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# X-RAY STUDIES OF SONE SULPHUR-CONTAINING ORGANIC COMPOUNDS. 

by

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A thesis submitted to the University of Keele in partial fulfilment of the requirements for the degree of Doctor of Philosophy.

[^0]FOREWARD.

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## ABSTRACT.

In this thesis, X-ray crystal structure analyses of the desaurin from acetophenone, 3-phenyl-l-propene-1,3-dione 1-(dimethyl mercaptole) and the 1,2,4-trithiacyclopentane from pinacolone, are described.

The X-ray analysis of the desaurin from acetophenone confirms the desaurin structure and shows that the molecule is the trans isomer. The 1,3-dithiete ring is planar, with carbonsulphur bond lengths of 1.766 and $1.764 \AA$, and a carbon-sulphurcarbon bond angle of $82.0^{\circ}$.

The structure of 3-phenyl-l-propene-1,3-dione 1-(dimethyl mercaptole) was determined employing the symbolic addition procedure, using one symbolic assignment. The sulphur-( $\left.\underline{s p}^{2}-\right)$ carbon bond lengths are 1.748 and $1.745 \AA$, and the sulphur-carbon-sulphur bond angle is $215.1^{\circ}$.

The diffraction study of the $1,2,4$-trithiacyclopentane from pinacolone confirms the thrithiacyclopentane structure and shows that the molecule is a cis isomer. The heterocyclic system is non-planar; the sulphur-sulphur bond length is $2.103 \AA$ and the mean carbon-sulphur bond length is $1.740 \AA^{\circ}$.

In each molecule, the $\alpha, \beta$-unsaturated carbonyl systems adopt an s-cis conformation and this system and the sulphur atoms are approximately coplanar. The intramolecular sulphur-oxygen distances in the three molecules are considerably less than the sum of the van der Waals radii of sulphur and oxygen; similar observations in related systems have been taken as evidence of sulphur-oxygen interaction and this interpretation is discussed.

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## CHAPTER 1.

## INTRODUCTION.

The desaurins and 1,2,4-trithiacyclopentanes.
Bergreen (1) found that when deoxybenzoin was treated with thiophosgene in the presence of sodium ethoxide, an insoluble yellow product was obtained, which he suggested was phenylbenzoylthioketen I. However Meyer. (2), after a molecular weight


I
determination, proposed the dimeric structure $I I_{a}$, and terned compounds of this type "desaurins" on account of their yellow


II $\quad a, R=R^{\prime}=P h$.
colour. Eleven other desaurins, including the desaurin from acetophenone III, were prepared by early German workers $(3,4,5,6)$,

[^1]whilst a number of other such compounds have more recently been prepared by Gompper and töpfl (7). The few original degradations


III
or other studies on the desaurins (8) provide little proof for such a novel heterocyclic system and it was not until a more recent re-examination of these compounds by Yates and Moore (9) that Meyer's structural proposal was shown to be correct.

It is evident that in the desaurins two geometrical isomers, cis and trans, are possible; thus for the desaurin from acetophenone III, the two isomers possible are IV and V. In the preparation
$\mathrm{Ph} \cdot \mathrm{CO}^{\circ}$


IV


V
of each desaurin only one geometrical isomer has been isolated; the nature of the isomer formed in each case cannot be easily ascertained except from a crystal structure analysis.

Kelber (6) and later Gompper (7) showed that oxidation

[^2]of benzoyldithioacetic acid with bromine or ammonium peroxydisulphate yields a yellow product, for which Gompper proposed the 1,2,4-trithiacyclopentane structure VI. By using suitable alternative ketones, various 1,2,4-trithiacyclopentanes,


VI


VII
including that from pinacolone VII, have been prepared utilising the method of Gompper. For each trithiacyclopentane, three geometrical isomers, (two cis and one trans), are possible; thus for the 1,2,4-trithiacyclopentane from pinacolone VII, the three isomers are $V I I_{a}, V I I_{b}$ and $V I I_{c}$.

$\mathrm{VII}_{a}$

$V I I_{b}$

$\mathrm{VII}_{c}$
A recent extensive examination of the desaurins and $1,2,4-$ trithiacyclopentanes by Lynch (10); has shown that certain

[^3]trithiacyclopentanes, notably those from acetophenone VI, p-methylacetophenone VIII and deoxybenzoin IX, when prepared using the


VIII


IX
method of Compeer, yield two isomeric forms, separable by fractional crystallisation. However, the preparation of the trithiacyclopentane from pinacolone yielded a single isomer, which was identified as one of the cis forms, $\mathrm{VII}_{a}$ or $\mathrm{VII}_{b}$, on the basis of its n.m.r. spectrum.

Lynch (10) also noted that the desaurins and trithiacyclopentane exhibit unusually weak carbonyl absorption in their infra red spectra. It is evident that if the $\alpha, \beta$-unsaturated carbonyl systems of these compounds exist in an s-cis conformation, the possibility of sulphur-oxygen interaction arises which could account for the unusual spectral properties found. Examples of close inframolecular approach of oxygen and sulphur atoms have been previously reported. An X-ray diffraction study of 2,5-dimethyl-dithiofurophthene, $X$ (Il), showed that in spite of a remarkably short


X


XI
sulphur-oxygen distance of $2.41 \AA$, (the sum of the van der Wails

[^4]radii of oxygen and sulphur is taken to be $3.05 \AA$ ), the molecular skeleton was planar. Attention was also drawn (ll) to the close structural similarity between X and dimethyl-thiothiophthene XI, (the structure analysis of which had been previously reported (12)), and the existence of a partial sulphur-oxygen bond was proposed.

A study of the methyl ester of o-nitrobenzenesulfenic acid, XII (13), showed a non-bonded sulphur-oxygen distance of $2.44 \AA$,


XII
which together with the planarity of the molecule, (except for the methyl group), was taken to indicate the existence of an attractive interaction between oxygen and sulphur.

X-ray diffraction studies on two compounds having some structural similarity to the desaurins and trithiacyclopentanes have recently been published ${ }^{\neq}$. The structure analysis of the
(12) S. Bezzi, M. Marmi and C. Garbulio, Nature, 182, 247 (1958).
(13) W. C. Hamilton and S. J. Laplaca, J. Am. Chem. Soc., 86, 2289 (1964).
\# These studies were published after the structure of the desaurin from acetophenone had been reported in short note form, and we are grateful to the authors for, supplying us with'structural details prior to publication.

2-methylene-1,3-dithiacyclobutane derivative, XIII (14), showed


XIII
the $\alpha, \beta$-unsaturated carbonyl system to adopt the s-cis conformation. In addition this system and the sulphur atoms were found to be approximately coplanar, (the mean deviation from the plane being $0.04 \AA$ ), which resulted in a short sulphur-oxygen distance of $2.63 \AA$. An X-ray diffraction study of 2-desylidene-1,3-dithiolane, XIV (15,16), yielded a somewhat similar result. The


XIV
Sl-CaC-C=O system is again planar and the sulphur-oxygen distance is correspondingly short at $2.70 \AA$. In this molecule, however, the second sulphur atom, $S 2$, is out of the least squares plane by $0.24 \AA$. The aims of the present work.

This thesis describes the crystal structure determinations, by single crystal X-ray diffraction methods, of the desaurin from

[^5]acetophenone III, the 1,2,4-trithiacyclopentane from pinacolone VII and of 3-phenyl-1-propene-1,3-dione 1-(dimethyl mercaptole) XV。


XV

The main purpose of the work is to confirm the desaurin and 1,2,4-trithiacyclopentane structures, to establish their stereochemistry and obtain dimensional details of the heterocyclic systems. A secondary purpose is to determine whether sulphuroxygen interaction, which had been hypothesized from spectral data, is present. 3-phenyl-l-propene-1,3-dione 1-(dimethyl mercaptole) bears a structural similarity to 'half' of the molecule of the desaurin from acetophenone, and consequently it was hoped that the structure analysis would provide some interesting comparisons to be made with the desaurin structure.

## CHAPTER 2.

## INTRODUCTION TO THE X-RAY CRYSTALLOGRAPHIC WORK.

Since a crystal can be considered as a continuous periodic electron density distribution in three-dimensions, this density can be represented by a three-dimensional Fourier series in which the structure factors, ( $\mathrm{hkl}^{\prime}$ 's), are the Fourier coefficients. The electron density at a point $(x, y, z)^{\ddagger}$ in the unit cell 18 given by,

$$
\begin{equation*}
p(x, y, z)=\frac{1}{V} \sum_{n} \sum_{-} \sum_{k}^{\infty} \sum_{-\infty} F_{h k l} \cdot \exp [-2 \pi i(h x+k y+l z)], \tag{2.1}
\end{equation*}
$$

where $V$ is the volume of the unit cell. The structure factor is, however, a complex quantity requiring both an amplitude and a phase to express it fully. Although the amplitude of each structure factor can be determined experimentally, the phase cannot, and therefore it is not possible to compute the electron density summation, (2.1), solely from experimental data.

## The Patterson function.

Perhaps the most widely-used method of overcoming the phase problem is by use of the Patterson function. Patterson ( 17,18 ) showed that if a Fourier summation is carried out using $\left|F_{h k l}\right|^{2}$ as coefficients, (which are phaseless), the resulting synthesis reveals information of the orientation and magnitudes

[^6]of interatomic vectors. The Patterson function,
\[

$$
\begin{equation*}
P(u, v, w)=\sum_{n}^{-} \sum_{-\infty}^{\infty} \sum_{k}\left|F_{h k l}\right|^{2} \cdot \cos 2 \pi(h u+k v+1 w) \tag{2.2}
\end{equation*}
$$

\]

exhibits peaks at vector distances, $[(u, v, w)]$, from the origin equal to vector distances between pairs of maxima in the electron density. Although it appears that the structure of any crystal can be easily solved from the Patterson function, the number of peaks in the vector map rises rapidly with the number of atoms in the unit cell; with $N$ atoms in the unit cell there are $N(N-1)$ interatomic vectors, each of which will contribute a peak in the vector map. The height of each Patterson peak is dependent upon the product of the scattering powers of the two atoms involved. In a multi light-atom problem, the Patterson function will therefore consist of a large number of vector peaks, each of approximately the same height. If the experimentally determined $\left|F_{\text {hkl }}\right|^{2}$ are used as coefficients in (2.2), these peaks are diffuse and many will coalesce making it extremely difficult to identify individual interatomic vectors. It is possible, however, to transform, (or 'sharpen'), the $\left|F_{h k l}\right|^{2}$ and thereby increase the resolving power of the Patterson synthesis. A simple sharpening technique (19) is to transform the experimental $\left|F_{n k l}\right|^{2}$ to the theoretical values, $\left(\left|{ }^{S_{F}}{ }_{h k l}\right|^{2}\right)$, they would have been if the crystal was composed of point atoms having no thermal motion,

$$
\begin{equation*}
\left|s_{F_{h k l}}\right|^{2}=\frac{\left|F_{h k l}\right|^{2}}{\sum_{i=1}^{\infty} f_{i}^{2} \cdot \exp \left(-2 B_{o v} \cdot \sin ^{2} \theta / \lambda^{2}\right)} \tag{2.3}
\end{equation*}
$$

(19) M. J. Buerger, "Vector Space", Wiley and Sons, New York, (1959), p. 64.
where, $f_{i}=$ the scattering factor of the ith atom in a unit cell containing $N$ atoms,
$B_{o v}=$ the overall isotropic temperature factor.
By using the sharpened coefficients, $\left|{ }^{s} F_{h k l}\right|^{2}$, in the Patterson synthesis, the electrons of the normal peak are concentrated into a 'sharp' peak, which leads not only to better resolution but also makes the weaker peaks more prominent. One disadvantage of using sharpened coefficients in the Patterson synthesis is that the Fourier series ceases to be convergent, which can result in spurious detail being produced in the background of the synthesis. Even with the superior resolution of a sharpened Patterson function, the solution of a multi light-atom problem from its Patterson is still extremely difficult. The problem can be simplified, however, if some features of the molecular structure are known. If the molecule contains a rigid group of atoms, and the dimensions of the rigid group are known, (for example a benzene ring), the complete structure analysis can be carried out by a systematic search in the Patterson function (20,21).

Although many light atom structures have been solved in the past by either interpreting a suitable sharpened Patterson function, or by using both the Patterson and a search technique, at the present time many structures of this type are subjected to direct phase determination methods, (see later).
(20) C. E. Nordman, Acta Cryst., 15, 18 (1962).
(21) C. E. Nordman and K. Nakatsu, J. Am. Chem.․Soc., 85, 353 (2963).

## The Patterson 'heavy-atom' method.

The greatest use of the Patterson synthesis is made when a small number of atoms in the unit cell are of relatively high atomic number, ('heavy'), compared with the majority. It is then possible to locate the heavy atoms easily from the Patterson function, since the vector peaks between different heavy atoms stand out predominently against the background of light atom interactions.

Having located the positions of the heavy atoms in this manner, there are then two main courses that may be followed to locate the remaining light atoms in the structure. The choice between the two methods depends upon the size of the contribution made by the heavy atoms to the structure factors. With one heavy atom per unit cell, the structure factor can be sub-divided into heavy atom, (H), and light atom, (L) contributions,

$$
F_{h k l}=f_{H} \exp 2 \pi i\left(h x_{H}+k y_{H}+l z_{H}\right)+\sum_{i=1}^{N-4} f_{L_{i}} \exp 2 \pi i\left(h x_{L_{i}}+k y_{L_{i}}+l z_{L_{i}}\right),(2.4)
$$

where $f_{H}$ and $f_{L_{i}}$ are the scattering factor of the heavy and ith light atoms respectively.

If $f_{H}$ is very much greater than $f_{L}$, the heavy atom contribution will tend to dominate the right hand side of equation (2.4), except in the relatively rare cases when all the light atoms of the residue scatter in phase to give a contribution greater than that of the heavy atom. In the case of a centrosymmetrical crystal, where all structure factors are real, the result is especially useful, as the signs of a high proportion of $\mathrm{F}_{\mathrm{hkl}}$ are determined
solely from the heavy atom contribution. Allocating the phases so determined to the observed structure amplitudes allows an initial electron density synthesis to be computed. Because the $\left|F_{h k l}\right|$ data contain a light atom contribution, this synthesis, in favourable cases, will reveal light atom locations in addition to returning the heavy atom positions. Lipson and Cochran (22) suggest that the heavy atom method is most powerful, (in a centrosymmetrical structure), when the sum of the squares of the atomic numbers of the heavy and light atoms are equal. Under these conditions approximately three quarters of the signs determined solely from the heavy atom contribution are correct; half the signs are correct, on average, because the two contributions will have the same sign, and half of the remainder, on average, will be correct because the contributions of the light atoms are opposite in sign to, but less than, the contribution from the heavy atoms. In any particular example the fraction of structure factors having the same sign as the heavy atom contribution can be assessed using the more exact relationships of Sim (23).

The application of the heavy atom method to non-centrosymmetrical structures is not as straightforward as in the centrosymmetrical case, but the basic principles outlined above remain the same. In this case each structure factor is complex, posessing a phase which can assume any value between 0 and $2 \pi$, and not merely 0 or $\pi$ as in the centrosymmetrical case. There
(22) H. Lipson and W. Cochran, "The Determination of Crystal Structures", G. Bell and Sons Ltd. London, (1953), p. 207. (23) G. A. Sim, Ácta Cryst., 10, 177 (1957).
is therefore a range of errors between the correct phases of the structure factors, and the phases calculated from the heavy atom positions. However for a proportion of the structure factors, the 'heavy atom phases' so determined will be sufficiently close to their correct value that an initial electron density synthesis will show some, if not all, of the light atom locations. For any specific example, the fraction of the structure factors, $N(\Delta)$, for which the phase calculated from the heavy atom contribution is within $\pm \Delta$ of its true value has been calculated by $\operatorname{Sim}(24)$.

When the heavy atom contribution does not greatly dominate the right hand side of equation (2.4), a higher fraction of the structure factors will have a phase different from that determined solely from the heavy atom locations. An initial electron density synthesis, based on the heavy atom phases, will therefore contain much spurious detail and unambiguous identification of the light atoms may be impossible. When this is the case, information about the positions of the light atoms can often be found by reverting back to the sharpened Patterson function and using the vector convergence method (25), or computing the Buerger minimum function (26).

The heavy atom method was employed in the structure analyses of the desaurin from acetophenone, (Chapter 5), and the 1,2,4-trithiacyclopentane from pinacolone, (Chapter 12).

[^7]Phase determination by direct methods.
Direct methods are those in which an attempt is made to derive the phases of the structure factors without previously postulating any atomic positions.

The first contribution in this field was made by Harker and Kasper (27), who showed that relationships could be formulated between certain $F_{\vec{h}}$ and $\left|F_{\vec{h}}\right|^{2}$ which imposed limitations on the phase of some structure factors. The relationships, in the form of inequalities which vary with the symmetry of the crystal under consideration, were applied in practice by Gillis (28), who succeeded in determining the signs of about forty structure factors for the oxalic acid dihydrate crystal, and by Kasper et al. (29), who used them to solve the structure of decaborane. Karle and Hauptman (30), using the criterion that the electron density must be everywhere non-negative, derived a set of inequality relationships which are valid for all the space groups, and showed that the introduction of specific symmetry elements to the set produced the Harker-Kasper inequalities which had been previously formulated.

