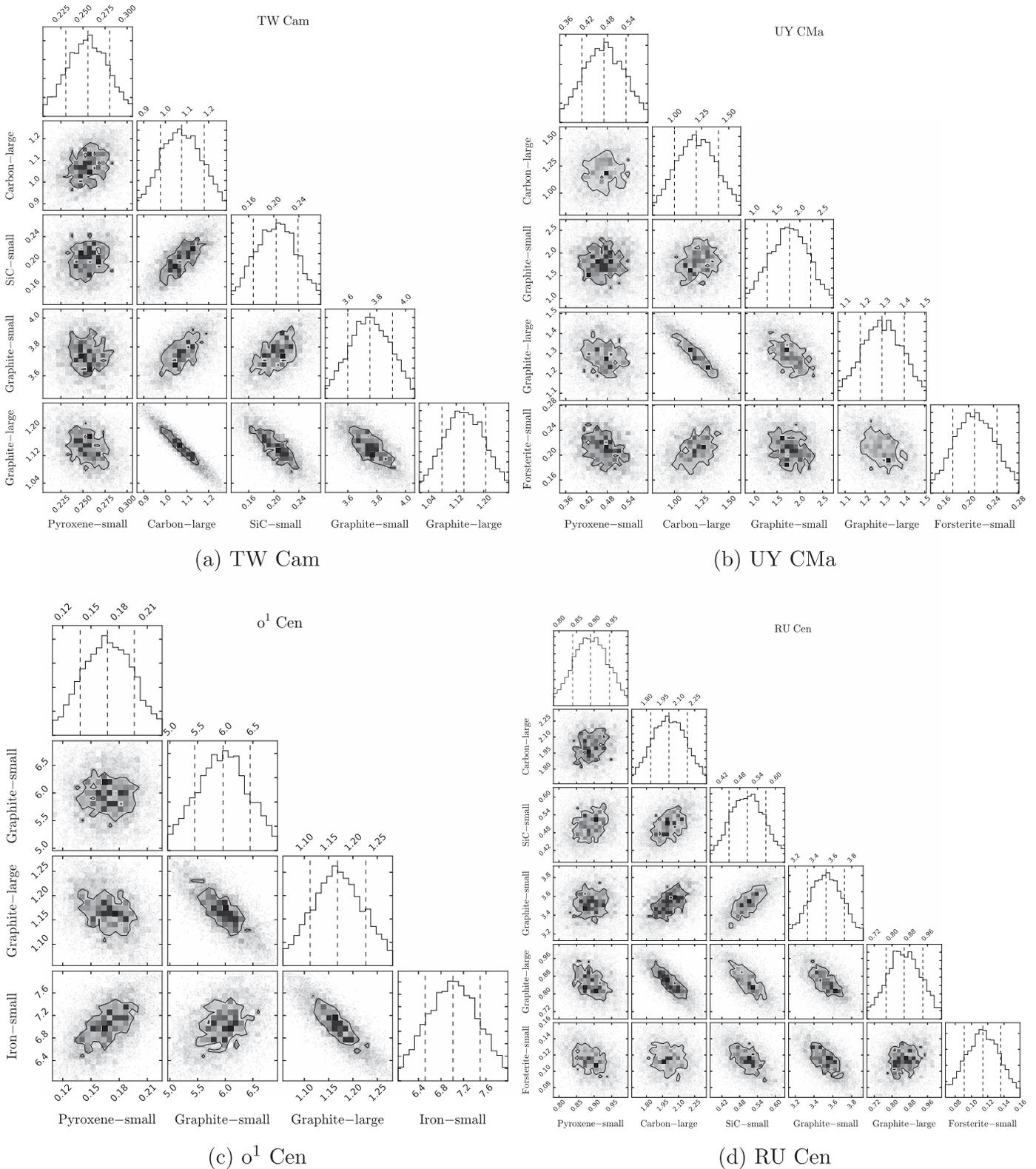
**Figure 3.** Best model fits for our sample of stars, showing the contribution of the different mineral species. Top: the observed SOFIA FORCAST spectra and  $1\sigma$  errors (black points) are plotted together with the best model fit (red curve) and the mineral species (colored curves). The data points between  $9.19\text{--}10.0\ \mu\text{m}$  have been removed as these are strongly affected by telluric ozone absorption. Bottom: the normalized residual spectra after dividing by the best model of the observed spectra. (An extended version of this figure is available.)

number of parameters at a minimum we only use a maximum of two Planck functions in our model.

After finding the best-fitting Planck functions, the best-fitting dust fraction coefficients,  $c_i$ , were found by a non-negative least-squares minimization of the model to the entire observed

FORCAST wavelength range. The reduced  $\chi^2$  of the spectrum is given by

$$\chi_{\text{red}}^2 = \frac{1}{N - M} \sum_{i=1}^N \left| \frac{F_{\text{model}}(\lambda_i) - F_{\text{obs}}(\lambda_i)}{\sigma_i} \right|^2, \quad (2)$$



**Figure 4.** Normalized probability distribution functions of the best-fit coefficients,  $c_i$ , of each mineral species after 5000 realizations of a Monte Carlo simulation with Gaussian noise distributions and the covariance between the coefficients. The dashed lines show the mean ( $\bar{c}_i$ ) and the  $1\sigma$  ( $\sigma_{\bar{c}_i}$ ) confidence levels. The contours show the  $1\sigma$  confidence levels. (An extended version of this figure is available.)

where  $N$  is the number of wavelength points,  $M$  is the number of fit parameters,  $F_{\text{model}}(\lambda_i)$  is the model flux at a given wavelength,  $F_{\text{obs}}(\lambda_i)$  is the observed flux at a given wavelength,

and  $\sigma_i$  is the absolute error of the observed flux at each wavelength  $\lambda_i$ . The reduced  $\chi^2$  values are summarized in Table 3.