A major impetus to the development of direct methods was initiated by Sayre (31), who examined relationships between the $F_{\vec{h}}$ 's of a structure and those of the same structure in which the atoms were replaced by 'squared atoms'. If the structure

[^8]consists of equal, resolved atoms, (so that the electron density and the squared density have an identical distribution of peaks and differ only in peak shape), Sayre showed, for a centrosymmetrical crystal, that the sign relation,
\[

$$
\begin{equation*}
s F_{\vec{h}}=s \sum_{\vec{k}} F_{\vec{k}} \cdot F_{\vec{h}-\vec{k}} \tag{2.5}
\end{equation*}
$$

\]

could be formulated, ( $s$ is an abbreviation for 'sign of'). Sayre illustrated the value of (2.5) by solution of the centric Okl projection of hydroxyproline (31), which he achieved by finding a consistent set of signs for some of the larger $F_{h}$ which satisfied the equation. Relationship (2.5) is true when the summation extends over all $\vec{k}$ and the atoms are equal and resolved; however Sayre, by his work on hydroxyproline, and Zachariasen (32), showed in practice that $(2.5)$ remained valid when only structure factors of large magnitude are used in the summation and the atoms of the molecule are of more than one chemical type.

Hauptman and Karle (33) derived, by mathematical analysis, a number of relationships between structure factors which led them to the conclusion that the phase problem was soluble if sufficient data were available. Since $\left|F_{h}\right|$ decrease as the Bragg angle increases, (because of the decline of scattering factor with increasing $\theta$ ), Hauptman and Karle in their work sharpened $\left|F_{\vec{h}}\right|$ to normalised structure amplitudes $\left|E_{h}\right|$,

[^9]\[

$$
\begin{equation*}
\left|E_{\hat{h}}\right|^{2}=\frac{\left|" F_{\vec{h}}\right|^{2}}{\varepsilon_{\vec{h}} \sum_{i=1}^{N} f_{i \vec{h}}^{2}} \tag{2.6}
\end{equation*}
$$

\]

where, $\left|" F_{h}\right|=$ the observed structure amplitude, placed on an absolute scale and suitably corrected for atomic thermal motion,

$$
\begin{aligned}
\varepsilon_{\vec{h}}= & \text { an integer which corrects for space group } \\
& \text { extinctions, }
\end{aligned}
$$

$f_{i}=$ the scattering factor of the fth atom in a unit cell containing $N$ atoms.

The 'probability equivalent' of the Sayre equation derived by Hauptman and Karle is,

$$
\begin{equation*}
\Sigma_{2}: \quad s E_{\vec{h}} \sim s \sum_{\vec{k}_{r}} E_{\vec{k}} \cdot E_{\vec{h}-\vec{k}} \tag{2.7}
\end{equation*}
$$

where the symbol $\vec{k}_{r}$ implies that $\vec{k}$ ranges only over those vectors associated with large $|E|$ values. Probability formulae for evaluating the reliability of a sign determined by application of $\Sigma_{2}$ were derived by Hauptman and Karla (33) and also by Cochran and Woolfson $(34,35)$. For $\Sigma_{2}$ the associated probability function, $\mathrm{P}+(\overrightarrow{\mathrm{h}})$, which determines the probability that the sign of $\mathrm{E}_{\mathrm{h}}$ be positive, is given by $(34,35)$,

$$
\begin{equation*}
P+(\vec{h})=1 / 2+1 / 2 \tanh \left[\frac{\sigma_{z}}{\sigma_{2}^{3 / 2}}\left|E_{\vec{h}}\right| \sum_{\vec{k}} E_{\vec{k}} \cdot E_{\vec{h}-\vec{k}}\right], \tag{2.8}
\end{equation*}
$$

where $\sigma_{n}=\sum_{i=1}^{N} z_{i}^{n} ; z_{i}$ being the atomic number of the fth atom.

[^10]It is evident, for a centrosymmetrical crystal, that the signs of normalised structure factors can be determined by using the $\Sigma_{2}$ relationship (2.7), and that a probability assessment of each generated sign may be made by employing the probability function (2.8). Application of $\Sigma_{2}$ is initialised employing certain large $\left|E_{\vec{k}}\right|$ 's whose phases may be assigned arbitrarily to specify the origin, (c.f. Chapter 9), and is supplemented, in the symbolic addition procedure $(36,37)$, by additional larger $\left|E_{\vec{k}}\right|$ 's whose phases are denoted by symbols. After repeated $\boldsymbol{\Sigma}_{2}$ application the signs of numerous $E_{\vec{h}}$ are generated, (either absolutely or in terms of the assigned symbols), and the normalised structure factors are then used as Fourier coefficients to generate an 'E-map' $(38)$. Karle et al. $(37,38)$ have illustrated that in general the atomic positions are readily obtained from a three-dimensional E-map computed with about ten normalised structure factors, (of larger $|E|)$, per atom of the asymmetric unit. If $p$ symbols have been assigned, there will be theoretically $2^{p}$ possible sign combinations to consider; however relationships amongst the symbols frequently occur as the phase determination proceeds and generally the number of E-maps it is necessary to compute is considerably reduced.

The symbolic addition procedure was employed in the structure analysis of 3-phenyl-l-propene-l,3-dione l-(dimethyl mercaptole), and the practical aspects of the procedure are described in Chapter 9.
(36) I. L. Karle and J. Karle, ibid., 16, 969. (1963).
(37) I. I. Karle, K. Britts and P. Gum, ibid., 27, 496 (1964).
(38) I. L. Karle, H. Hauptman, J. Karle and A. B. Wing, ibid., 11, 257 (1958).

## Least squares refinement of structure parameters.

The method of least squares can be used in structure refinement (39) by minimising some function of the observed and calculated structure amplitudes with respect to the atomic parameters. The function most commonly minimised is,

$$
\begin{equation*}
R^{\prime}=\sum_{q} w(h k l)\left[\left|F_{0}(h k l)\right|-\left|F_{c}(h k l)\right|\right]^{2}, \tag{2.9}
\end{equation*}
$$

where $\left|F_{0}\right|$ and $\left|F_{c}\right|$ are the observed and calculated structure amplitudes, $q$ is the total number of independent $\left|F_{0}\right|$ data and $w(h k l)$ is a weight applied to each tern which depends upon the reliability of the $\left|F_{0}(h k l)\right|$ datum. Although strictly w(hkl) is proportional to $l /\left(\sigma\left|F_{0}(h k l)\right|\right)^{2}$, (where $\sigma\left|F_{0}\right|$ is the standard deviation of the observed structure amplitude), frequently, in crystal structure refinement, only relative weights can be estimated and it is often assumed that $\sigma\left|F_{0}\right|$ depends only on $\left|F_{0}\right|$.

The trial set of $\underline{n}$ atomic parameters, $\left(p_{1}, p_{2} \ldots p_{n}\right)$, require corrections to be applied, $\left(\Delta p_{i}\right)$, to adjust them to their 'most probable' values, ( $\left.p_{1}+\Delta p_{1}, \ldots p_{n}+\Delta p_{n}\right)$, which will result in the required minimisation of (2.9). If the parameter shifts required are small, (so that Taylor's expansion to first order may be employed), (2.9) can be rewritten as,
$R^{\prime}=\sum_{q} w\left[\left|F_{0}\right|-\left|F_{c}^{\prime}\right|+\frac{\partial F_{c}^{\prime}}{\partial p_{1}} \cdot \Delta p_{1}+\frac{\partial F_{c}^{\prime}}{\partial p_{2}} \cdot \Delta p_{2} \ldots .+\frac{\partial F_{c}^{\prime}}{\partial p_{n}} \cdot \Delta p_{n}\right]^{2}(2.10)$
where $F_{c}^{\prime}$ is the value calculated using the trial set of parameters.
(39) E. W. Hughes, J. Am. Chem. Soc., 63, 1737 (1941).

Function (2.10) will be minimum when the partial derivatives with respect to each of the $n$ parameters vanish, i.e.,

$$
\begin{equation*}
\frac{\partial R^{\prime}}{\partial P_{1}}=0 ; \frac{\partial R^{\prime}}{\partial P_{2}}=0 \ldots \ldots \cdot \frac{\partial R^{\prime}}{\partial P_{n}}=0 \tag{2.21}
\end{equation*}
$$

Therefore the $\underline{n}$ conditions of type (2.11) may be written as a series of $\underline{n}$ equations, (the "normal equations"),

$$
\sum_{q} w\left[\left(\frac{\partial F_{c}^{\prime}}{\partial p_{1}} \left\lvert\,\left(\frac{\partial F_{c}^{\prime}}{\partial p_{1}}\right) \Delta p_{1}+\left(\frac{\partial F_{c}^{\prime}}{\partial p_{1}} \left\lvert\,\left(\frac{\partial F_{c}^{\prime}}{\partial p_{2}} \left\lvert\, \Delta p_{2}+\ldots\left(\left|\frac{\partial F_{c}^{\prime}}{\partial p_{1}}\right|\left|\frac{\partial F_{c}^{\prime}}{\partial p_{n}}\right| \frac{\Delta p_{n}}{}\right]\right.\right.\right.\right.\right.\right.\right.
$$

$$
=-\sum_{q} w\left(\left|F_{0}\right|-\left|F_{c}\right|\right)\left|\frac{\partial F_{c}}{\partial p_{i}}\right|
$$

$$
\sum_{q} w\left[\left|\frac{\partial F_{c}}{\partial p_{n}}\right|\left|\frac{\partial F_{c}}{\partial p_{1}}\right| \Delta p_{1}+\left|\frac{\partial F_{c}^{c}}{\partial p_{n}}\right|\left|\frac{\partial F_{c}^{c}}{\partial p_{2}}\right| \Delta p_{2}+\cdots\left(\frac{\partial F_{c}}{\partial p_{n}}\left|\frac{\partial F_{c}^{\prime}}{\partial p_{n}}\right| \Delta p_{n}\right]\right.
$$

$$
=-\sum_{q} w\left(\left|F_{0}\right|-\left|F_{c}\right|\right)\left|, \frac{\partial F_{c}^{1}}{\partial p_{n}}\right|
$$

The normal equations represent a set of $\underline{n}$ simultaneous equations In $n$ unknowns and may therefore be solved for $\Delta p_{1}, \Delta p_{2} \ldots \Delta p_{n}$. Since (2.10) is an approximation, the required minimisation of the function (2.9) is not normally achieved by a single cycle of least squares refinement. However, by applying the calculated $\Delta p_{i}$ and repeating the procedure, additional improvement in the parameters can be obtained. The standard deviation of each atomic parameter is calculable from the least squares residuals, and
in practice refinement is continued until the parameter shifts calculated represent only a fraction of their standard deviation.

Computational work.
All computations were carried out on the I.B.M. 7094
at the University of Toronto.
The program used for structure factor calculations and
full matrix least squares refinement was a modified version of OR-FLS $(40,41)$. One modification, due to Shiono, was that reflection data were read into the program one at a time from magnetic tape; consequently there was no restriction on the number of observations. In the orizinal version of OR-FLS the value of the atomic scattering factor for each reflection was obtained by linear interpolation from a stored table of 32 values for ( $\sin \theta / \lambda$ ) from 0.0 to 1.55 in 0.05 intervals. Because of the coarse mesh of this table, values interpolated at low angles . could be significantly in error, and a second Shiono modification was to replace the original 'f' table with the 51 scattering factor values for $\sin \theta$ from 0.0 to 2.0 in intervals of 0.02 . The 51 unit table for sulphur, carbon and oxygen were prepared using the Aitken polynomial method (42) on the scattering factor values listed in International Tables (43) ${ }^{\ddagger}$. To the scattering curve for sulphur 0.3 electrons were added over the whole $\sin \theta$
(40) W. R. Busing, K. O. Nartin and H. A Levy, 'Oak Ridge National Laboratory, Oak Ridge, Tennessee, (1962).
(41) R. Shiono, Crystallography Laboratory, University of Pittsburgh, Pittsburgh, (1963).
(42) See, for example, T. A. Bak and J. Lichtenberg, Wathematics for Scientists", Benjamin, (1966), p. 459.
(43) International Tables for X-Ray Crystallogrephy, Volume III, The Kymoch Press, Birmingham, (1962), p. $201 ;$ (a) p. 214. * A listing of the interpolation program is given in Appendix $V$.
range to allow for the real part of the anomalous scattering ( $43_{a}$ ). Patterson, electron density, difference and the E-map syntheses were computed employing "A three-dimensional Fourier synthesis program" written by Shiono (41).

## EXPERINENTAL - PART I

THE X-RAY STRUCTURE ANALYSIS OF THE DESAURIN FROM ACETOPHENONE.

## CHAPTER 3.

## PRELIMINARY X-RAY WORK.

The desaurin from acetophenone, prepared by the method of Kelber and Schwarz (5), crystallises from nitrobenzene as small plates (44). The major face exhibited is (001), using the axial description given below.

X-ray equipment.
This portion of the work was carried out using a Norelco X-ray generator, operated at 45 kV and approximately 14 mA ; nickel filtered copper radiation was used throughout. The recording device used was a Stoe integrating Weissenberg goniometer. Ilford industrial G X-ray film was used, being processed at $20^{\circ} \mathrm{C}$ with Kodak D19B developer and Kodak F-5 acid-hardening fixer.

## Crystal density.

The density of the crystals was determined by flotation in a benzene and carbon tetrachloride mixture at room temperature, the density of the liquid mixture in which the crystals remained suspended being found by use of a specific gravity bottle. The average value of the crystal density was found to be $1.45 \mathrm{~g} . \mathrm{cc}^{-1}$.

## Unit cell dimensions and space group.

Unit cell dimensions were originally estimated from zero level Weissenberg photographs and were later determined more accurately from new sets of zero level photographs calibrated
(44) We are grateful to Dr. T. R. Lynch for providing crystals of the desaurin from acetophenone.
with aluminium powder lines. A least squares refinement procedure was used.

$$
\begin{array}{ll}
\text { System monoclinic }(y \text { axis unique }) \\
a=5.470 \AA & (\sigma=0.008) \\
b=4.848 \AA & (\sigma=0.012) \\
c=28.299 \AA & (\sigma=0.008) \\
\beta=96.90 & (\sigma=0.1) \\
v=745.0 \AA^{3} &
\end{array}
$$

The following systematic absences were noted on zero level and equi-inclination Weissenberg photographs;
hkI : no general absences,
hol : absent for $1=2 N+1$,
OkO : absent for $k=2 N+1$.

The space group is therefore $\mathrm{P}_{1} / \mathrm{c}$.

## The number of molecules per unit cell.

The number of molecules per unit cell can be expressed by the formula,
where,

$$
\begin{aligned}
n & =\frac{\alpha_{\text {meas }} \cdot V}{1.66 \cdot M} \\
\alpha_{\text {meas }} & =\text { the crystal density, } \\
V & \left.=\text { the volume of the unit cell, (in } \AA^{3}\right), \\
M & =\text { the molecular weight of the compound. }
\end{aligned}
$$

Using deas $=1.45 \mathrm{g.cc}{ }^{-1}, M=324$ and $V=745.0 \AA^{3}$, the calculated value of $n$ is 2.01. Thus $Z$, the number of molecules per unit cell, is 2.

## CHAPTER 4.

## COLIEGTION OF THREE-DIMENSIONAL INTENSITY DATA.

## Optimum crystal size.

The linear absorption coefficient, $\mu$, is defined by,

$$
\mu=d \sum \rho \mu_{m}
$$

where,

$$
\begin{aligned}
\mathrm{d}= & \text { the crystal density }, \\
\rho= & \text { the fractional weight of each element }, \\
\mathfrak{p}_{m}= & \text { the mass absorption coefficient for } \\
& \text { each element. }
\end{aligned}
$$

Using the mass absorption coefficients listed in International
Tables (45), the calculated value of $\mu$ is $33.15 \mathrm{~cm}^{-1}$. Buerger (46) showed that the X-ray reflections with the longest path through the crystal reach their maximum intensity when the crystal diameter is $2 / \mu$, and suggested that this value be taken as the optimum thickness, $t$, for crystals used in X-ray diffraction work. For the desaurin, $t$ is 0.61 mms . The crystals from which intensity data were collected were much smaller than this optimum size, (the approximate dimensions of the largest crystal used were $0.2_{0} \times 0.2_{0} \times 0.0_{2}$ mms), and no absorption corrections were applied.

Collection of three-dimensional intensity data.
Three-dimensional intensity data were collected utilising
(45) International Tables for X-ray Crystallography, Volume III, The Kynoch Press, Birmingham, (1962), p.162.
(46) M. J. Buerger, "X-ray Crystailography", Wiley and Sons Inc., New York, N.Y., (1940), p.180.
the equi-inclination Weissenberg method. To ensure that the intensity data recorded came within the range where they could be estimated easily by eye, the multi-film technique was employed (47); thus a pack of four films, interleaved with aluminium foil ( 0.001 inches thick), was used for each exposure. The intensity of each reflection was measured by visual comparison with an intensity wedge consisting of twenty-eight 'spots' whose relative intensities were known. This wedge was generated by allowing a reflection from the crystal to pass through its reflecting position for varying numbers of times, the translational motion of the filmholder being retained. The crystal used to prepare the wedge was that from which intensity data were collected, so that the shape of the spots on the wedge were typical of the shape of the reflections on the films to be measured.