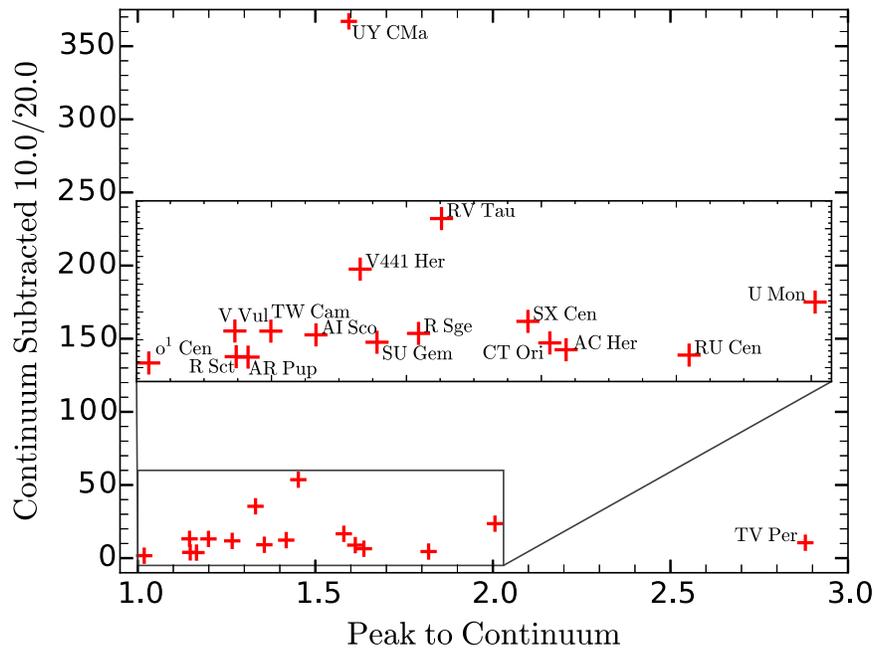


Figure 5. Ratio of the continuum subtracted flux at 10 and 20  $\mu\text{m}$  vs. the peak-to-continuum ratio of the 10  $\mu\text{m}$  silicate feature.

Errors on the dust fraction coefficients were calculated from 5000 realizations of a Monte Carlo simulation with Gaussian noise distributions. We omitted a mineral species from the fit if the error,  $\sigma_{\bar{c}_i}$ , on  $\bar{c}_i$  was greater than the value of  $\bar{c}_i$  itself. Figure 4 illustrates the distributions of the best-fitting dust fraction coefficients,  $c_i$ , and the covariance between the coefficients (Foreman-Mackey 2016). Most of the coefficients show little to no correlation with the exception of the carbon, graphite, and metallic iron species which are strongly anti-correlated in most cases. Even though this model is only an approximation, it gives a good fit overall to the observed spectra (see Figure 3). A full radiative transfer model that could account for a temperature gradient in the disk that may produce a better fitting result is beyond the scope of this work.

#### 4.3. Grain Size Distribution

To study the grain size distribution, we used two dust grain sizes in our model with radii of 0.1 and 2.0  $\mu\text{m}$ . These sizes were chosen based on the work of Bouwman et al. (2001) and Honda et al. (2004) which found that, in the 10  $\mu\text{m}$  spectral region, 0.1  $\mu\text{m}$  grains sufficiently describe grains with  $a < 1.0 \mu\text{m}$  while 1.5–2.0  $\mu\text{m}$  grains sufficiently describe grains with  $a > 1.5 \mu\text{m}$ . Larger-sized grains were not considered as the emission features from larger grains become too weak to distinguish from the continuum emission. Because the grains with radii of 2.0  $\mu\text{m}$  are in the Mie scattering regime, we used the python module *pymiecoated* (Leinonen 2012) to calculate the mass absorption coefficients for the 2.0  $\mu\text{m}$  grains. *pymiecoated* computes the scattering properties of single- and dual-layered spheres in the Mie regime using the results of Bohren & Huffman (1983) and the optical constants of bulk materials. The mass absorption coefficient for crystalline forsterite was taken directly from Koike et al. (2003). Because we did not have the optical constants for this mineral we were unable to calculate the mass absorption coefficient for larger-sized grains in the Mie scattering regime.