Initially, recrystallisation of the desaurin from nitrobenzene produced crystals which were generally multiple, together with some very small crystals which were used for data collection. Crystals were mounted about the $a$ and $\underline{b}$ axes and two series of photographs were recorded, hOl through h3l, and Okl through 4 kl . For each layer the intensity of each reflection was measured on all four films of the pack (48), an average film factor determined, and the average intensity for each plane calculated. The intensity data were corrected for Lorentz and polarisation factors employing the graphical method devised by Cochran (49). The Lp corrected
(47) J. M. Robertson, J. Sci. Instr., 20, 175 (1943).
(48) We are grateful to Dr. T. R. Lynch for these initial measurements.
(49) W. Cochran, J. Sci. Instr., 25, 253 (1948).
intensities, measured on different photographs, were placed on the same arbitary scale by comparing $\left|F_{h k l}\right|^{2}$ of common reflections, the Okl film being taken as the arbitary reference standard. The routes by which this inter-scaling was achieved is shown in Figure 4.1.

Absolute scaling and overall temperature factor.
The intensity data for the crystal, (now on an arbitary scale), were placed on an approximately absolute scale by the method due to Wilson (50). The theory depends on the approximation that,

$$
\left.\left.\langle | F\right|^{2}\right\rangle=\left\langle\sum_{i=1}^{N} f_{i}^{2}\right\rangle
$$

where $\left.\left.\langle | F\right|^{2}\right\rangle$ is the average absolute intensity and ' $f_{i}$ is the scattering factor of the ith atom.
Since $\quad\left|F_{0}\right|^{2}=K .|F|^{2}$,
then $\left.\left.\left.\langle | F_{0}\right|^{2}\right\rangle=\left.K\langle | F\right|^{2}\right\rangle$ and $\left.\left.\langle | F_{0}\right|^{2}\right\rangle=K .\left\langle\sum f_{i}^{2}\right\rangle$
where $\left|F_{0}\right|^{2}$ are the corrected observed intensities, on some arbitary scale, and $K$ is a scaling factor with respect to the absolute data. Since under the conditions of observation the atoms possess thermal motion, a temperature factor, $B_{\text {or }}$, must be introduced to modify the theoretical scattering factors accordlngly, i.e.,

$$
f_{i}=f_{i} \exp \left(-E_{o v} \cdot R H O\right)
$$

where RHO $=(\sin \theta / \lambda)^{2}$. and
$f_{i}=$ the scattering factor for a non-vibrating atom.


| Layer | K' $^{\prime}$ | Layer | K' $^{\prime}$ |
| :---: | :---: | :---: | :---: |
| h01 | 2.51 | 0 kl | 1.00 |
| h11 | 2.15 | 1 kl | 1.11 |
| h21 | 0.85 | 2 kl | 1.03 |
| h31 | 2.29 | 3 kl | 2.30 |
|  |  | 4 kl | 1.08 |

The routes by which inter-film scaling was achieved and the scaling factors determined ( $K^{\prime}$ ).

## Figure 4.1

Therefore,

$$
\begin{aligned}
& \left.\left.\langle | F_{0}\right|^{2}\right\rangle=K \cdot\left\langle\sum_{i=1}^{N} f_{i}^{2}\right\rangle \cdot \exp \left(-2 B_{O V} \cdot R H O\right) \\
& \ln \left[\frac{\left.\left.\langle | F_{0}\right|^{2}\right\rangle}{\left\langle\sum_{i} f_{i}^{2}\right\rangle}\right]=\ln (K)-2 B_{o v} \cdot R H O
\end{aligned}
$$

The hOl zone was divided into five ranges, each containing approximately an equal number of reflections, and the value of $\ln \left[\frac{\left.\left.\langle | F_{0}\right|^{2}\right\rangle}{\left\langle\sum_{i} f_{i}\right\rangle}\right]$
was calculated and plotted against the average

RHO for each range. The resultant plot is shown in Figure 4.2, and yielded values of 14.8 for the scale factor, (with respect to $\left|F_{0}\right|^{2},[\equiv 1 / K]$ ), and $5.0 A^{2}$ for the overall temperature factor, $B_{o v}$. It was noted that almost all reflections for which $2 \theta$ (CuK $)>80^{\circ}$ had immeasurably small intensities; consequently it was decided to exclude these data for initial stages of the structure analysis. Data Set $1,\left(2 \theta(C u K \alpha) \leqslant 80^{\circ}\right)$, comprising 430 independent reflections of which 334 had measurable intensities, were used for the identification of atomic locations using Patterson and electron density syntheses. Further recrystallisations of the desaurin, (from nitrobenzene), yielded a few larger crystals and at a later stage in the structure analysis layers hol through hll were re-collected, (Data Set 2). Intensity data, measured on these photographs, were corrected for Lorentz and polarisation factors using a program written for the I.B.M. 7094 ${ }^{*}$, and were elevated to the absolute scale by comparing $\left|F_{h k l}\right|^{2}$ common to Data Set 1. Of the 1493 independent reflections recorded in Data Set 2; 1020 had measurable intensities.

[^11]
## CHAPTER 5.

## THREE-DIMENSIONAL PATTERSON AND FOURIER SYNTHESES.

The three-dimensional Patterson function.
The Patterson function was computed for a quarter of the unit cell defined by $-0.5 \leqslant u \leqslant 0.5,0.0 \leqslant v \leqslant 0.5$ and $0.0 \leqslant w \leqslant 0.5$, in intervals of 0.05 in $u$ and $v$, and 0.01 in $w$. The $\left|F_{h k l}\right|^{2}$ of Data Set 1 were used as the Fourier coefficients and the computation was carried out using the general Fourier program. $\left|F_{000}\right|^{2}$ was omitted from the calculation.

As stated in Chapter 3, the desaurin from acetophenone crystallises in space group $\mathrm{P} 2_{1} / \mathrm{c}$ with two molecules, (four sulphur atoms), per unit cell. If one sulphur atom is situated at the general position ( $x, y, z$ ) in the unit cell, the remaining three sulphur atoms must be located at space group equivalent positions $(\bar{x}, \bar{y}, \bar{z}),\left(x, \frac{1}{2}-y, \frac{1}{2}+z\right)$ and $\left(\bar{x}, \frac{1}{2}+y, \frac{1}{2}-z\right)$. Peaks in the Patterson function which represent sulphur-sulphur vectors will therefore be located at,


In reported structures of six-membered sulphur-containing heterocyclics, where a divalent sulphur atom is bonded to an
sp$^{2}$-hybridised carbon atom, sulphur-carbon bond lengths in the order 1.76-1.78 $\AA^{\circ}$ and sulphur-carbon-sulphur bond angles of approximately $100^{\circ}$ have been found $(80,81)$. By applying these dimensions to the 1,3 -dithiete ring of the desaurin, it is therefore expected that the sulphur-sulphur inversion peaks, (f's 1 and 2), will be of the order $2.4 \AA$ to $2.8 \AA$ distant from the origin.

Inspection of the Patterson function revealed a prominent peak about $2.6 \AA$ from the origin, which obviously represented the intramolecular sulphur-sulphur vector. This peak and the corresponding peaks on the Harker line, $(0, v, 1 / 2)$, and the Harker section, ( $u, 1 / 2, w)$, which complete the sulphur-sulphur vector set, are listed below. The numbering of the peaks below is in keeping with the system used in Table 5.1; the peak heights listed were measured from the lowest negative trough in the computed function.

|  | U | V | $W$ | Peak <br> height. |
| :--- | ---: | :---: | :---: | :--- |
| $1,(2)$ | -0.380 | 0.000 | 0.046 | 300 |
| 3. | 0.000 | 0.500 | 0.500 | 650 |
| 4. | 0.380 | 0.500 | 0.454 | 280 |

By referring to Table 5.1 it is seen that these vector peaks correspond to an asymmetric sulphur atom situated at (-0.190, 0.000, 0.023) in the unit cell. Since the sulphur atom has a 'pseudo-special' location $(x, 0, z)$, some of the sulphursulphur vectors have a weight different from those given in

Table 5.1. In this special case the two inversion peaks coalesce into a single peak on the mirror plane at $V=0$ having two fold weight, whilst the peak on the Harker line becomes fourfold at $(0,1 / 2,1 / 2)$. Thus it is evident that the vector set identified has relative peak heights which are consistent with the theoretical weights of the interactions.

Structure factor calculations and three-dimensional electron density syntheses.

For space group P21/c, the calculated structure factors, $F_{c}$, are given by,
$(k+1)=2 n \quad ; \quad F_{c}=4 \sum_{i=1}^{N} f_{i} \cos 2 \pi\left(h X_{i}+1 Z_{i}\right) \cdot \cos 2 \pi k Y_{i}$, $(k+1)=2 n+1 ; F_{c}=-4 \sum_{i=1}^{N} f_{i} \sin 2 \pi\left(h X_{i}+1 Z_{i}\right) . \sin 2 \pi k Y_{i}$,
where $f_{i}$ and ( $X_{i}, Y_{i}, z_{i}$ ) are the scattering factor and the fractional coordinates pertaining to the ith atom of an asymmetric unit containing N atoms.

Structure factors were calculated employing a modified version of OR-FLS and the general Fourier program was used to compute electron density syntheses. The two programs were run sequentially, the input data for the Fourier being prepared by OR-FLS and stored out-of-core between programs. The reflection. data used in this portion of the work were those of Data Set 1 and each electron density synthesis was computed within the - bounds $-0.5 \leqslant x \leqslant 0.5,-0.5 \leqslant y \leqslant 0.5$ and $0.0 \leqslant z \leqslant 0.5$, in intervals of 0.025 in $x$ and $y$ and 0.005 in $z$. The synthesis was calculated for twice the asymmetric volume required by the space group as
this enabled an easier comparison to be made between 'true' and 'false' electron density peaks, (discussed below). Isotropic temperature factors used in the structure factor calculations were $4.0 \AA^{\circ}$ for the sulphur atom and $5.0 \AA^{2}$ for the light atoms. The residual index,

$$
R=\frac{\sum\left(\left|F_{0}\right|-\left|F_{c}\right|\right)}{\sum\left|F_{0}\right|},
$$

has reflections of zero measured intensity included with $\left|F_{0}\right|=0.0$.
Structure factors were calculated using the positional parameters of the sulphur atom determined from the Patterson function, $(R=0.66)$. Since the sulphur atoms are located at $\underline{y}=0$, (and $\underline{y}=1 / 2$ ), they make no contribution to structure factors having ( $k+l$ ) $=2 n+1$, (see equation (5.1)), and the initial electron density synthesis, phased by the sulphur atom, has therefore (020) as a false mirror plane. From the distribution of peaks in this synthesis and packing considerations, it was evident the $\alpha, \beta$-unsaturated carbonyl system is in an s-cis conformation, (with the plane of the benzene ring lying approximately parallel to the long $c$ axis). Attention was first focussed on seven peaks corresponding to atoms of the $C(1)-C(2)-$ $C(3)-0^{\ddagger}$ system and the false symmetry related group, (C) (3) lay on the mirror plane); however an unambiguous subdivision of the seven peaks into two unsaturated carbonyl systems related by the false mirror plane proved difficult. Since the sulphur
\# The atomic numbering system employed is shown in Figure 5.1.


The atomic numbering system used.
atoms are not confined to (020) by any space group requirements, the possibility arises that these atoms are located only 'closeto' these planes, (the Patterson function having insufficient resolution to show a small finite character in the $y$ coordinate of the atoms). The asymmetric sulphur atom was allocated $y=0.01$, hoping that this would sufficiently weight the phases of structure factors for which $(k+\ell)=2 n+1$, that one set of three peaks corresponding to a molecular $C(1), C(2)$ and 0 system would increase in density whilst the 'false symmetry' related set would show a reduction in peak density. However the only significant difference in the subsequent electron density synthesis was that the density at one oxygen location was noticeably greater than the density at the 'false symmetry' position. A third electron density synthesis, phased by the sulphur, $(\underline{y}=0.01)$, and 'true' oxygen atoms, $(R=0.58)$, revealed unambiguously the $C(1)$ and $C(2)$ positions, which together with 0 and $C(3)$, (still centred on $y=1 / 2)$, completed the conjugated carbonyl system of the molecule. A fourth structure factor calculation, using $S, O, C(1)$ and $C(2)$, $(R=0.51)$, and electron density synthesis revealed the locations of $C(4)$ and $C(7)$, and still showed $C(3)$ centred on $y=1 / 2,(a$ location which was now accepted). These three atoms were included in a fifth structure factor calculation, $(R=0.43)$, and the remaining four carbon atoms, which completed the asymmetric unit, were identified from the subsequent electron density synthesis. A final electron density synthesis, phased by all eleven atoms of the asymmetric unit, ( $\quad(=0.33$ ), was devoid of spurious detail;
the fractional coordinates of the atoms estimated from this synthesis are listed in Table 5.2. At this stage all the gross features of the molecular structure had been established and $R$ was $0.28^{\neq}$. Least squares refinement of the atomic parameters was now initiated.
$\neq$ The $\left|F_{0}\right|$ were rescaled by equating $\Sigma\left|F_{0}\right|$ and $\Sigma\left|F_{c}\right|$.

| Atom | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ |
| :--- | :--- | :--- | :--- |
| S | -0.210 | 0.010 | 0.023 |
| 0 | -0.175 | 0.345 | 0.097 |
| $\mathrm{C}(1)$ | 0.085 | 0.168 | 0.027 |
| $\mathrm{C}(2)$ | 0.245 | 0.332 | 0.0625 |
| $\mathrm{C}(3)$ | 0.040 | 0.500 | 0.097 |
| $\mathrm{C}(4)$ | 0.173 | 0.660 | 0.136 |
| $\mathrm{C}(5)$ | 0.0175 | 0.707 | 0.169 |
| $\mathrm{C}(6)$ | 0.110 | 0.875 | 0.205 |
| $\mathrm{C}(7)$ | 0.315 | 1.000 | 0.205 |
| $\mathrm{C}(8)$ | 0.465 | 0.992 | 0.162 |
| $\mathrm{C}(9)$ | 0.376 | 0.805 | 0.129 |

Atomic coordinates estimated from the final three-dimensional electron density synthesis.

Table 5.2.

## CHAPTER 6.

## REFINEMENT OF THE STRUCTURE.

The program employed for least squares refinement was OR-FLS, which had previously been used for basic structure factor calculations and which had now been converted into Fortran IV language. This version of OR-FLS was able to handle up to 45 atoms and a maximum of 450 parameters, of which 165 could be refined simultaneously.

The function minimised was $\sum w\left(\left|F_{0}\right|-\left|F_{c}\right|\right)^{2}$. The weighting scheme used throughout the refinement was that due to Hughes (39),

$$
\begin{array}{ll}
w=1 /\left(\left.F_{0}\right|^{2}\right. & \text { if }\left|F_{0}\right| \geqslant 4 F_{\min }, \\
w=1 /\left(4 F_{\min }\right)^{2} & \text { if }\left|F_{0}\right|<4 F_{\min }, \\
w=0.0 & \text { if } F_{0}=0.0 .
\end{array}
$$

For least squares cycles using isotropic temperature factors $F_{\text {min }}$ was set as 1.0 , however it was found that with $F_{\min }=3.0$ R reduced to its lowest value and this latter allocation for $F_{\min }$ was used for cycles in which anisotropic temperature factors were employed.

Two residual indices are quoted throughout this part of the work; $R_{1}$ is the index obtained using all experimentally recorded reflection data, (1493 for Data Set 2; the 478 reflectfons too faint to be measured being included, (initially), with $\left.\left|F_{0}\right|=0.0\right) . \quad R_{2}$ is the index obtained when only those reflections
which had measurable intensities are used in the summation. Refinement employing isotropic temperature factors.

The initial positional parameters used in the refinement were those obtained from the final electron density synthesis, (Table 5.2), with isotropic temperature factors specified as 4.OA ${ }^{2}$ for the sulphur atom and 5.OA ${ }^{2}$ for oxygen and carbon atoms. Two cycles of refinement were performed using Data Set 1 ; in the first only positional parameters were allowed to vary, whilst in the second all 44 atomic parameters were refined. At this stage the more extensive Data Set 2 were introduced for the final phases of the refinement. Although these data had been adjusted to the absolute scale via Data Set 1 , each layer hol through h4l was assigned an individual scale factor, which were treated as variables in the least squares refinement. After a further four cycles of refinement of the 49 parameters, $R_{1}$ and $R_{2}$ had reduced to 0.193 and 0.153 respectively. The maximum shift calculated in the last cycle represented about half of the standard deviation of the parameter, and it was considered that conversion to anisotropic temperature factors could now be safely made. The atomic parameters at this stage are listed in Table 6.1.