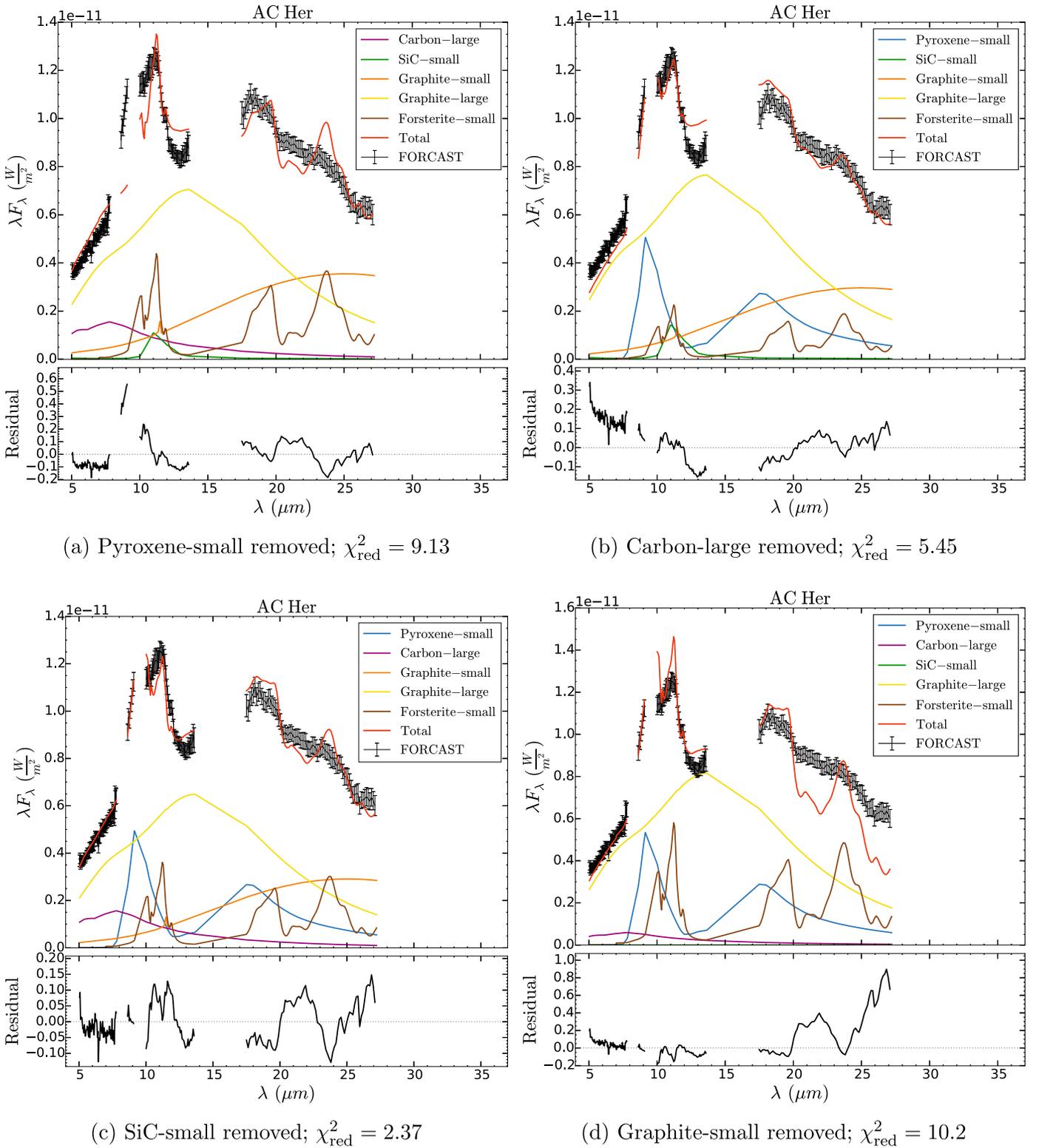
## 5. Results

The results of the spectral decomposition modeling are shown in Figure 3. Details of individual sources are provided in Table 3. The spectroscopic features present in these sources span a broad range of dust properties and characteristics. All of the RV Tauri stars in our sample, with the exceptions of UY CMa, TX Per, and R Sct, have been reported as disk sources (Gezer et al. 2015). Two sources, o<sup>1</sup> Cen and V Vul, display a simple blackbody continuum. However, model fits demonstrate that they both also exhibit a weak IR excess (see Figures 1(c) and (r)), suggesting that these systems may be in the final stages of dissipation. The remaining RV Tauri stars all exhibit emission from carbon-rich minerals with varying degrees of amorphous and crystalline silicates.

Two of our SRd sources, V441 Her and TV Per, show prominent silicate features. The presence of these features suggests that these sources, too, may have dusty disks akin to those of the RV Tauri stars. Alternatively, these features might also arise from normal dusty outflows. In the case of V441 Her, the 10  $\mu\text{m}$  feature is strong while the 20  $\mu\text{m}$  feature is weak, as is expected for relatively fresh and unprocessed amorphous silicates (Nuth & Hecht 1990). In contrast, TV Per exhibits strong 10 and 20  $\mu\text{m}$  silicate emission, suggestive of prolonged exposure to hard radiation. We discuss this more fully in Section 6.

## 6. Discussion

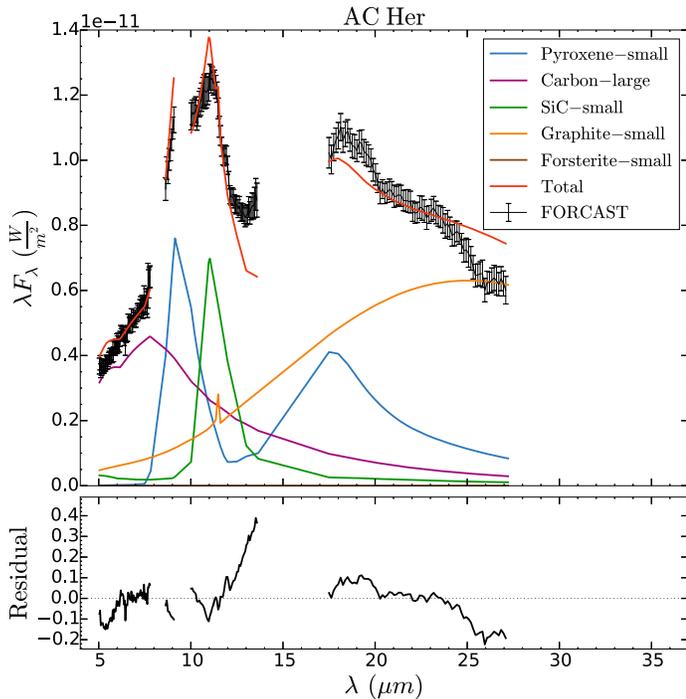
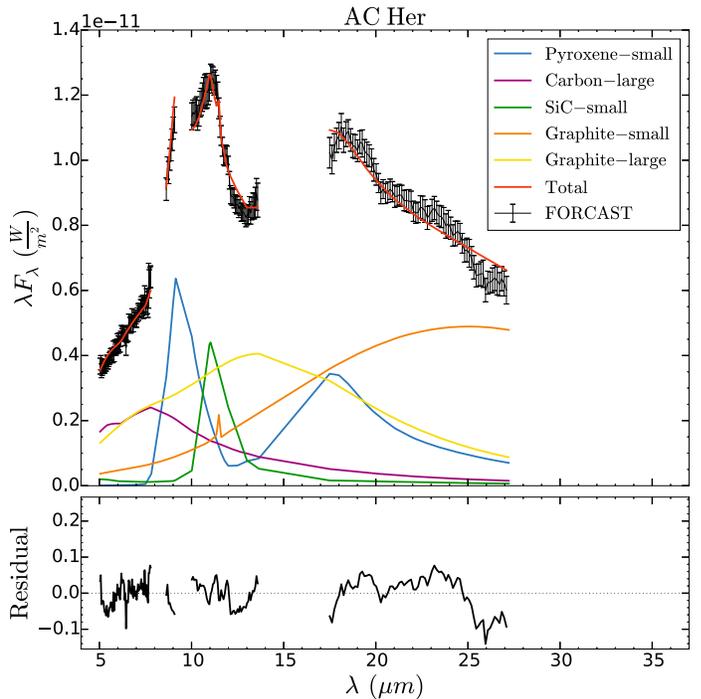
The 10 and 20  $\mu\text{m}$  emission features can be used to quantify the grain size and age of the circumstellar dust (van Boekel et al. 2003, 2005; Juhász et al. 2010). The peak-to-continuum ratio of the 10  $\mu\text{m}$  feature can be used as a measure of the amount of grain growth because larger grains will produce a less prominent feature. In addition, the continuum-subtracted 10/20  $\mu\text{m}$  flux ratio has been shown to decrease monotonically with increased processing and therefore can be used to indicate the age of the circumstellar silicates (Nuth & Hecht 1990).



**Figure 6.** Effect on the fit when removing dust species from the best model of AC Her.

Older, more processed grains will have a lower 10/20  $\mu\text{m}$  ratio. A plot of these two ratios for all our program stars can be seen in Figure 5. Most of the sources show a low peak-to-continuum value (i.e.,  $<2.0$ ) and a low 10/20  $\mu\text{m}$  ratio (i.e.,  $<50$ ) indicating that the grains are relatively large and have undergone significant processing. This supports the idea that

the dust is constrained to a Keplerian disk. There are two outliers in Figure 5, TV Per and UY CMA. TV Per has a high peak-to-continuum value and a small 10/20  $\mu\text{m}$  ratio, indicating that the dust grains are small and old. The small grain size is consistent with our model, which predicts a small spherical grain volume fraction of  $\sim 74\%$  for TV Per. This suggests that

(e) Graphite-large removed;  $\chi^2_{\text{red}} = 5.46$ (f) Forsterite removed;  $\chi^2_{\text{red}} = 1.01$ **Figure 6.** (Continued.)

the circumstellar environment around TV Per is such that the grains are unable to grow to large sizes. UY CMa has a low peak-to-continuum value and a large  $10/20 \mu\text{m}$  ratio, indicating that the dust grains are both large and young. This is also consistent with our model which predicts a small grain volume fraction of  $\sim 50\%$  around UY CMa. Robinson & Hyland (1977) and Mitchell & Robinson (1981) found that for low optical depths, some of the circumstellar silicate dust may be in absorption rather than emission if it is at a low temperature. The viewing angle of the disks will also affect the peak-to-continuum ratio of the  $10 \mu\text{m}$  feature (Crapsi et al. 2008). As mentioned in Section 4.2, factors like the grain size and shape will affect the observed emission features. The optical depth, viewing angle, temperature, and particle size of the grains may mask the  $10/20 \mu\text{m}$  ratio and could contribute to the high volume fraction of small grains found by our model.