## Refinement employing anisotropic temperature factors.

The anisotropic temperature factor, $T(h k l)_{i}$, applied to the contribution by atom $i$ to $F_{c}(h k I)$ is, $T(h k I)_{i}=\exp -1 / 4\left(h^{2} a^{* 2} B_{11_{i}}+k^{2} b^{* 2} B_{22_{i}}+1^{2} c^{* 2} B_{33 i}\right.$

$$
\left.+2 h k a * b * B_{12_{i}}+2 h l a * c * B_{13 i}+2 k I b * c^{*} B_{23_{i}}\right)
$$

|  | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ | $\mathrm{B}_{\mathrm{iso}}\left(\mathrm{A}^{2}\right)$. |
| :--- | :--- | :--- | :--- | :--- |
| S | -0.1965 | 0.0024 | 0.0236 | 3.598 |
| 0 | -0.1502 | 0.3365 | 0.0980 | 4.757 |
| $\mathrm{C}(1)$ | 0.0950 | 0.1608 | 0.0263 | 3.460 |
| $C(2)$ | 0.1948 | 0.3391 | 0.0583 | 3.586 |
| $C(3)$ | 0.0549 | 0.4254 | 0.0959 | 3.371 |
| $C(4)$ | 0.1588 | 0.6310 | 0.1326 | 3.319 |
| $C(5)$ | 0.0367 | 0.6746 | 0.1727 | 4.320 |
| $C(6)$ | 0.1223 | 0.8591 | 0.2071 | 5.182 |
| $C(7)$ | 0.3330 | 1.0141 | 0.2020 | 4.964 |
| $C(9)$ | 0.4603 | 0.9754 | 0.1628 | 4.949 |

Atomic parameters at the conclusion of the refinement using isotropic temperature factors.

## Table 6.1

where $a^{*}, b^{*}$ and $c^{*}$ are reciprocal cell parameters. The Hughes weighting scheme was retained with $F_{\text {min }}$ reset to 3.0 , and reflections of zero measured intensity were again included with $\left|F_{o}\right|=0.0$. On the initial cycle of refinement a singularity was produced in the matrix due to correlation between the $B_{22}$ thermal parameter of the sulphur atom and the scale factors; consequently these parameters were not allowed to refine simultaneously. After four cycles of refinement $R_{1}$ and $R_{2}$ had improved to 0.164 and 0.125 respectively and the parameter shifts were insignificantly small. Reflections of zero observed intensity were now allocated the value $F_{u n}$, suggested by Hamilton (5l) as being the statistically probable value for an unobserved reflection in a centrosymmetrical space group;

$$
F_{u n}(h k l)=\left[\frac{I_{\min }}{3} \cdot \frac{1}{L p_{h k l}} \cdot s\right]^{1 / 2}
$$

where, $\quad I_{\min }=$ the minimum observed intensity on the film,
$L_{p k I}=$ the Lorentz-polaristion factor,
$\mathrm{S} \quad=\mathrm{the}$ film scale factor.
After a further two cycles of refinement $R_{1}$ and $R_{2}$ improved to 0.127 and 0.111 respectively; parameter shifts calculated in the last cycle were at most $25 \%$ of the standard deviation of the parameter and refinement was considered complete. The final positional and thermal parameters and their estimated standard
(51) W. C. Hamilton, Acta Cryst., 8, 185 (1955).
deviations, calculated from the least squares residuals, are given in Tables 6.2 and 6.3.

## The location of hydrogen atoms.

The observed structure amplitudes possess a small contribution from the hydrogen atoms of the molecule. As the positional and thermal parameters of the non-hydrogen atoms were now known accurately, a Fourier synthesis using ( $F_{0}-F_{c}$ ) coefficients, (the difference synthesis), will show the residual electron density in the cell and should exhibit maxima at the locations of the hydrogen atoms.

Peaks attributable to the six hydrogen atoms of the asymmetric half molecule were easily identified in the computed difference synthesis. Their fractional coordinates are listed in Table 6.4, together with the 'theoretical' hydrogen positions calculated from the carbon atom locations assuming a carbonhydrogen bond length of $1.075 \AA$.

A final structure factor calculation was computed with the $F_{c}{ }^{\prime \prime} s$ based on the parameters of all seventeen atoms of the asymmetric unit. In this calculation the hydrogen atoms, (located from the difference synthesis), were allocated the anisotropic thermal motion of the carbon atoms to which they are bonded, and the hydrogen scattering curve used was that given by Mclieeny (52). The final values of $R_{1}$ and $R_{2}$ were 0.124 and
(52) R. McWeeny, ibid., 4, 513 (1951).


The final fractional coordinates and standard deviations of the non-hydrogen atoms.

Table 6.2

|  | $B_{11}$ | $\sigma B_{11}$ | $B_{22}$ | $\sigma B_{22}$ | $B_{33}$ | $\sigma B_{33}$ | $B_{12}$ | $\sigma B_{12}$ | $B_{13}$ | $\sigma B_{13}$ | $B_{23}$ | $\sigma B_{23}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| S | 2.82 | 0.06 | 3.83 | 0.07 | 4.20 | 0.07 | -0.39 | 0.05 | 0.55 | 0.04 | -0.53 | 0.05 |
| 0 | 3.38 | 0.17 | 5.73 | 0.28 | 5.81 | 0.22 | -1.45 | 0.17 | 1.11 | 0.14 | -1.49 | 0.20 |
| $C(1)$ | 2.41 | 0.18 | 3.86 | 0.31 | 4.05 | 0.23 | 0.02 | 0.18 | 0.24 | 0.15 | 0.12 | 0.21 |
| $C(2)$ | 3.05 | 0.20 | 3.52 | 0.31 | 4.68 | 0.25 | 0.07 | 0.19 | 0.18 | 0.17 | -0.11 | 0.22 |
| $C(3)$ | 3.04 | 0.20 | 3.65 | 0.31 | 4.16 | 0.24 | 0.06 | 0.19 | 0.21 | 0.16 | -0.18 | 0.21 |
| $C(4)$ | 3.15 | 0.20 | 2.55 | 0.27 | 4.33 | 0.23 | 0.33 | 0.18 | 0.19 | 0.16 | 0.22 | 0.20 |
| $C(5)$ | 4.68 | 0.27 | 3.40 | 0.33 | 4.77 | 0.27 | 0.23 | 0.23 | 0.74 | 0.21 | -0.30 | 0.24 |
| $C(6)$ | 6.02 | 0.34 | 3.89 | 0.36 | 4.71 | 0.28 | 0.04 | 0.28 | 0.43 | 0.24 | -0.28 | 0.26 |
| $C(7)$ | 5.13 | 0.31 | 3.55 | 0.35 | 5.58 | 0.32 | 0.54 | 0.24 | -0.85 | 0.25 | -0.28 | 0.25 |
| $C(8)$ | 4.37 | 0.27 | 3.66 | 0.35 | 6.49 | 0.36 | 0.14 | 0.23 | -0.28 | 0.25 | -0.13 | 0.27 |
| $C(9)$ | 4.08 | 0.24 | 3.03 | 0.31 | 5.31 | 0.28 | -0.41 | 0.21 | 0.41 | 0.20 | -0.02 | 0.24 |

The final anisotropic thermal parameters, ( $\AA^{\circ}$ ), and their standard deviations.
Table 6.3

|  | Attached to <br> Carbon | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{C}$ | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{H}(1)$ | $\mathrm{C}(2)$ | 0.372 | 0.421 | 0.055 | 0.3793 | 0.4184 | 0.0561 |
| $\mathrm{H}(2)$ | $\mathrm{C}(5)$ | -0.085 | 0.529 | 0.175 | -0.1277 | 0.5546 | 0.1762 |
| $\mathrm{H}(3)$ | $\mathrm{C}(6)$ | -0.058 | 0.914 | 0.221 | 0.0259 | 0.8890 | 0.2372 |
| $\mathrm{H}(4)$ | $\mathrm{C}(7)$ | 0.404 | 1.150 | 0.227 | 0.3963 | 1.1616 | 0.2289 |
| $\mathrm{H}(5)$ | $\mathrm{C}(8)$ | 0.614 | 1.081 | 0.159 | 0.6220 | 1.0925 | 0.1589 |
| $\mathrm{H}(6)$ | $\mathrm{C}(9)$ | 0.487 | 0.783 | 0.098 | 0.4683 | 0.7496 | 0.0972 |

The hydrogen positions obtained from the difference synthesis.

The theoretical hydrogen positions calculated from the carbon locations.

Table 6.4
0.110 respectively. The observed and calculated structure factors are tabulated in Appendix I. Superimposed sections of a final three-dimensional electron density distribution, taken through the atomic centres parallel to (ONO), are shown in Figure 6.1, together with a drawing of the corresponding molecuar unit.


Composite electron density normal to [010]. Contour intervals at the carbon and oxygen atoms $1.0 \mathrm{e}^{-3}$, at the sulphur atom, $2.0 \mathrm{e}^{\AA^{-3}}$, (the zero contour is omitted).

## EXPERIMENTAL - PART II

THE X-RAY STRUCTURE ANALYSIS OF 3-PHENYL-1-PROPENE-1,3-DIONE 1-(DINETHYL MERCAPTOLE).

## CHAPTER 7.

## PRELIMINARY X-RAY WORK AND INTENSITY DATA COLLECTION.

3-phenyl-1-propene-1,3-dione 1-(dimethyl mercaptole) ${ }^{\ddagger}$, prepared by the method of Kelber (6), crystallises from an ethyl acetate and pentane mixture as yellow needles (53). The crystals were not sufficiently well formed to study the morphology, except to note that the needle axis is parallel to $\underline{b}$.

## X-ray equipment.

Preliminary photographic work was carried out using a Philips 1009 X-ray generator and a Stoe integrating Weissenberg goniometer. Three-dimensional intensity data were collected using a Picker automated four-circle diffractometer. Nickel filtered copper radiation was used throughout the structure analysis.

## Crystal density.

The density of the crystals, determined by flotation in a hexane and carbon tetrachloride mixture at room temperature, was found to be $1.26_{5} \mathrm{g.cc} .{ }^{-1}$. For PDN, therefore, the linear absorption coefficient, $\mu$, is $36.01 \mathrm{~cm}^{-1}$.

Unit cell dimensions and space group.
Cell dimensions were originally estimated from zero level Weissenberg photographs and were later refined using
measurements from the diffractometer. A least squares refinement procedure was used on observed theta angles for forty-two reflections.

$$
\begin{array}{ll}
\text { System monoclinic }(y \text { axis unique }) \\
a=5.972 \AA & (\sigma=0.007) \\
b=10.382 \AA & (\sigma=0.025) \\
c=19.536 \AA & (\sigma=0.011) \\
\beta=109.420 & (\sigma=0.02) \\
V=1142.3 \AA^{3} &
\end{array}
$$

The following systematic absences were found on zero level and equi-inclination Weissenberg photographs,

$$
\begin{aligned}
& \text { hkl - no general absences, } \\
& \text { hO1 - absent for } 1=2 N+1, \\
& \text { OkO - absent for } k=2 N+1
\end{aligned}
$$

The space group is therefore identified as $\mathrm{P}_{1} / \mathrm{c}$.

The number of molecules per unit cell.
Using deas $=1.26_{5} \mathrm{~g} \cdot \mathrm{cc}{ }^{-1}, V=2142.3 \AA^{3}$, and molecular weight, M , of 224.3,

$$
\underline{n}=\frac{d_{\text {meas }} \cdot v}{1.66 \cdot M}=3.89 .
$$

Thus there are four molecules of PDM per unit cell.

## The Picker four circle diffractometer.

A diagramatic representation of the four angles of the diffractometer is shown in Figure 7.1, which also defines the senses of $2 \theta, \chi$ and $\omega$. The zero position for $\varnothing$ is chosen


Schematic representation of a four-circle diffractometer.
Figure 7.1.
arbitrarily and the figure shows the direction of rotation which increases this angle. The diffractometer with all angles set to zero is shown schematically in Figure 7.2, which also defines the laboratory orthogonal axes system employed.

The intensity measurement for each reflection is carried out by moving the crystal through the Bragg reflecting position by motion about the $\omega$-axis. The detector rotates about the $2 \theta-$ axis, (which is coincident with the w-axis), its motion being restricted to the equatorial plane. This diffraction geometry is the commonly used normal-beam equatorial set-up.

The symmetrical-A setting, $(\omega=\theta)$, was used in all data collection. How the angular settings bring a reciprocal lattice point to the reflecting position on the surface of the Ewald sphere is shown in Figure 7.3. Figure 7.3, shows the Ewald sphere and a smaller sphere drawn through the reciprocal lattice point $P$ with its centre at the origin of the reciprocal lattice. $A \emptyset$ rotation brings $P$ to position $Q$ in the $x$-plane, and a $x$ rotation moves it to position $R$ on the $Y$ orthogonal axis. Figure 7.32 is a plan view on to the equatorial plane. The reciprocal lattice point, now at $R$, is brought to the surface of the Ewald sphere by rotating $\omega$ through the angle $\theta$, the detector being positioned at $2 \theta$ to receive the diffracted beam. With the Picker instrument in the symmetrical-A type mode, the $\omega$ and $2 \theta$ movements are linked in a $1: 2$ ratio, and the positioning of the crystal and detector for measurement of intensity


The diffractometer with all angles set to zero, showing the laboratory orthogonal axes system employed.

Figure 7.2.


Figure 7.31: Normal-beam equatorial geometry.


Figure 7.32.: The symmetrical-A setting.
The sequence of angular movements which bring the reciprocal lattice point $P$ to the reflecting position.
datum for each reflection is achieved by three angular settings, $\chi, \varnothing$ and $2 \theta$.

The detector used for intensity measurement was a Picker $2811 B$ scintillation counter, consisting of a thallium activated sodium iodide crystal and a photomultiplier tube, The diffracted X-radiation was filtered by a nickel slide, 10.001 inches thick), placed in front of the detector window. To guarantee that the scintillation counter was being used in the range where quanta entering and counts recorded have a linear relationship, a threshold point of 10,000 counts per second was used with the detector. If the count rate for any diffracted beam was greater than 10,000 c.p.s. attenuators were automatically introduced into the beam to reduce the count rate to below the threshold level. These attenuators were prepared using aluminium foil of different thicknesses. The attenuator factor for each was determined by positioning the detector to receive a moderately strong diffracted beam from the crystal, and counting for a fixed time period with and without the attenuator in position, (the crystal and detector remaining stationary). The attenuator factors so determined are given in the following Table, (CuKa).

| Attenuator | Attenuator fact |
| :---: | ---: |
| 1 | 1.35228 |
| 2 | 1.82866 |
| 3 | 2.47286 |
| 4 | 3.34400 |
| 5 | 4.52203 |

Collection of intensity data.
All intensity data were collected on the diffractometer employing the $\omega_{-2 \theta}$ scan mode, (moving-crystal, moving-counter). A variable scan range was used during the collection to allow for the increasing angular separation of the $K \alpha_{1}$ and $K \alpha_{2}$ components of the diffracted $X$-ray spectrum. The scan range, $\Delta 2 \theta$, for each reflection was calculated using the formulation of Alexander and Smith (54),

$$
\Delta 2 \theta=1.80+0.86 \tan \theta_{\mathrm{m}},(\mathrm{CuK} \mathrm{\alpha}),
$$

where $\theta_{m}$ is computed using the weighted mean wavelength of the $\underline{K} \alpha$ radiation. Thus the scan limits for each measurement are,

$$
\begin{aligned}
& 2 \theta_{1}=2 \theta_{m}-\Delta 2 \theta / 2 \\
& 2 \theta_{2}=2 \theta_{m}+\Delta 2 \theta / 2
\end{aligned}
$$

The scan rate used was one degree per minute and background counts were taken at both $2 \theta_{1}$ and $2 \theta_{2}$ for twenty seconds. The maximum $\omega$ angle attainable with the diffractometer is about $67^{\circ}$, above which the $x$-circle enters into collision with the X -ray tube housing; consequently the intensity data collection was limited to reflections with $2 \theta_{m}(C u K a) \leqslant 130^{\circ}$, (which allows about $4^{0}$ for scan completion).

The dimensions of the crystal used for intensity data collection were 0.3 mm in length, (parallel to b ), and $0.1 \times 0.1 \mathrm{~mm}$ in cross section. The crystal was mounted with b* parallel to the
(54) L. E. Alexander and G. S. Smith, Acta Cryst., 17, 1195 (1964).
$\emptyset$-axis of the goniostat, and $c^{*}$ was labelled $\emptyset=0$. Refined cell dimensions were obtained, (p.46), from which angular settings $\chi, \phi, 2 \theta_{1}$ and $2 \theta_{2}$ were calculated (55) for all independent reflections, ( hkl and $h \mathrm{kl}$ ), within the range $0<2 \theta_{\mathrm{m}}(\mathrm{CuK} \mathrm{\alpha}) \leqslant 130^{\circ}$. The digital input-output unit to the diffractometer is an I.B.M. 026 printing card punch, the input card for each reflection being pre-punched with indices and the required setting angles. The following output card, (initially blank), contains a duplication of this information and also records the attenuator number used, the total count over the scan range, (PK), and the twenty second background counts at $2 \theta_{1}$ and $2 \theta_{2},\left(B_{1}\right.$ and $B_{2}$ respectively). The input cards were sorted by hand to keep the time required for angular settings to a minimum. The $h k l$ and $h k \bar{l}$ series were run separately as the two major groups, and within each group data were collected in hnl levels. The setting cards in each hnl level were further sorted so that the 1 index changed more rapidly than $h$ and such that the sum of the differences between the indices of sucessive planes was a minimum. The average time taken for the complete angular settings and count measurements for each reflection was four minutes, of which twenty seconds, on average, represented the time taken to set the required angles.