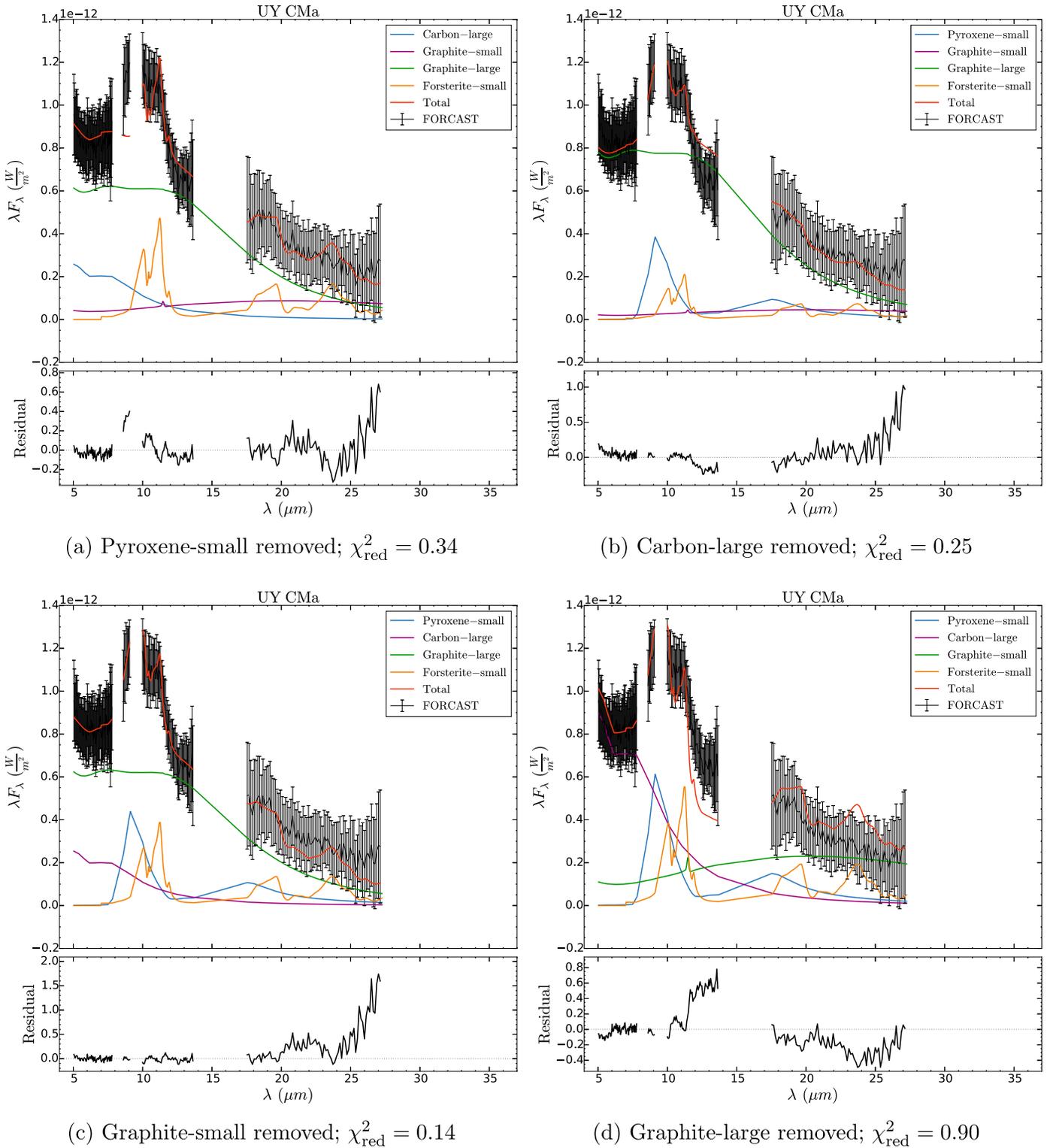
### 6.1. Crystallinity

Although Molster et al. (2002a, 2002b) showed that the silicate crystallinity fraction in disk sources was much higher than that observed in outflow sources, we do not see a high silicate crystallinity fraction in any of the objects in our sample listed as “disk” SEDs in Table 1 and for all of the FORCAST spectra there are no obvious crystalline emission features present. One explanation is that crystalline silicate material is not abundant in any of the stars observed. However, crystalline olivines have been detected around AC Her and are thought to be present at the  $\sim 10\text{--}50\%$  level around AR Pup and U Mon (de Ruyter et al. 2005; Blommaert et al. 2014). The lack of strong crystalline silicate emission features does not necessarily indicate a lack of presence. If a temperature difference exists between the amorphous and crystalline silicates, it is possible

to include up to 40% of crystalline silicates in the circumstellar dust without seeing crystalline features in the spectra (Kemper et al. 2001). Another possible explanation for this observation is that crystalline silicates are generally colder than amorphous silicates, which could mean that the grains are not co-spatial or that they have different optical properties. In fact, when modeling the circumstellar material around AC Her, Hillen et al. (2015) found the spatial distribution of the forsterite to be different from the amorphous dust. The difference in optical properties could be due in part to the different iron content of each material, which increases the opacity in the near-IR significantly (Dorschner et al. 1995; Molster et al. 2002c). Additionally, the spectral features will be less prominent if the crystalline grains are larger than the amorphous grains. Therefore, if the crystalline silicates are only moderately abundant (i.e.,  $\lesssim 40\%$ ), cooler, and larger than the amorphous silicates, the spectral features of the crystalline silicates could easily be masked by the amorphous silicates. Interestingly, Blommaert et al. (2014) did not detect crystalline olivines around AR Pup and U Mon in the mid-IR and interpreted this as an indication that the crystalline olivines are hot ( $\sim 600 \text{ K}$ ) around these two stars. If they are indeed hot, then the crystalline olivine abundances must be relatively low or the grains must be large around these two stars for them to go undetected by FORCAST.