## Data reduction.

With a twenty second background count at each end of the scan, the average background signal is $\left(B_{1}+B_{2}\right) / 40$ counts per
(55) We are grateful to Dr. J. Trotter, University of British Columbia, for the program used to calculate setting angles.
second. By assuming that this background count rate is characteristic for the total scan, the integrated intensity, $P K$, is background corrected and attenuator scaled by,

$$
I=\left[P K-\frac{60 \cdot \Delta 2 \theta \cdot\left(B_{1}+B_{2}\right)}{40}\right] \cdot A_{t},
$$

(where $60 . \Delta 2 \theta$ represents the scan time, in seconds). As all intensity data are measured in the equatorial plane, the Lorentzpolarisation factor applicable is,

$$
L p=\frac{1+\cos ^{2} 2 \theta}{2 \sin 2 \theta}
$$

No absorption corrections were applied, thus,

$$
\left|F_{h k l}\right|^{2}=I_{h k l} \cdot \frac{1}{L p_{h k l}}
$$

The standard deviation of each observed structure amplitude, $\sigma F_{0}$, was calculated by employing the statistical treatment formulated by Abrahams et al. $(56,57,58)$. With background measurements taken for a total of forty seconds and a scan rate of one degree per minute, then,

$$
\sigma F_{h k l}=\frac{1}{2\left|F_{h k l}\right|} \cdot \sqrt{\sigma^{2}\left|F_{h k l}\right|^{2}}
$$

where,

[^12]$$
\sigma^{2}\left|F_{h k l}\right|^{2}=\left[\frac{1}{L_{p h k}}\right]\left[P K+\left[\frac{\Delta 2 \theta, 60}{40}\right]^{2} \cdot\left(B_{1}+B_{2}\right)\right] \cdot A_{t}+c_{o r}\left|F_{h k l}\right|^{4}
$$
$\mathrm{C}_{\text {or }}$ incorporates many instumental features, such as the short time X -ray supply and counter circuit instability, and in this work was assigned the value 0.002 , (this being the average of the values found by Abrahams).

Intensity data were reduced, in the manner described above, to observed structure amplitudes, each with a corresponding standard deviation, employing a program written for the I.B.M. 7094 (59). If the peak corrected for background, (I), was zero or negative, the reflection was allocated a background corrected intensity value of two counts, reduced in the normal manner, and termed unobserved. In cases where the observed structure amplitude was finite, but had $\sigma F_{0} \geqslant\left|F_{0}\right|$, the calculated values of $\sigma F_{0}$ and $\left|F_{0}\right|$ were retained, but the reflection was again termed unobserved.

Intensity measurements were made on 1882 independent reflections within the range $0<2 \theta_{\mathrm{m}}(\mathrm{CuKa}) \leqslant 133^{\circ}$. After data reduction 322 of these reflections, $(17 \%)$, were classified as unobserved. Throughout the course of the work two residual index values are used. $R_{1}$ is the residual index obtained using all 1882 reflections for which intensity measurements were made,
(59) The data reduction progran was written by Mr. K. S. Dichmann, Chemistry Department, University of Toronto.
reflections classified as unobserved being treated as described above; whilst $R_{2}$ is the index obtained when reflections termed unobserved are excluded from the summation.

## CHAPTER 8.

## THE GENERATION OF NORMALISED STRUCTURE AMPLITUDES.

Generation of the $K$-curve.
The observed structure amplitudes were placed on the absolute scale using the K-curve method described by Karle and Hauptman (60) and later by Norment (61).

The magnitude of the normalised structure factor, $\left|E_{h k l}\right|$, is defined by equation (2.6), which can be extended to the form,

$$
\begin{equation*}
\left|E_{h k l}\right|=\left[\frac{\left|F_{h k l}\right|^{2} \cdot K\left(s_{h k l}\right)}{\varepsilon_{h k l} \sum_{i=1}^{N} f_{i}^{2} h k l}\right]^{1 / 2} \tag{8.1}
\end{equation*}
$$

where, $\varepsilon_{\mathrm{hkl}}$ is a quantity which depends upon the systematic absences of the space group of the crystal; for $P 2_{1} / c \in$ has the value two for hOl and OkO sets, and is otherwise unity,
$f_{i} h k l$ is the scattering factor of the ith atom in a unit cell containing a total of N atoms,
$\left|F_{h k l}\right|$ is the observed structure amplitude,
$K\left(S_{h k l}\right)$ is a function of $\sin \theta$ which places each $\left|F_{h k l}\right|^{2}$ datum on the absolute scale.
(60) J. Karle and H. Hauptman, Acta Cryst., 6, 473 (1953).
(61) H. G. Norment, U. S. Naval Research Laboratory Report 5739, Washington, D. C., (1961).

The theory of the method depends on the approximation $\langle E \mid 2\rangle=1$. Initially the reciprocal lattice, $\left(0<2 \theta_{\mathrm{m}} \leqslant 130^{\circ}\right)$, was divided into ten ranges of equal volume so that each sub-interval contained approximately the same number of reflections. Each of the ten ranges was characterised by a minimum and maximum $\sin \theta,\left(S_{1}\right.$ and $S_{2}$ respectively), and by a mean $\sin \theta, S_{m}$, where,

$$
s_{m}=\left[\frac{s_{1}^{3}+s_{2}^{3}}{2}\right]^{\frac{4}{3}}
$$

The value of $K\left(S_{m}\right)$ was calculated for each range, where,

$$
\begin{equation*}
K\left(S_{m}\right)=\sum_{s_{1}}^{s_{2} \varepsilon_{h k l} \sum_{i=1}^{N} f_{i}^{2} h k I} \frac{\sum_{s_{1}}\left|F_{h k l}\right|^{2}}{} \text {. } \tag{8.2}
\end{equation*}
$$

The procedure was repeated using fifteen and twenty ranges. The calculated values of $K\left(S_{m}\right)$ and $S_{m}$ are listed in Table 8.1, and the K-curve plot of $K\left(S_{m}\right)$ against $S_{m}$ is given in Figure 8.1 , which also shows a monotonically increasing curve drawn through the experimental points.

It is assumed that the curve through the experimental points can be formulated by an exponential function of the type,

$$
\begin{equation*}
K\left(S_{m}\right)=A \exp B\left(S_{m}\right)^{C} \tag{8.3}
\end{equation*}
$$

The 'best fit' values of $A$ and $B$ were calculated by a least squares method for the seven equations of type (8.3) which had the exponent C specified from 1.0 to 4.0 in intervals of 0.5 . Also computed for each equation was the sum of the squares of the deviations of the forty-five experimental points from the formulated curve,

| 10 Ranges |  | 15 Ranges |  | 20 Ranges |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $S_{m}$ | $K\left(S_{m}\right)$ | $S_{\text {rn }}$ | $K\left(S_{m}\right)$ | $S_{m}$ | $K\left(S_{m}\right)$ |
| 0.3316 | 1.0475 | 0.2896 | 2.0433 | 0.2632 | 0.9543 |
| 0.4782 | 1.5979 | 0.4177 | 1.2318 | 0.3795 | 1.2713 |
| 0.5670 | 2.1831 | 0.4953 | 1.6330 | 0.4500 | 1.5552 |
| 0.6343 | 2.1418 | 0.5541 | 2.2017 | 0.5034 | 1.6500 |
| 0.6897 | 3.3541 | 0.6025 | 2.0707 | 0.5474 | 2.2824 |
| 0.7374 | 4.7696 | 0.6642 | 2.2253 | 0.5853 | 2.0813 |
| 0.7796 | 6.0392 | 0.6811 | 3.4261 | 0.6188 | 1.7930 |
| 0.8177 | 5.4467 | 0.7143 | 3.8616 | 0.6490 | 2.5529 |
| 0.8525 | 6.1488 | 0.7448 | 4.7001 | 0.6767 | 3.2657 |
| 0.8847 | 6.5652 | 0.7729 | 6.0601 | 0.7022 | 3.4617 |
|  |  | 0.7991 | 4.9021 | 0.7260 | 4.9055 |
|  |  | 0.8237 | 6.7563 | 0.7484 | 4.6289 |
|  |  | 0.8469 | 6.2195 | 0.7695 | 5.3104 |
|  | - | 0.8689 | 6.3475 | 0.7895 | 6.8578 |
|  |  | 0.8899 | 6.4606 | 0.8085 | 4.4094 |
|  |  |  |  | 0.8267 | 7.4067 |
|  |  |  |  | 0.8441 | 6.6556 |
|  |  |  |  | 0.8608 | 5.6656 |
|  |  |  |  | 0.8769 | 7.5611 |
|  |  |  |  | 0.8924 | 5.7065 |

The computed values of $K\left(S_{m}\right)$ and $S_{m}$.

> Table 8.1.


The K-curve plot.

Table 8.1

SUM. The coefficients of the seven curves and the corresponding values of SUM are listed in Table 8.2. A minimum occurs for SUM between the two curves formulated with $C=1.5$ and $C=2.0$, so that the exponential functions derived with these coefficients are most representative of the experimental point distribution.

The K-curve method is essentially a modified Wilson function ${ }^{\neq}$, if $c$ is specified as two, then the relationships between the scaling quantities determined by this method and those of Wilson's original treatment are,

$$
\begin{aligned}
& A=K^{\prime} \\
& B=\frac{2 B_{O V}}{\lambda^{2}}
\end{aligned}
$$

where $K^{\prime}$ is the absolute scaling factor, (with respect to $\left|F_{0}\right|^{2}$ ), and $B_{o v}$ is the overall temperature factor of Wilson's method.

## The generation of normalised structure amplitudes.

The magnitudes of the normalised structure factors, $\left|E_{n k i}\right| s$ were computed using the expanded form of equation ( 8.1 ),

$$
\begin{equation*}
\left|E_{h k I}\right|=\left[\frac{\left|F_{h k l}\right|^{2} \cdot A \operatorname{expB}\left(\sin \theta_{h k I}\right)^{c}}{\varepsilon_{h k I} \sum_{i=1}^{N} f_{i}^{2} h k I}\right]^{1 / 2} \tag{8.4}
\end{equation*}
$$

The program used to compute the normalised structure factor magnitudes was also coded to calculate the statistical averages, $\left.\left.\langle | E\right|^{2}\right\rangle,\langle ||E|^{2}-1| \rangle$ and $\langle | E\rangle$; and the percentage distribution of $|E|$.

[^13]| $\underline{A}$ | $\underline{B}$ | $\underline{C}$ (specified) | SUM |
| :---: | :---: | :--- | ---: |
| 0.3058 | 3.5355 | 1.0 | 3.128 |
| 0.5918 | 3.0369 | 1.5 | 2.777 |
| 0.8320 | 2.8586 | 2.0 | 2.991 |
| 1.0295 | 2.8081 | 2.5 | 3.729 |
| 1.1947 | 2.8221 | 3.0 | 5.155 |
| 1.3356 | 2.8741 | 3.5 | 7.635 |
| 1.4580 | 2.9514 | 4.0 | 11.796 |

The 'best fit' coefficients $A$ and $B$ computed for the seven curves. SUM is the sum of the squares of the deviation of the forty-five experimental points from the formulated curve.

Table 8.2

Normalised structure factor magnitudes were computed using the two 'best-fitting' K-curves, $(C=1.5$ and $C=2.0$, Table 8.2). The statistical averages and the $|\mathrm{E}|$ distribution produced are shown in Table 8.3, together with the theoretical values of these quantities for both acentric and centric space groups. From Table 8.3 it can be seen that the averages and |El distribution are almost identical for the two K-curve formulations and a detailed inspection of the complete $|E|$ listings showed that individual $|E|$ 's were, in general, insignificantly different in the two sets; the only major differences apparent occur at low $\sin \theta$ when $K\left(S_{m}\right)$ approximates to $A$. The symbolic addition procedure was carried out using normalised structure factor magnitudes generated with the second $K$-curve, ( $C=2.0$ ), as this made the thermal correction factor analogous with the normally assumed Debye form.

For application of the symbolic addition procedure, the normalised structure factor magnitudes were listed in descending order and classified by their indices into the eight groups of type ggg, gug, ugg, ggu, uuu, uug, ugu and guu, (where g = even and $u=$ odd). There are 91 planes with $|E| \geqslant 2.0$ and 242 with $|E| \geqslant 1.5$. The fact that the distribution of the $|E| ' s$ within the eight parity groups is reasonably uniform, (shown in Table 8.4), indicates that there is no specialisation in the location of the sulphur atoms, (as had been found in the desaurin from acetophenone structure), and is a favourable indication for the success


A
B
C
$<|E|>$
0.835
$\langle | E^{2}-1| \rangle$
0.922
$\left\langle\mathrm{E}^{2}\right\rangle$
1.016
$\%|E| \geqslant 3.0$
0.2
$\%|E| \geqslant 2.0$
4.9
$\%|E| \geqslant 1.0$
32.1
0.8320
0.835
0.798
0.886
0.922
0.968
0.736
1.017
1.000
1.000
0.2
0.3
$0.0_{1}$
32.2

Theoretical

2.8586
2.0
4.8
5.0
2.8
32.0
37.0

The distribution of normalised structure factor magnitudes and statistical averages found together with the theoretical values.

$$
\text { Table } 8.3
$$

| h k 1 | $\|E\| \geqslant 2.0$ | $\|E\| \geqslant 1.5$ |
| :---: | :---: | :---: |
| gug | 9 | 40 |
| gu ${ }^{\text {a }}$ | 13 | 31 |
| g g g | 13 | 26 |
| $g \mathrm{~g} \mathrm{u}$ | 14 | 29 |
| u ug | 8 | 25 |
| u u u | 10 | 26 |
| ugg | 11 | 32 |
| ugu | 13 | 33 |
| $\left[\begin{array}{lll} g & u & g \\ g & g & u \\ u & u & g \\ u & g & u \end{array}\right]$ | 44 | 127 |
| $\left.\sum \begin{array}{lll} g & u & u \\ g & g & g \\ u & u & u \\ u & g & g \end{array} \right\rvert\,$ | 47 | 115 |

$\left.\sum \begin{array}{lll}g & u & g \\ g & g & u \\ u & u & g \\ u & g & u\end{array}\right]$
44
127
$\left.\sum \begin{array}{lll}g & u & u \\ g & g & g \\ u & u & u \\ u & g & g\end{array}\right]$
47 115

The distribution of normalised structure factor magnitudes throughout the eight sub-groups.

Table 8.4
of the symbolic addition procedure.
The program used to generate the K-curve and compute normalised structure factor magnitudes was written in Fortran IV for the I.B.M. 7094 ${ }^{*}$.
\# A listing of this program is given in Appendix V .

## CHAPTER 9.

## THE SYMBOLIC ADDITION PROCEDURE.

The practical aspects of the symbolic addition procedure follow the method outlined by Karle and Karle (62). The only phase relationship used was $\Sigma_{2}$, which for a centrosymmetrical crystal is formulated as,

$$
s E_{\vec{h}} \sim s \sum_{\vec{k}} E_{\vec{k}} \cdot E_{\vec{h}-\vec{k}} \text {, (where 's' means 'sign of'). }
$$

The hyperbolic tangent probability formula was used to evaluate the reliability of phases generated by application of the $\Sigma_{2}$ relationship $(34,35)$. The probability function, $\mathrm{P}+(\overrightarrow{\mathrm{h}})$, which determines the probability that the sign of $\mathrm{E}_{\mathrm{h}}$ be positive, is defined by,

$$
P+(\vec{h})=1 / 2+1 / 2 \tanh \left[\frac{\sigma_{3}}{\sigma_{2}^{3 / 2}}\left|E_{\vec{h}}\right| \cdot \sum_{\vec{k}} E_{\vec{k}} \cdot E_{\vec{h}-\vec{k}}\right],
$$

where $\sigma_{n}=\sum_{i=1}^{N} z_{i}^{n} ; z_{i}$ being the atomic number of the fth atom in a unit cell containing $N$ atoms. For $P D N \quad \sigma_{2}=4013$ and

$$
\sigma_{3}=46,284 .
$$

The $\Sigma_{2}$ listing and combinational count.
A $\sum_{2}$ listing (62) of all combinations of $\vec{k}$ and $\vec{h}-\vec{k}$ for a given $\vec{h}$ was computed, considering only the 91 reflections with $|E| \geqslant 2.0$, (which were expanded to $364|E \vec{h}|$ using symmetry related reflections). This listing facilitated a concise
(62) J. Karle and I. L. Karle, Acta Cryst., 21, 849 (1966).
means of applying the $\Sigma_{2}$ formula by hand for this group of reflections. Also computed were the number of times each of the reflections entered into a $\Sigma_{2}$ combination with other members of the set (63).

## Origin specification.

Working in three-dimensions in the monoclinic system, (with a primitive unit cell), the origin is fixed by assigning phases arbitrarily to three linearly independent reflections $(64,65)$. The three $E_{\vec{k}}$ selected to specify the origin were,

| $h$ | $k$ | 1 | $s$ | $\|E\|$ | number of <br> combinations. |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | 2 | -9 | + | 2.63 | 60 |
| 5 | 1 | -2 | + | 2.83 | 33 |
| 2 | 3 | 1 | + | 2.90 | 36 |

These three linearly independent reflections were chosen on the basis that they have large $|E|$ value, (hence any new phases generated by $\Sigma_{2}$ application will have a high associated probability), and that they form numerous $\Sigma_{2}$ combinations amongst themselves and with other reflections of larger $|E|$.

## The symbolic addition procedure.

$\Sigma_{2}$ generation with the three origin specifying reflections

[^14]yielded signs for five additional $\mathrm{E}_{\mathrm{h}}$ on the first cycle. Although each of the generated signs was the result of a single $\Sigma_{2}$ combination, the associated probability values were all greater than 0.99. These five normalised structure factors were added to the original three to form the input for the next cycle. After three cycles of $\Sigma_{2}$ application, in this manner, a total of thirteen reflections, lexclusive of those used to specify the origin), could be allocated an absolute phase, (+ or - ). If the generated sign was the result of a single $\Sigma_{2}$ combination, a probability minimum of 0.975 was imposed; single generations of lower probability were placed to one side in the hope that further $\Sigma_{2}$ combinations would be formed at a later stage, which would produce a corresponding increase in the probability factor. One- and two-dimensional, data were treated with great care during the procedure, as it has been previously found that exceptions to the $\Sigma_{2}$ relationship are much more numerous with these type of reflections.

Since no additional signs could be generated after these three cycles a symbolic assignment was necessary and by inspection of the $\Sigma_{2}$ listing 2.10 .2 was chosen and allocated the symbol A. This reflection was chosen on the basis that it would enter into many $\Sigma_{2}$ combinations with the sixteen reflections whose phases were already known, and that it had a large $\mid \mathrm{El}$ of 2.64. After nine additional cycles of $\Sigma_{2}$ all 91 reflections with $|E| \geqslant 2.0$ could be allocated phases of + , - or $\pm$ A. Since $A$ must be plus or minus, phases determined as $+A^{2}$ and $-A^{2}$ were carried through $a s+$ and - respectively. In the final three cycles of the
procedure it was noted that phases generated for a few of the reflections were of mixed absolute and symbol type, (either $+A$ and - or $-A$ and + ), indicating that the relationship $A=-$ was probable. Karle and Karle (62) point out that symbols are often found to be definable in terms of others or in terms of absolute values, but suggest that any relationships indicated should not be assumed until the completion of the procedure. The relationship that A was negative was therefore not accepted at this stage, and reflections for which generated phases were of mixed absolute and symbolic type were placed aside and not used in $\Sigma_{2}$ generation when the list was extended.

Now that phases could be allocated to all reflections with $|E| \geqslant 2.0$, the list was extended by including the 151 planes for which 2.0$\rangle|E| \geqslant 1.5$. The input for the extended cycle were $84 \mathrm{E}_{\vec{k}}$ for which phases had been generated in the first part of the procedure, (seven reflections being excluded because of low probability or because phases indicated were of a mixed absolute and symbol type). As the $|E|$ values in the new list were smaller, with the result that the triple product for a single $\Sigma_{2}$ combination and hence the probability were lower, a criterion was adopted that there must be at least five independent $\Sigma_{2}$ combinations for any new reflection before the indicated phase would be accepted. On the first pass phases could be assigned to $107 \mathrm{E}_{\overrightarrow{\mathrm{h}}}$ in the new list using this criterion (66). Numerous of the phases generated had both symbolic and
(66) This final $\mathrm{E}_{\text {a }}$ cycle was computed using a nrogram written
by Mr. F. T. C. Leung.
absolute contributors, indicating again that $A$ was negative, and this assignment was now accepted.

A summary of the $\sum 2$ cycles is given in Appendix II. The E map computation.

The 198 signs determined for reflections with larger $|E|$ were already in excess of the ten to twelve normalised structure factors per atom of the asymmetric unit generally accepted as being sufficient to produce a well resolved E-map $(37,38)$. Using the 198 normalised structure factors as Fourier coefficients, the three-dimensional E-map was computed for a quarter of the unit cell defined by $-0.5 \leqslant x \leqslant 0.5,0.0 \leqslant y \leqslant 0.5$ and $0.0 \leqslant z \leqslant 0.5$, in intervals of 0.05 in $x, 0.025$ in $y$ and 0.015 in 2. The resultant synthesis clearly revealed the location of all the atoms of the molecule, with the exception of the two carbon atoms of the methyl groups. Around the sulphur atoms there were subsiduary diffraction peaks, (in the region $1.7 \AA$ to $2.2 \AA$ distant from the sulphur atoms), similar to those found in the E-map of allylthiourea (67), and unambiguous identification of the remaining carbon atoms was not possible. The fractional coordinates of the twelve atoms whose locations could be identified from the E-map are listed in Table 9.1. A composite drawing of the E-map, formed by superimposing sections taken through the peak centres normal to $c^{*}$, is shown in Figure 9.1, where also is shown the atomic numbering system employed. The

[^15]| Atom | X/a | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ |
| :--- | :---: | :---: | :---: |
| $\mathrm{S}(1)$ | 0.530 | 0.160 | 0.487 |
| $\mathrm{~S}(2)$ | 0.728 | 0.412 | 0.465 |
| 0 | 0.073 | 0.153 | 0.393 |
| $C(1)$ | 0.475 | 0.310 | 0.438 |
| $C(2)$ | 0.280 | 0.333 | 0.383 |
| $C(3)$ | -0.075 | 0.258 | 0.360 |
| $C(4)$ | -0.110 | 0.310 | 0.250 |
| $C(5)$ | -0.320 | 0.438 | 0.197 |
| $C(6)$ | -0.535 | 0.383 | 0.184 |
| $C(7)$ | -0.540 | 0.265 | 0.275 |

The atomic coordinates estimated from the E-map.


Composite E-map normal to $\mathrm{c}^{*}$. The contours are spaced at equal arbitary intervals.
additional peaks around the sulphur atoms are shown in Figure 9.1 with the label $A$.

In the allylthiourea problem the subsiduary peaks close to the sulphur atom were removed by increasing the number of coefficients in the E-map computation. However we decided that rather than generate more phases using $\Sigma_{2}$, an electron density synthesis, using observed structure amplitudes phased by the twelve atoms whose positions were postulated, should unambiguously identify the carbon atoms of the methyl groups and at the same time confirm the atomic locations obtained from the E-map. Structure factors were computed using the atomic locations listed in Table 9.1 and allocating an overall temperature factor of $3.4 \AA^{2}$, (calculated from the B coefficient of the K-curve). Reflections tagged unobserved were excluded from the calculation which yielded a residual index of 0.38 . The subsequent electron density synthesis ${ }^{\ddagger}$ clearly revealed the positions of the carbon atoms of the methyl groups. The fractional coordinates of all fourteen atoms of the molecule were assessed from the synthesis, and are listed in Table 9.2. Refinement of the atomic parameters by the method of least squares was now initiated.

[^16]| Atom | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ |
| :--- | :--- | :--- | :--- |
| $\mathrm{S}(1)$ | 0.530 | $0.157_{5}$ | 0.483 |
| $\mathrm{~S}(2)$ | 0.732 | 0.405 | 0.468 |
| 0 | 0.065 | $0.157_{5}$ | 0.390 |
| $\mathrm{C}(1)$ | 0.480 | 0.304 | 0.438 |
| $\mathrm{c}(2)$ | 0.286 | 0.334 | $0.382_{5}$ |
| $\mathrm{c}(3)$ | 0.070 | 0.261 | $0.358_{5}$ |
| $\mathrm{c}(4)$ | -0.132 | $0.297_{5}$ | $0.298_{5}$ |
| $\mathrm{C}(5)$ | -0.117 | $0.398_{5}$ | 0.253 |
| $\mathrm{C}(6)$ | -0.3 .13 | 0.446 | $0.196_{5}$ |
| $\mathrm{C}(7)$ | -0.533 | $0.382_{5}$ | 0.182 |
| $\mathrm{C}(8)$ | -0.555 | $0.277_{5}$ | 0.228 |
| $\mathrm{C}(9)$ | -0.356 | 0.234 | 0.287 |
| $\mathrm{Me}(10)$ | 0.835 | 0.149 | 0.5375 |
| $\mathrm{Me}(11)$ | 0.630 | 0.550 | 0.401 |

Atomic coordinates estimated from the three-dimensional electron density synthosis.

## CHAPTER 10.

## REFINEMENT OF THE STRUCTURE.

The least squares program used was the same version of OR-FLS which had previously been employed in the refinement of the desaurin from acetophenone structure. All recorded reflection data were used in the refinement except in the final two cycles with anisotropic temperature factors when reflections classified as unobserved were excluded. Each observed structure amplitude was weighted using its standard deviation, $\sigma F_{0}$, calculated in the data reduction procedure,


An overall scale factor was treated as a variable throughout the refinement; the function minimised was $\sum w\left(\left|F_{0}\right|-\left|F_{c}\right|\right)^{2}$. Refinement with isotropic temperature factors.

The initial positional parameters used in the refinement were those obtained from the electron density synthesis, (Table 9.2), and each atom was allocated an isotropic temperature factor of $3.4 \AA^{2}$. One cycle of refinement was computed with the thermal parameters of the atoms fixed and was followed by three cycles in which all fifty-six positional and thermal parameters were refined, $R_{1}$ and $R_{2}$ improving to 0.18 and 0.16 respectively. The maximum parameter shift calculated in the last cycle represented about a quarter of the standard deviation of the parameter and
conversion of the thermal parameters into anisotropic mode was now considered in order. The atomic parameters at this stage are listed in Table 10.1.

## Refinement with anisotropic temperature factors.

Four cycles of least squares refinement, performed on the one hundred twenty-seven variables using all recorded reflection data, proceeded normally. A further two cycles, computed with reflections tagged unobserved excluded, produced only small parameter shifts and illustrated that the weighting of these reflections was such that their inclusion or exclusion had little effect during the least squares refinement. The shifts calculated in the final cycle were at most $25 \%$ of the standard deviation of the parameter, and refinement was considered complete. The residual index $R_{2}$ was 0.112 . The final positional and thermal parameters of the fourteen non-hydrogen atoms of the molecule, and their estimated standard deviations, calculated from the least squares residuals, are listed in Tables 10.2 and 10.3. $B_{i j}$, in Table 10.3, are coefficients in the anisotropic temperature expression,

$$
\begin{aligned}
& \exp -1 / 4\left(h^{2} a^{* 2} B_{11}+k^{2} b^{* 2} B_{22}+1^{2} c^{* 2} B_{33}\right. \\
& \left.\quad+2 h k a^{*} b^{*} B_{12}+2 h l a^{*} c^{*} B_{13}+2 k l b^{*} c * B_{23}\right)
\end{aligned}
$$

Location of the hydrogen atoms.
The final atomic parameters, (Tables 10.2 and 10.3), were used to set-up coefficients for a three-dimensional difference synthesis, from which peaks attributable to the twelve hydrogen

| Atom | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ | $\mathrm{B}_{\mathrm{iso}}\left(\mathrm{A}^{2}\right)$ |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{S}(1)$ | 0.5257 | 0.1589 | 0.4812 | 3.613 |
| $\mathrm{~S}(2)$ | 0.7298 | 0.4106 | 0.4651 | 3.906 |
| 0 | 0.0653 | 0.1523 | 0.3892 | 4.114 |
| $\mathrm{C}(1)$ | 0.4831 | 0.3100 | 0.4377 | 3.167 |
| $\mathrm{C}(2)$ | 0.2740 | 0.3428 | 0.3845 | 3.417 |
| $\mathrm{C}(3)$ | 0.0733 | 0.2600 | 0.3603 | 3.257 |
| $\mathrm{C}(4)$ | -0.1406 | 0.3020 | 0.2969 | 3.486 |
| $\mathrm{C}(5)$ | -0.1225 | 0.3998 | 0.2512 | 4.359 |
| $\mathrm{C}(6)$ | -0.3287 | 0.4342 | 0.1924 | 4.328 |
| $\mathrm{C}(7)$ | -0.5451 | 0.3724 | 0.1846 | 4.371 |
| $\mathrm{C}(8)$ | -0.5548 | 0.2730 | 0.2304 | 4.378 |
| $\mathrm{C}(9)$ | -0.3560 | 0.2355 | 0.2879 | 3.487 |
| $\mathrm{Me}(10)$ | 0.8455 | 0.1558 | 0.5351 | 4.470 |
| $\mathrm{Me}(11)$ | 0.6412 | 0.5507 | 0.4042 | 4.131 |

Atomic parameters at the completion of the isotropic least squares refinement.

## Table 10.1

| Atom | $\sigma x / a$ |  | $\sigma \mathrm{F} / \mathrm{b}$ |  |  | $\sigma Z / c$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | X/a | $\times 10^{4}$ | $\mathrm{Y} / \mathrm{b}$ | $\times 10^{4}$ | 2/c | $\times 10^{5}$ |
| S(1) | 0.5264 | 5 | 0.1585 | 3 | 0.4812 | 15 |
| S(2) | 0.7300 | 5 | 0.4107 | 3 | 0.4649 | 15 |
| 0 | 0.0684 | 13 | 0.1511 | 7 | 0.3895 | 38 |
| c(1) | 0.4843 | 17 | 0.3079 | 9 | 0.4378 | 55 |
| C(2) | 0.2739 | 18 | 0.3433 | 10 | 0.3857 | 58 |
| c(3) | 0.0729 | 18 | 0.2580 | 9 | 0.3610 | 56 |
| c(4) | -0.1386 | 19 | 0.3024 | 10 | 0.2982 | 55 |
| C(5) | -0.1223 | 20 | 0.3986 | 11 | 0.2509 | 59 |
| c(6) | -0.3262 | 21 | 0.4356 | 11 | 0.1934 | 60 |
| c(7) | -0.5435 | 22 | 0.3732 | 13 | 0.1840 | 62 |
| c(8) | -0.5569 | 21 | 0.2729 | 11 | 0.2312 | 63. |
| c(9) | -0.3546 | 19 | 0.2363 | 11 | 0.2874 | 61 |
| $\mathrm{Me}(10)$ | 0.8446 | 19 | 0.1558 | 12 | 0.5355 | 64 |
| Me(11) | 0.6406 | 20 | 0.5491 | 10 | 0.4045 | 60 |

The final fractional coordinates and standard deviations of the non-hydrogen atoms.

Table 10.2


| $s(1)$ | 3.34 | 0.15 | 3.03 | 0.12 | 4.03 | 0.14 | 0.06 | 0.11 | 0.15 | 0.11 | 0.96 | 0.11 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $s(2)$ | 3.53 | 0.15 | 3.28 | 0.13 | 4.36 | 0.15 | -0.56 | 0.12 | 0.07 | 0.11 | 0.64 | 0.11 |
| 0 | 3.56 | 0.39 | 2.84 | 0.33 | 5.24 | 0.39 | 0.15 | 0.30 | -0.13 | 0.32 | 2.04 | 0.31 |
| $c(1)$ | 2.49 | 0.47 | 2.67 | 0.46 | 2.97 | 0.49 | -0.62 | 0.38 | 0.24 | 0.40 | -0.51 | 0.39 |
| $c(2)$ | 3.29 | 0.51 | 2.60 | 0.44 | 4.37 | 0.57 | 0.32 | 0.42 | 1.15 | 0.45 | 0.44 | 0.45 | $\begin{array}{llllllllllllll}c(3) & 3.74 & 0.55 & 2.10 & 0.44 & 3.48 & 0.50 & 0.19 & 0.40 & 0.72 & 0.42 & -0.49 & 0.42\end{array}$ $\begin{array}{lllllllllllll}c(4) & 3.65 & 0.56 & 2.39 & 0.47 & 3.01 & 0.49 & 0.76 & 0.42 & 0.97 & 0.43 & 0.00 & 0.40\end{array}$ $\begin{array}{lllllllllllll}C(5) & 3.93 & 0.62 & 4.96 & 0.59 & 3.61 & 0.53 & 1.38 & 0.53 & 0.90 & 0.49 & 1.45 & 0.50\end{array}$ $\begin{array}{lllllllllllll}c(6) & 3.71 & 0.62 & 4.28 & 0.64 & 4.25 & 0.57 & 0.30 & 0.54 & 0.09 & 0.48 & -0.86 & 0.50\end{array}$ $c(7) . \begin{array}{llllllllllll}4.03 & 0.67 & 5.63 & 0.79 & 4.13 & 0.60 & 1.32 & 0.58 & 0.39 & 0.51 & -0.92 & 0.54\end{array}$ $\begin{array}{lllllllllllll}c(8) & 3.86 & 0.66 & 3.70 & 0.57 & 4.10 & 0.60 & -0.05 & 0.49 & -0.25 & 0.52 & -0.94 & 0.49\end{array}$ $\begin{array}{lllllllllllll}c & c & 2.86 & 0.54 & 4.03 & 0.59 & 4.50 & 0.57 & -1.48 & 0.47 & 1.11 & 0.47 & -2.79\end{array} 0.50$ $\begin{array}{lllllllllllll}\mathrm{Me}(10) & 2.75 & 0.53 & 3.10 & 0.60 & 6.02 & 0.72 & 0.44 & 0.49 & -1.66 & 0.50 & 0.44 & 0.57\end{array}$ $\begin{array}{lllllllllllll}\text { Me(11) } & 4.99 & 0.65 & 2.88 & 0.48 & 5.64 & 0.62 & 0.68 & 0.48 & 0.97 & 0.51 & 2.99 & 0.47\end{array}$

The final anisotropic thermal parameters, ( $\AA^{2}$ ), and their standard deviations.
atoms of the molecule were identified. Their fractional coordinates are listed in Table 10.4, together with the theoretical values calculated from the carbon atom positions assuming a carbon-hydrogen bond length of $1.075 \AA^{\circ}$. In this calculation one hydrogen atom of each methyl group was considered staggered with respect to the corresponding $S-C(1)$ direction. The numbering system used for the hydrogen atoms is shown in Figure 10.1.