### 6.2. Dual Chemistry

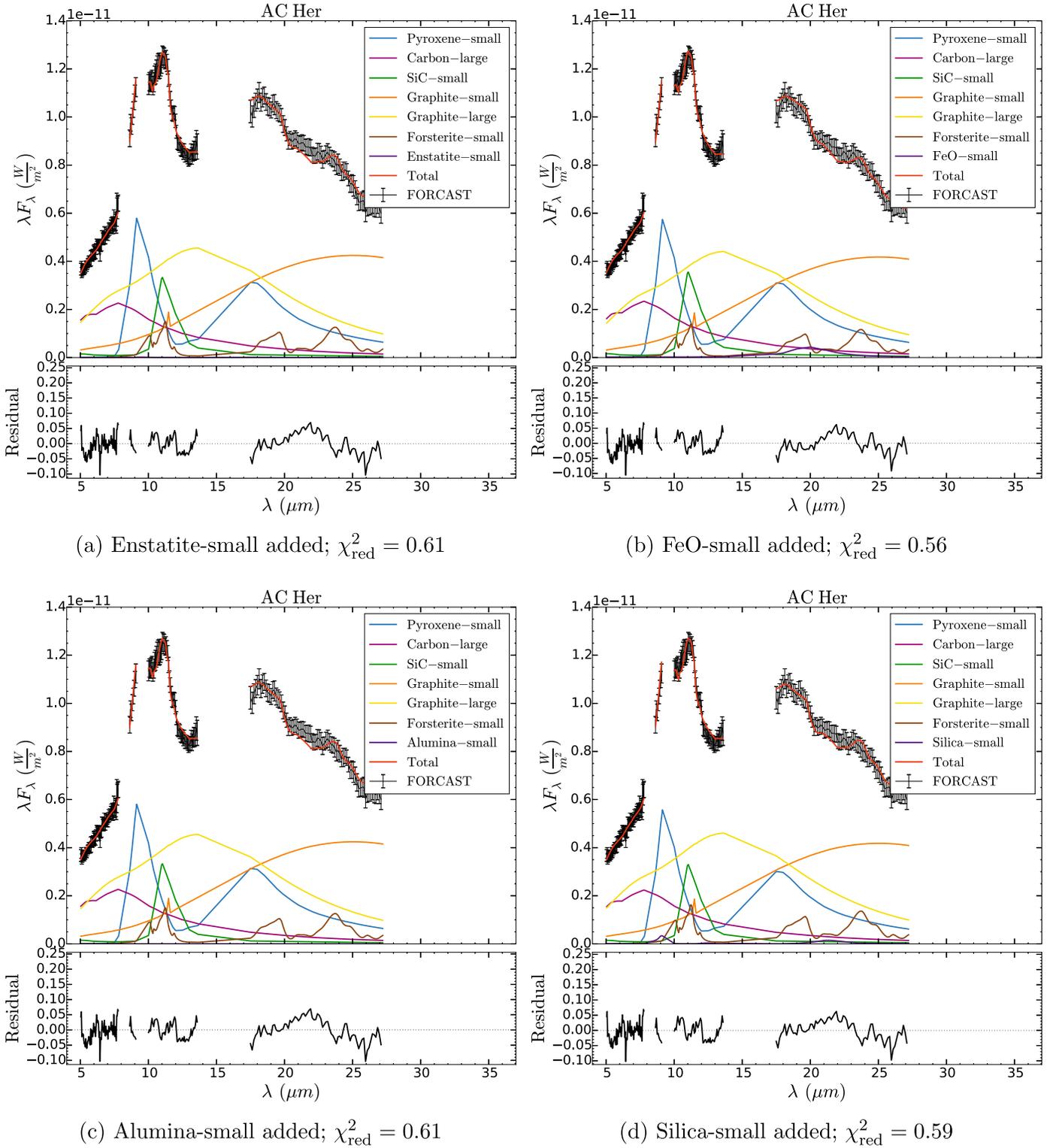
Our model predicts that most of the dust is carbon-rich with some oxygen-rich silicates. This dual formation of carbon- and oxygen-rich minerals has been observed in several classical novae, namely V1280 Sco (Sakon et al. 2016), V705 Cas (Evans et al. 2005), V842 Cen (Smith et al. 1994), and QV Vul (Gehrz et al. 1992) as well as IRAS 09425-6040, a carbon



**Figure 7.** Effect on the fit when removing dust species from the best model of UY CMa. (An extended version of this figure is available.)

AGB star which shows circumstellar silicate dust features (Suh 2016). The formation of both carbon- and oxygen-rich dust could be due to a chemical gradient in the wind as the stars evolve from oxygen-rich to carbon-rich after undergoing C dredge-up processes due to a recent AGB thermal pulse. Suh (2016) successfully modeled the dust envelope around IRAS 09425-6040 with an outer oxygen-rich shell and an inner carbon-rich shell, validating this hypothesis. Similarly,

the carbon-rich PN BD +30° 3639 exhibits spatially separated carbon-rich PAHs and oxygen-rich silicate dust (Guzman-Ramirez et al. 2015). The post-AGB binary HR 4049 is a peculiar example of a depleted oxygen-rich star with a featureless mid-IR spectrum possibly resulting from amorphous carbon masking the spectral features from silicates (Acke et al. 2013). While it is possible we may be observing the stars in transition from oxygen-rich to carbon-rich, it would require

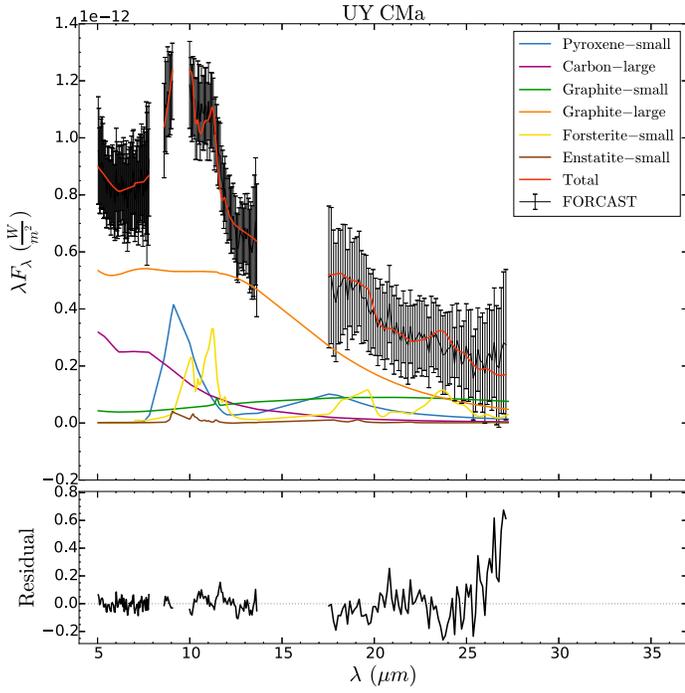


**Figure 8.** Effect on the fit when adding dust species to the best model of AC Her. The best-fitting coefficient,  $c_i$ , for alumina-small in panel (c) was 0.00.

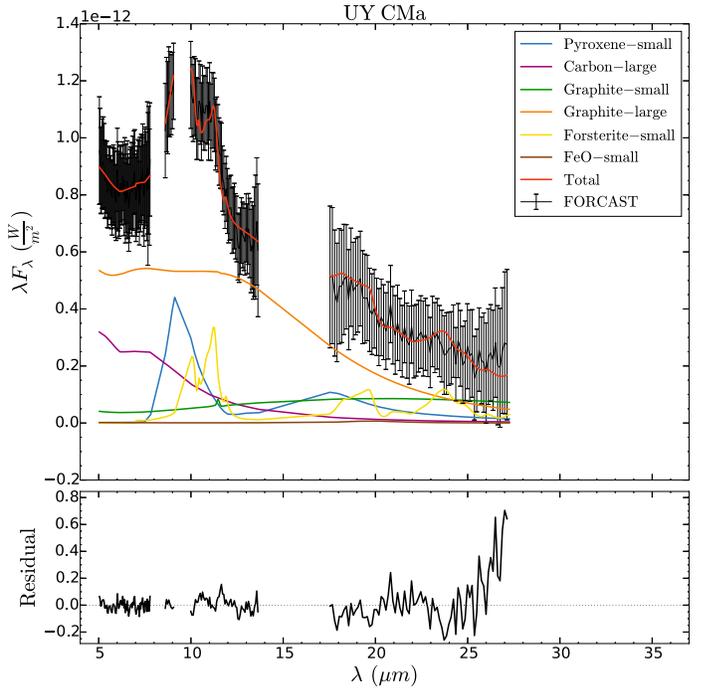
that all of these stars result from a narrow range of masses that terminate AGB evolution just as the carbon exceeds the oxygen abundance. A more plausible explanation for the dual chemistry mineralogy is that the dust formed in a common envelope environment of a binary system where the carbon and oxygen abundances can rapidly change. This mechanism has been invoked as the possible origin of post-AGB disks (Kashi & Soker 2011; Lü et al. 2013; Hardy et al. 2016).

### 6.3. Viewing Effects

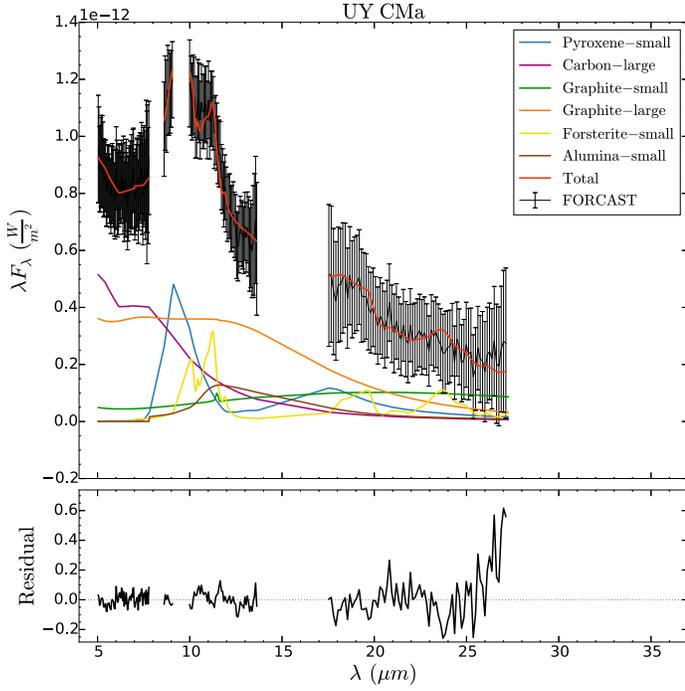
Most of the FORCAST continua are well described by two Planck functions, suggesting that we are viewing the systems from a nearly face-on orientation and see both the inner ( $\sim 1000$  K) and outer ( $\sim 250$  K) regions of the disks. Our results are corroborated by Hillen et al. (2015) who used a radiative transfer code to model the dust around AC Her as a highly



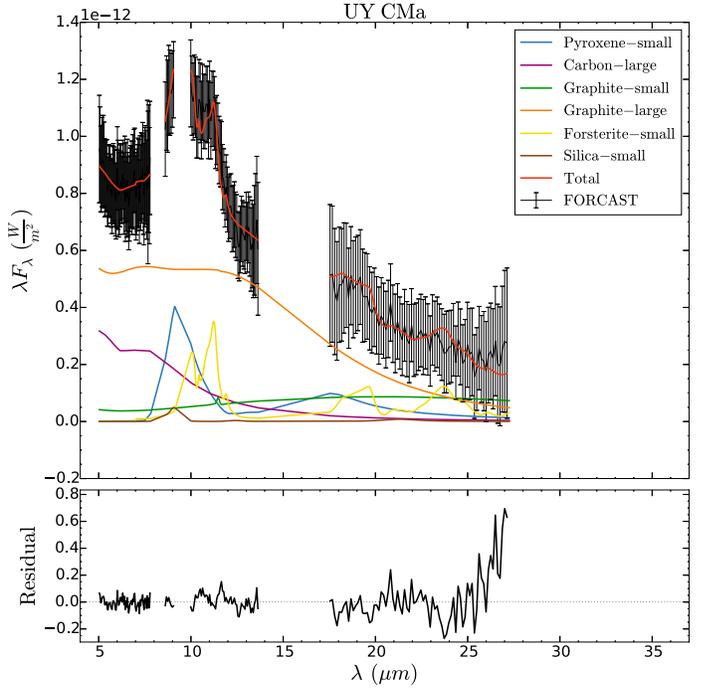
(a) Enstatite-small added;  $\chi_{\text{red}}^2 = 0.085$



(b) FeO-small added;  $\chi_{\text{red}}^2 = 0.085$



(c) Alumina-small added;  $\chi_{\text{red}}^2 = 0.080$



(d) Silica-small added;  $\chi_{\text{red}}^2 = 0.085$

**Figure 9.** Effect on the fit when adding dust species to the best model of UY CMa.

evolved (i.e., mm sized grains), circumstellar disk with an inclination of  $50^\circ \pm 8^\circ$  and found good agreement with observations. Bujarrabal et al. (2007) detected an extended bipolar outflow and an unresolved, compact (presumably disk) component around V441 Her with an inclination of  $\sim 75^\circ$ . The typical uniform disk diameter of the N-band emission region of the RV Tauri stars interferometrically observed by Hillen et al.

(2017) is  $\sim 40$  mas, making it difficult to verify the inclination angle of other systems.