Hydrogen atoms, located from the difference synthesis and assigned the anisotropic thermal motion of the carbon atom to which they are bonded, were included in a final structure factor calculation. The hydrogen atom scattering factor used was that listed in International Tables (43). The observed and calculated structure factors are tabulated in Appendix III, and yield $R_{1}$ and $R_{2}$ of 0.119 and 0.101 respectively. Superimposed sections of a final three-dimensional electron density distribution, taken through the atomic centres normal to $c^{*}$, are shown in Figure 10.2.

Analysis of the phases generated by the symbolic addition procedure.
The phases generated by the symbolic addition procedure for the 198 normalised structure factors which constituted the input of the E-map were compared with those calculated from the final atomic parameters. It was revealed that the phase of 254, $(|E|=1.58)$, was the only incorrect allocation made during the symbolic addition procedure. Since it cannot be considered that this single incorrect phase could have significantly affected the re sultant E-map, the additional diffraction maxima around the sulphur atoms apparent in the synthesis appear the result of a lack of coefficients in the computation.

| Atom | Attached to carbon | X/a | Y/b | 2/c | X/a | Y/b | Z/c |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| H(1) | c(2) | 0.245 | 0.422 | 0.351 | 0.2593 | 0.4382 | 0.3626 |
| H(2) | C(5) | 0.024 | 0.475 | 0.261 | 0.0451 | 0.4447 | 0.2581 |
| H(3) | c(6) | -0.390 | 0.483 | 0.149 | -0.3159 | 0.5111 | 0.1569 |
| H(4) | c(7) | -0.685 | 0.416 | 0.144 | -0.7001 | 0.4021 | 0.1406 |
| $\mathrm{H}(5)$ | C(8) | -0.720 | 0.219 | 0.230 | -0.7230 | 0.2250 | 0.2237 |
| H(6) | C(9) | -0.355 | 0.178 | 0.340 | -0.3622 | 0.1580 | 0.3226 |
| H(7) | $\mathrm{Me}(10)$ | 0.852 | 0.068 | 0.533 | 0.8706 | 0.0639 | 0.5619 |
| H(8) | Me(10) | 0.972 | 0.158 | 0.495 | 0.9520 | 0.1645 | 0.5011 |
| H(9) | Me (10) | 0.885 | 0.227 | 0.570 | 0.8874 | 0.2325 | 0.5747 |
| H(10) | Me(11) | 0.805 | 0.600 | 0.416 | 0.7918 | 0.6121 | 0.4211 |
| H(11) | Me(11) | 0.600 | 0.490 | 0.354 | 0.6012 | 0.5203 | 0.3489 |
| H(12) | Me(11) | 0.530 | 0.618 | 0.421 | 0.4878 | 0.5954 | 0.4109 |
|  |  | Hydrogen positions obtained from the difference synthesis. |  |  | Theoretical hydrogen positions calculated from the carbon locations. |  |  |



The atomic numbering system used.
Figure 10.1


Composite electron density normal to c*. Contour intervals at carbon and oxygen atoms $2.0 e^{-3}$, at the sulphur atoms $2.0 \mathrm{e}^{\AA^{-3}}$. The zero contour is omitted.

## EXPERIMENTAL - PART III

THE STRUCTURE ANALYSIS OF THE 1,2,4TRITHIACYCLOPENTANE FROM PINACOLONE.

## CHAPTER 11.

## PRELIMINARY X-RAY YORK AND INTENSITY DATA COLLECTION.

The 1,2,4-trithiacyclopentane from pinacolone, prepared by the method of Gompper (7), crystallises from a chloroform and hexane mixture as yellow needles (68). The crystals have a as the needle axis with (Oll) and (O1I) developed, (Figure 11.1, p.82).

## X-ray equipment.

Preliminary photographic work was carried out using a Philips 1009 X-ray generator and a Stoe integrating Weissenberg goniometer. Three-dimensional intensity data were collected using a Picker automated four-circle diffractometer. Nickel filtered copper radiation was used throughout the structure analysis. Unit cell dimensions.

Cell dimensions were originally obtained from zero level Weissenberg photographs and were later determined more accurately using measurements from the diffractometer. The crystal, (later used for intensity data collection), was mounted with $b^{*}$ parallel to the $\varnothing$ axis of the goniostat, and at $\mathcal{X}=0, \beta^{*}$ was determined directly as the $\varnothing$ angular displacement between $a^{*}$ and $c^{*}$. The axial lengths were calculated from average observed theta angles for two h 00 , three $0 k 0$ and four 001 reflections,
(68) We are grateful to Dr . P. Yates for supplying crystals of the $1,2,4$-trithiacyclopentane from pinacolone.

$$
\begin{aligned}
& \text { System monoclinic (y axis unique) } \\
& \mathrm{a}=5.93_{6} \AA\left( \pm 0.00_{7}\right)^{\ddagger} \\
& \mathrm{b}=13.02_{3} \AA\left( \pm 0.01_{0}\right) \\
& \mathrm{c}=22.732 \AA\left( \pm 0.01_{5}\right) \\
& \beta=108.0_{2} 0\left( \pm 0.0_{2}\right) \\
& \mathrm{v}=1757.3 \AA^{3}
\end{aligned}
$$

Because the cell dimensions were determined using relatively few measurements from the diffractometer, their accuracy were assessed by comparing observed and calculated values of $2 \theta$ and $\emptyset$ for twenty reflections of the hOl and Okl zones. The maximum deviation found was $0.02^{\circ}$.

## Crystal density.

The average density of the crystals, ( $\mathrm{d}_{\text {meas }}$ ), determined by flotation in an acetone and carbon tetrachloride mixture at room temperature, was found to be $1.22 \mathrm{~g} . \mathrm{cc}^{-1}$. Using this value for the crystal density the linear absorption coefficient, $\mu$, is $37.5 \mathrm{~cm}^{-1}$.

Space group and number of molecules per unit cell.
The systematic absences found on zero level and equiinclination Weissenberg photographs were,
hkl - no general absences,
hol - absent for $1=2 \mathrm{~N}+1$,
OkO - absent for $k=2 N+9$.

[^17]The space group is therefore identified as $\mathrm{P} 21 / \mathrm{c}$. Using $d_{\text {meas }}=1.22$ g.cc-1, $V=1757.3 \AA^{\circ}$ and molecular weight, M , $=316.5$,

$$
n=\frac{d_{\text {meas }} \cdot v}{1.66 \cdot M}=4.07 .
$$

Thus there are four molecules of the 1,2,4-trithiacyclopentane from pinacolone per unit cell.

Collection of three-dimensional intensity data.
An acicular crystal was cut into a more symmetrical shape and used for intensity data collection. The approximate dimensions of the crystal were,


Figure 11.1

Absorption is therefore fairly low, ( $\mu_{\mathrm{CuK}}=37.5 \mathrm{~cm}^{-1}$ ), and no corrections were applied.

Three dimensional intensity data were collected using a Picker automated four-circle diffractometer ${ }^{\boldsymbol{}}$. The crystal was mounted with $b^{*}$ parallel to the $\varnothing$ axis of the goniostat and $a^{*}$ was specified as $\emptyset=0$. All data were collected employing the $\omega-2 \theta$ scan mode, (moving-crystal, moving-counter), using a scan rate of one degree per minute. The scan range for each reflection

[^18]was calculated using the standard equations of Alexander and Smith, (equations 7.1), and background measurements were taken at both $2 \theta_{1}$ and $2 \theta_{2}$ for twenty seconds. As previously stated in Chapter 7, the maximum value of $\omega$ safely attainable with the Picker diffractometer is $67^{\circ}$; consequently the intensity data collection was restricted to reflections for which $2 \theta_{m}\left(\mathrm{CuK}_{\alpha}\right)$. $\leqslant 130^{\circ}$. The setting cards, which form the input to the diffractometer, were ordered by hand to minimise the time required for angular slew between measurements. The hkl and hkI series were run separately as the two major groups and within each data were collected in hnl levels, with the 1 index changing more rapidly than h. Crystal orientation and system stability were checked during the automatic phases of the collection by taking intensity measurements for twelve 'test' reflections after each hnl level was completed. The maximum deviation of a test measurement, from its mean value over the total collection, was found to be $2 \%$. The procedure whereby intensity data, measured using the diffractometer, were reduced to observed structure amplitudes was analogous to that employed in the structure analysis of $\mathrm{PDM}{ }^{\ddagger}$; the statistical treatment due to Abrahams being retained to calculate the standard deviation, $\sigma F_{0}$, of each $\left|F_{0}\right|$ datum. The classification and treatment of unobserved reflections followed the same lines as that employed in the PDPi data reduction, (page 53). Intensity measurements were made on 2778 independent

[^19]reflections within the range $0<2 \theta_{\mathrm{m}} \leqslant 130^{\circ}$; after data reduction 530 of these reflections, ( $19.1 \%$ ), were classified as unobserved. Absolute scaling and everall temperature factor.

The K-curve method ${ }^{\ddagger}$ was employed to place the observed structure amplitudes on the absolute scale and ascertain the average isotropic temperature factor.

Initially the reciprocal lattice, $\left(0<2 \theta_{\mathrm{m}} \leqslant 130^{\circ}\right)$, was divided into ten sub-intervals, (of equal volume), and the value of $K\left(S_{m}\right)$ calculated for each, $\left[K\left(S_{m}\right)\right.$ is defined by equation (8.2)]. The procedure was repeated using fifteen and twenty ranges. The calculated values of $K\left(S_{m}\right)$ and the corresponding values of $S_{m}$, (the mean sin $\theta$ for the range), are listed in Table 11.l; the resultant K-curve plot, of $K\left(S_{m}\right)$ against $S_{m}$, is shown in Figure 11.2 . Again it was assumed that the K-curve could be fitted by the exponential function (8.3), viz,

$$
\begin{equation*}
K\left(S_{m}\right)=A \operatorname{expB}\left(S_{m}\right)^{C} \tag{11.1}
\end{equation*}
$$

Seven equations of type (11.1) were formulated, having C specified from 1.0 to 4.0 in intervals of 0.5 , and for each the best fit values of $A$ and $B$ were calculated using a least squares method. The coefficients of the seven curves and the corresponding values of SUM, (the sum of the squares of the deviations of the 45 experimental points from the formulated curve), are listed in Table 11.2.

キ The K-curve method is discussed in Chapter 8, p. 55.

| 10 Ranges |  | 15 Ranges |  | 20 Ranges |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $S_{m}$ | $K\left(S_{m}\right)$ | $\mathrm{S}_{\mathrm{m}}$ | $K\left(S_{m}\right)$ | $S_{m}$ | $K\left(S_{m}\right)$ |
| 0.3316 | 2.1115 | 0.2897 | 1.1175 | 0.2632 | 1.1477 |
| 0.4782 | 1.5665 | 0.4177 | 1.2829 | 0.3795 | 1.0509 |
| 0.5670 | 1.8403 | 0.4953 | 1.5543 | 0.4500 | 1.4947 |
| 0.6343 | 2.0667 | 0.5541 | 1.8439 | 0.5034 | 1.6626 |
| 0.6897 | 3.5055 | 0.6025 | 1.8837 | 0.5474 | 1.8259 |
| 0.7374 | 4.3026 | 0.6442 | 2.1470 | 0.5853 | 1.8568 |
| 0.7796 | 4.5198 | 0.6811 | 3.5645 | 0.6188 | 2.0326 |
| 0.8177 | 5.6951 | 0.7143 | 3.7734 | 0.6490 | 2.1028 |
| 0.8525 | 6.8810 | 0.7448 | 4.2729 | 0.6767 | 3.3627 |
| 0.8847 | 7.5412 | 0.7729 | 4.4546 | 0.7022 | 3.6781 |
|  |  | 0.7991 | 4.7466 | 0.7260 | 4.5747 |
|  |  | 0.8237 | 6.2379 | 0.7484 | 4.0572 |
|  |  | 0.8469 | 7.0279 | 0.7695 | 4.9995 |
|  | - | 0.8689 | 6.7931 | 0.7895 | 4.1570 |
|  |  | 0.8899 | 7.8390 | 0.8085 | 5.2954 |
|  |  |  |  | 0.8267 | 6.1872 |
|  |  |  |  | 0.8442 | 6.7282 |
|  |  |  |  | 0.8608 | 7.0533 |
|  |  |  |  | 0.8769 | 7.1286 |
|  |  |  |  | 0.8924 | 8.0494 |

Computed values of $K\left(S_{m}\right)$ and $S_{m}$ used for the $K$-curve generation.
Table 11.1


The K-curve plot.

Figure 11.2

| A | B | C | SUM |
| :---: | :---: | :---: | :---: |
| 0.2827 | 3.6127 | 1.0 | 3.069 |
| 0.5468 | 3.1293 | 1.5 | 2.028 |
| 0.7681 | 2.9687 | 2.0 | 1.669 |
| 0.9498 | 2.9376 | 2.5 | 1.638 |
| 1.1016 | 2.9725 | 3.0 | 1.840 |
| 1.2311 | 3.0470 | 3.5 | 2.276 |
| 1.3437 | 3.1482 | 4.0 | 3.013 |

The 'best fit' coefficients $A$ and $B$ for the seven curves. SUM is the sum of the squares of the deviations of the fortyfive experimental points from the formulated curve.

Table 11.2

Generation of normalised structure amplitudes.
Normalised structure factor magnitudes were computed ${ }^{\text {* }}$ using the two 'best-fitting' K-curves, ( $C=2.0$ and $C=2.5$; Table 11.2). The statistical averages and the percentage distribution of $|E|$ produced are given in Table 11.3 , together with the theoretical values of these quantities. As is evident from Table ll.3, the computed averages and $|E|$ distribution for the two K-curve formulations are not significantly different and in all further work the curve with $C=2.0$ was used. This K-curve therefore yields,

$$
\begin{aligned}
& K^{\prime \prime}=\sqrt{0.7681} \\
& B_{o V}=\frac{2.9687 . \lambda^{2}}{2}=3.876 \\
&=3.53 \AA^{2}
\end{aligned}
$$

where $K^{\prime \prime}$ is the scale factor required to place the observed structure amplitudes on the absolute scale and $B_{o v}$ is the average isotropic temperature factor.

All reflections having $|E| \geqslant 1.5$ were classified by their indices into the eight parity groups and produced the distribution shown in Table 11.4. There are 129 reflections having $|E| \geqslant 2.0$, but only 43 of these are contained in parity groups for which $(k+1)$ is odd. When reflections having $|E| \geqslant 1.5$ are considered, this specialisation in $|E|$ distribution is less severe but still apparent. For the purpose of comparison a similar distribution

[^20]|  | Experimental |  | Theoretical. |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $\begin{gathered} \text { K-curve } \\ 1 . \end{gathered}$ | $\begin{gathered} \text { K-curve } \\ 2 . \end{gathered}$ | Centric | Acentric |
| A | 0.7681 | 0.9498 |  |  |
| B | 2.9687 | 2.9376 |  |  |
| C | 2.0 | 2.5 |  |  |
| $\langle \| E\rangle$ | 0.821 | 0.821 | 0.798 | 0.886 |
| $\langle \| E^{2}-1\| \rangle$ | 0.944 | 0.944 | 0.968 | 0.736 |
| $\left\langle\mathrm{E}^{2}\right\rangle$ | 1.010 | 1.010 | 1.000 | 1.000 |
| \% $/$ \| $\mid \geqslant 3.0$ | 0.2 | 0.1 | 0.3 | 0.01 |
| $\%\|E\| \geqslant 2.0$ | 4.6 | 4.5 | 5.0 | 1.8 |
| $\%\|E\| \geqslant 1.0$ | 31.9 | 31.9 | 32.0 | 37.0 |

The distribution of normalised structure amplitudes and statistical averages found, and the theoretical values of these quantities.