#### 6.4. Limitations of the Fit

It is worth mentioning the various difficulties encountered in the modeling of the dust species present around our program

stars and the limitations of our simplified model. As mentioned in Section 3, less than half of the program stars were observed with the G6 grism (28.7–37.1  $\mu\text{m}$ ). The addition of these data would have aided the modeling and identification of the minerals as crystalline silicates have prominent emission features in this region. Similarly, the lack of coverage from 14–17  $\mu\text{m}$  between grisms G3 and G5 made it more difficult to constrain the abundance of amorphous silicates which have emission features around 17  $\mu\text{m}$ . Amorphous carbon and metallic iron, on the other hand, lack prominent IR features and our model could be fitting these species to the overall dust continuum or temperature gradient, thereby increasing the relative abundances.

Because our model included 13 dust species and some of the program stars had relatively low signal-to-noise FORCAST spectra, it was difficult to confirm the uniqueness of our models. For spectra with high signal-to-noise ratios, excluding dust species from the model had a noticeably negative impact on the goodness of fit (see Figure 6) whereas, for low signal-to-noise spectra, the exclusion of dust species changed the overall shape of the fit but the  $\chi_{\text{red}}^2$  values changed very little (see Figure 7). For both high and low signal-to-noise spectra, adding additional dust species improves the goodness of fit very little, if at all. We checked this by including Mg-rich crystalline enstatite (Jaeger et al. 1998a), iron oxide (FeO; Henning et al. 1995), amorphous alumina (porous; Begemann et al. 1997), and amorphous silica (Henning & Mutschke 1997) to our model one at a time. The  $\chi_{\text{red}}^2$  values either remained the same or marginally improved for all of the program stars. The volume fraction of the added dust species depended on the signal-to-noise ratio of the spectra, with lower signal-to-noise spectra generally including 1%–4% by volume and higher signal-to-noise spectra including  $\leq 1\%$  by volume. Figure 8 illustrates the effect of adding additional dust species to AC Her (high signal-to-noise spectra) and Figure 9 illustrates the effect of adding additional dust species to UY CMA (low signal-to-noise spectra).

### 6.5. Crystallinity of the ISM

An upper limit on the degree of crystallinity of silicates in the diffuse ISM has been estimated by Kemper et al. (2004) to be  $0.2\% \pm 0.2\%$  by mass. This estimate is similar to the average crystalline silicate fraction we find for our program stars of  $0.4\% \pm 0.05\%$  by volume. We concur with Kemper et al. (2004) that this suggests crystalline material is either diluted in the ISM by other amorphous grain-producing processes such as supernovae, or there is an amorphization process that occurs in the ISM on a shorter timescale than the destruction timescale, possibly heavy ion bombardment. Our mineralogy model of each star predicts that the majority of the dust is in the form of graphite and amorphous carbon. On average our model predicts  $80\% \pm 1\%$  graphite and amorphous carbon, and  $57\% \pm 1\%$  graphite around each star by volume. This large volume fraction of graphite and carbon around post-AGB stars may help explain the 2175 Å bump observed in the interstellar extinction curve which is possibly due to these two species (Rouleau et al. 1997; Duley & Seahra 1998; Bradley et al. 2005).

## 7. Conclusion

We have presented a first look at data obtained with SOFIA FORCAST of 15 RV Tauri and 3 SRd variable stars. These data have demonstrated the diversity of dust features present in these systems, possibly tracing the evolutionary track from post-AGB star to PN. These observations of IR excess support the hypothesis that the systems in question are at an advanced stage in their transition to PNe. Our main conclusions can be summarized as follows.

1. Almost all of the stars observed display a 10  $\mu\text{m}$  and/or 20  $\mu\text{m}$  emission feature. For most of the stars observed, the FORCAST continua are well described by two Planck functions, one at  $\sim 1000$  K and one at  $\sim 250$  K, with a majority of the dust ( $\sim 97\%$ ) in the cooler form. A single Planck function fit the underlying continuum of o<sup>1</sup> Cen and V Vul, indicating that these systems may be in the final stages of disk dissipation.
2. Our mineralogy model indicates the presence of both carbon-rich and oxygen-rich dust species with a majority of the dust,  $80\% \pm 1\%$  by volume on average, in the form of amorphous carbon and graphite. All of the stars display this dual chemistry circumstellar dust. This requires that either these stars result from a narrow range of masses that terminate AGB evolution just as the carbon exceeds the oxygen abundance, or the formation process is not single-star evolution. We speculate that the formation process is common envelope evolution.
3. The spectra do not exhibit any obvious crystalline emission features and our model only predicts UY CMA, RU Cen, and AC Her to have crystalline forsterite at volume fractions of  $4\% \pm 0.9\%$ ,  $1\% \pm 0.3\%$ , and  $1\% \pm 0.1\%$ , respectively.
4. Most of the spectra show a low peak-to-continuum value (i.e.,  $< 2.0$ ) and a low 10/20  $\mu\text{m}$  ratio (i.e.,  $< 50$ ), indicating that the grains are relatively large and have undergone significant processing, supporting the hypothesis that the dust is constrained to a Keplerian disk and that we are viewing the heavily processed, central regions of the disk from a nearly face-on orientation.
5. The average composition of the SRd variables contains 8% more small carbon dust and less graphite (14% less of the small species and 5% less of the large) than the average composition of the RV Tauri stars. Of the three SRd variables modeled in this work, none of them contained the large carbon species; on average the RV Tauri stars contained 13% by volume. Overall the average volume fraction of large grains for the SRd variables was 16% compared to 30% for the RV Tauri stars. The paucity of large grains around any of the SRd variables supports the hypothesis that these stars are single-star systems.
6. Between the featureless IR dust species, amorphous carbon is included in more of our models (16 out of 17) than metallic iron (4 out of 17).

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*Facility:* SOFIA (FORCAST).

*Software:* pymiecoated (Leinonen 2012).

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