Table 11.3
h kl $|E| \geqslant 2.0$

$$
|E| \geqslant 1.5
$$g g g2652

g gu 1443
tug ..... 938
gu u ..... 27 ..... 56
ugh ..... 16 ..... 52
u gu 11 ..... 38
u Mg ..... 9 ..... 34
$\mathbf{u} \mathbf{u} \mathbf{u}$ 27 ..... 51$\left.\sum \begin{array}{lll}g & g & g \\ u & g & g \\ g & u & u \\ u & u & u\end{array} \right\rvert\,$86211$\left.\sum \begin{array}{lll}g & g & u \\ u & g & u \\ g & u & g \\ u & u & g\end{array} \right\rvert\,$43153

The distribution of $|E|$ within the eight sub-groups, for the trithiacyclopentane from pinacolone.

Table 11.4
of the normalised structure factor magnitudes for the desaurin from acetophenone is shown in Table $11.5^{*}$. The effect of the specialised $\underline{y}$ coordinate of the sulphur atoms of the desaurin is very apparent, in that of 73 reflections having $|E| \geqslant 2.0$ and 210 reflections having $|E| \geqslant 1.5$, only 2 and 13 respectively are contained in parity groups for which ( $k+1$ ) is odd. Although the $|E|$ distribution amongst the eight parity groups is much less severely specialised for the 1,2,4-trithiacyclopentane from pinacolone, it is sufficient to indicate that one or two sulphur atoms of the molecule are probably located 'close' to (020).

* The normalised structure factor magnitudes for the desaurin from acetophenone were computed after completion of the structure analysis.

| hki | $\|E\| \geqslant 2.0$ | $\|E\| \geqslant 1.5$ |
| :---: | :---: | :---: |
| u u u | 18 | 40 |
| u ug | 1 | 1 |
| $u \mathrm{~g} g$ | 17 | 59 |
| $u \mathrm{gu}$ | 0 | 5 |
| g g g | 16 | 54 |
| $\mathrm{g} \mathrm{g} \mathrm{u}^{\text {a }}$ | 1 | 5 |
| $\mathrm{gu} u$ | 20 | 44 |
| gug | 0 | 2 |

$\left\langle\begin{array}{lll}u & u & g \\ u & g & u \\ g & g & u \\ g & u & g\end{array}\right|-$
2
13


71
197

The distribution of $|E|$ within the eight sub-groups, for the desaurin from acetophenone.

## CHAPTER 12.

## THREE $\rightarrow$ DIMENSIONAL PATTERSON AND ELECTRON DENSITY SYNTHESES

In Chapter 2 it was noted that resolution of peaks in the Patterson function is improved by employing sharpened coefficients, ( $\left|{ }^{\mathrm{F}_{\mathrm{hkl}}}\right|^{2}$ ), formulated by equation (2.3), viz,

$$
\left|{ }^{s_{F}}{ }_{h k l}\right|^{2}=\frac{\left|F_{h k l}\right|^{2} \exp \left(2 B_{o v} \cdot R H O\right)}{\sum_{i=1}^{N} f_{i n k l}^{2}}, \quad \text { (12.1) }
$$

where,
$\left|F_{n k l}\right|=$ the observed structure amplitude,
RHO $=(\sin \theta / \lambda)^{2}$,
Bor $=$ the overall isotropic temperature factor,
$f_{i} h k l=$ the scattering factor of the fth atom in a unit cell containing N atoms.

By using the K-curve coefficients, $A$ and $B$, (with $C$ specified as 2.0), equation (12.1) can be rewritten as,

$$
\begin{equation*}
\left|{ }^{s} F_{h k l}\right|^{2}=\frac{\left|F_{h k l}\right|^{2} \cdot A \exp \left[B(\sin \theta)^{2.0}\right]}{\sum_{i=1}^{N} f_{i}{ }^{2} h k l} \tag{12.2}
\end{equation*}
$$

and comparison with equation (8.4) therefore yields,

$$
\left.\right|^{s_{\mathrm{F}}^{\mathrm{hkI}}} 1^{2}-\varepsilon_{\mathrm{hkl}}\left|\mathrm{E}_{\mathrm{hkl}}\right|^{2}
$$

A sharpened Patterson function can therefore be computed by employing $\varepsilon|E|^{2}$ coefficients in the summation.

The origin peak apparent in the Patterson function is removed (69) by employing in the synthesis coefficients of type,

$$
\begin{equation*}
\left[\left|F_{h k l}\right|^{2}-\sum_{i=1}^{N} f_{i}^{2} h k I\right] \tag{12.3}
\end{equation*}
$$

where $\left|{ }^{\prime} F_{h k l}\right|$ is the observed structure amplitude adjusted to the same scale as $f_{i} h k l$, (usually absolute for non-vibrating atoms). By expressing $\left|F_{h k l}\right|^{2}$ in the form $\left\{\left|F_{h k l}\right|^{2} \cdot A \exp B(\sin \theta)^{C}\right\}$, and sharpening the coefficients, (by dividing through by $\sum f_{i}^{2} h k l$ ), the origin-removed sharpened Patterson function can therefore be computed by employing $\left[\varepsilon_{n k l}\left|E_{h k l}\right|^{2}-1\right]$ terms in the summation. The program used to compute the K-curve and calculate normalised structure factor magnitudes, was also coded to produce a binary data tape of $\left(\varepsilon|E|^{2}-1\right)$ and $\varepsilon|E|^{2}$ coefficients. The prepared data tape was utilised as input by a subsequent general Fourier program whence either the sharpened Patterson function or the origin-removed sharpened Patterson function was computed.

## The three-dimensional Patterson function.

The origin-removed sharpened Patterson function was computed for the quarter of the unit cell defined by $-0.5 \leq u \leq 0.5$, $0.0 \leqslant v \leqslant 0.5$ and $0.0 \leqslant w \leqslant 0.5$, in intervals of 0.05 in $u, 0.02$ in $v$ and 0.01 in $w$.

The equivalent positions of space group $\mathrm{P} 2_{1} / \mathrm{c}$ are $\pm(x, y, z)$ and $\pm\left(x, \frac{1}{2}-y, \frac{1}{2}+z\right)$. Peaks in the Patterson function which represent vectors between symmetry related atoms will
(69) See, for example, M. J. Buerger, "Vector Space", Wiley and Sons, New York, (1959), p. 56.
therefore be located at,

|  |  | weight |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $\left.\begin{array}{l}  \pm(2 x, 2 y, 2 z) \\ +(2 x, \overline{2 y}, 2 z) \end{array}\right]$ | Inversion peaks | 1 x | Table 12,1 |
|  |  |  |  |  |
| c. | $\pm\left(2 x, \frac{1}{2}, \frac{1}{2}+2 z\right)$ | Harker section | $2 x$ |  |
| D. | $\pm\left(0, \frac{1}{2}-2 y, \frac{1}{2}\right)$ | Harker line | 2 x |  |

The desaurin from acetophenone crystallises in space group P 2 / c with two molecules, (four sulphur atoms), per unit cell. Thus the only sulphur-sulphur peaks apparent in the

Patterson function are those representing vectors between symmetry related atoms, and by using Table 5.1, ( $\equiv$ Table 12.1), the real space coordinates of the sulphur atoms were easily determined. In the 1,2,4-trithiacyclopentane from pinacolone there are three sulphur atoms per asymmetric molecule and interpretation of the Patterson function will therefore be made more difficult by sulphur 'cross-vector' peaks, (i.e. peaks representing the vectors between sulphur atoms unrelated by space group symmetry). The twelve sulphur atoms in the unit cell will give rise to one hundred and thirty-two sulphur-sulphur vector peaks in the complete Patterson function. From Table 12.1 it can be seen that of these peaks, six will occur on both the Harker section and Harker line, (each of two-fold weight), and twelve inversion peaks will occur in the body of the unit cell. The remaining forty-eight peaks, of two-fold weight, represent the cross-vector sulphur-sulphur interactions. In the quarter of the unit cell for which the

Patterson function was computed, there will be, therefore, three peaks on both the Harker section and Harker line, three single weight inversion peaks and twelve double weight cross~ vector peaks. The most easily identified of the cross-vector peaks will be those representing intramolecular interactions. Although no structural details of the 1,2,4-trithiacyclopentane system were available, (at this time), it was possible to assess approximate dimensions by utilising bond distances and angles found in related sulphur-containing systems. In l,2-dithiolane-4-carboxylic acid (70) and 1,2-dithiane-3,6-dicarboxylic acid (71), disulphide bond lengths of 2.07 and $2.09 \AA$, (respectively), were found, with heterocyclic C-S-S angles in the order 92-990. The structure analyses of the desaurin from acetophenone and PDM yield $C\left(\underline{s p}^{2}\right)-S$ bond distances of $2.75-1.77 \AA$; in addition the $S(I I)-C\left(\underline{s p}^{2}\right)-S(I I)$ bond angle found in PDM is $115^{\circ}$. Structure analyses of six-membered sulphur-containing heterocyclics $(80,81)$, have yielded $C\left(\underline{s p}^{2}\right)-S(I I)-C\left(\underline{s p}^{2}\right)$ angles of $\sim 1000$. The atomic numbering system employed for the trithiacyclopentane ring is,


By applying the dimensions quoted above to this system, approximate sulphur-sulphur distances in the molecule will be,
(70) O. Foss, A. Hordvik and J. Sletten, Acta Chem. Scand., 20, 1169 ( 1966 )
(71) O. Foss, K. Johnson and T. Reistad, ibid., 18, 2345 (1964).

* Structural data for PDM and the desaurin are given in Chapter 14.


In the computed portion of the Patterson function, three prominent peaks, characteristic of intramolecular sulphursulphur vectors, were found, (Table 12.2).

| \# | U | V | W | Peak <br> height | Distance from <br> origin $(\AA)$ | Interpret- <br> ation. |
| ---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 0.100 | 0.100 | 0.080 | 793 | 2.17 | S1-S2 |
| 18 | 0.525 | 0.000 | 0.055 | 1150 | 2.97 | S1-S3 |
| 6 | -0.425 | 0.090 | 0.025 | 631 | 2.99 | S2-S3 |

Peaks representing intramolecular S-S vectors.

## Table 12.2

Peaks 1 and 6 have the normal two-fold weighing associated with the interaction, whereas peak 18, situated on the mirror plane at $V=0$, will be of four-fold weight. The specialised location of peak 18 indicates that S1 and S3 must have approximately the same $y$ coordinate. This set therefore fixes the orientation of the three sulphur atoms of the heterocyclic ring.

Three peaks were observed on the Harker section having relative positions consistent with this orientation, (Table 12.3).

| \#/ | U | $\mathbf{V}$ | W | Peak <br> height | Interpretation. |  |
| :--- | ---: | :---: | :---: | :---: | :---: | :---: |
| 12 | -0.60 | 0.50 | 0.21 | 607 | S1-S1 |  |
| 15 | -0.40 | 0.50 | 0.37 | 695 | S2-S2 |  |
| 14 | 0.45 | 0.50 | 0.32 | 751 |  | S3-S3 |

Peaks on the Harker section ( $u, 1 / 2, w$ ).

Inspection of the Harker line showed only two prominent peaks, (Table 12.4).

| \# | U | v | w | Peak height | Interpretation. |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 16 | 0 | 0.28 | 1/2 | 682 | S2-s2 |
| 17 | 0 | 0.47 | 1/2 | 1052 | $-\mathrm{H}_{\mathrm{S} 3-\mathrm{S} 3}^{\mathrm{Sl}-\mathrm{S1}}$ |
|  | Peaks on the Harker line ( $0, \mathrm{v}, 1 / 2$ ) |  |  |  |  |

Peak 16 has the normal two-fold weight characteristic of the interaction, whilst peak 17 is formed by a superimposition of two such peaks, (the analysis of intramolecular sulphur-sulphur vectors had already indicated that SI and S3 have approximately the same $y$ coordinate).

By using Table 12.1, (C and D, in which the Harker peaks are expressed in terms of general, real space, coordinates of the atoms), together with the intramolecular vector set, the fractional coordinates of the sulphur atoms of the asymmetric molecule were taken as,

|  | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ |
| :--- | :---: | :---: | :---: |
| S1 | -0.300 | 0.010 | -0.145 |
| S2 | -0.200 | 0.110 | -0.065 |
| S3 | 0.225 | 0.020 | -0.090 |

Prominent peaks in the Patterson function corresponding to all sulphur-sulphur cross-vectors indicated that the sulphur atoms were correctly located in the unit cell. The coordinates of the highest twenty-two peaks in the computed portion of the

Patterson function, and the sulphur-sulphur vectors to which they correspond, are given in Table 12.6. The subscript applied to the sulphur atom, in Table 12.6, indicates the equivalent position occupied, and uses the key,
$S_{1}$; $x, y, \quad z \quad$ (the coordinates listed in Table 12.5)
$S_{2}:-x, \quad-y, \quad-z$
$S_{3}: \quad x, \frac{1}{2}-y, \frac{1}{2}+z$
$S_{4} ;-x, \frac{1}{2}+y, \frac{1}{2}-z$
From Table 12.6 it can be seen that the relative peak heights, (each measured from the lowest negative trough in the computed function), compare quite well with the theoretical weight. Peak 13 is 'non-Harker', and represents Sl-S3 vectors; the effective four-fold weighting is again due to the $y$ coordinates of these atoms being approximately the same. The SI inversion peak, (\#19), has double weighting due to the atom lying close to the plane at $y=0$. Peak 11 is too near to the origin, ( $\sim 0.7 \AA$ ), to represent any interatomic vector, and must result from slight errors in the origin removal procedure. Although the resolution in the Patterson function was remarkably good, and many peaks representing vectors between sulphur and light atoms were clearly evident, it was decided to compute electron density syntheses to locate the remaining sixteen light atoms in the molecule.

## Three-dimensional electron density syntheses.

In the desaurin from acetophenone, the sulphur atoms lie very close to (020), causing pseudosymmetry in the electron

| \# | $u$ | v | W | Peak height | Corresponding S - S vectors | $\underset{\mathbf{x}}{\text { Weight }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1. | 0.10 | 0.10 | 0.08 | 793 | SI-S2 intramolecular | 2 |
| 2. | $\overline{0.05}$ | 0.12 | 0.15 | 693 | $\mathrm{S3} 4-\mathrm{S} 2_{3} ; \mathrm{S2}_{4}-\mathrm{S} 3_{3}$ | 2 |
| 3. | $\overline{0.10}$ | 0.38 | 0.42 | 686 | $\mathrm{SI}_{3}-\mathrm{S} 1_{1} ; \mathrm{S3}_{2}-\mathrm{SI} 4$ | 2 |
| 4. | 0.05 | 0.42 | 0.35 | 656 | $\mathrm{S3},-\mathrm{S} 2_{4} ; \mathrm{S2}_{3}-\mathrm{S} 32$ | 2 |
| 5. | 0.45 | 0.38 | 0.48 | 632 | $\mathrm{S3} 3-\mathrm{S} 2_{1} ; \mathrm{S2}_{2}-\mathrm{S} 3_{4}$ | 2 |
| 6. | $\overline{0.425}^{5}$ | 0.09 | 0.025 | 631 | S2-S3 intramolecular | 2 |
| 7. | 0.05 | 0.04 | 0.23 | 606 | $\mathrm{S3}_{4}-\mathrm{Sl}_{3} ; \mathrm{Sl}_{4}-\mathrm{S} 3_{3}$ | 2 |
| 8. | 0.45 | 0.46 | 0.44 | 563 | $\mathrm{S3}_{2}-\mathrm{Sl}_{4} ; \mathrm{Sl}_{3}-\mathrm{S} 3_{1}$ | 2 |
| 9. | 0.40 | 0.22 | 0.13 | 425 | S2 inversion | 2 |
| 10. | $\overline{0.45}$ | 0.04 | 0.18 | 422 | S3 inversion | 1 |
| 11. | 0.00 | 0.04 | 0.02 | 404 |  |  |
| 12. | 0.40 | 0.50 | 0.21 | 607 | Sl Harker (section) | 2 |
| 13: | $\overline{0.05}$ | 0.50 | 0.27 | 1127 |  | 4* |
| 14. | 0.45 | 0.50 | 0.32 | 751 | S3 Harker (section) | 2 |
| 15. | $\overline{0.40}$ | 0.50 | 0.37 | 695 | S2 Harker (section) | 2 |
| 16. | 0.00 | 0.28 | 0.50 | 686 | S2 Harker (line) | 2 |
| 17. | 0.00 | 0.47 | 0.50 | 1052 | $\left[\begin{array}{l} \text { S1 } \\ \text { S3 } \end{array}\right] \text { Harker (Iine) }$ | 4* |
| 18. | $\overline{0.475}$ | 0.00 | 0.055 | 1150 | S1-S3 intramolecular | 4* |
| 19. | 0.425 | 0.00 | 0.285 | 590 | SI inversion | 2* |
| 20. | 0.35 | 0.00 | $0.30_{5}$ | 490 |  |  |
| 21. | 0.50 | 0.12 | 0.21 | 723 | $\mathrm{S2}_{4}-\mathrm{S1} 1_{3} ; \mathrm{S1}_{4}-\mathrm{S} 2_{3}$ | 2 |
| 22. | 0.50 | 0.41 | 0.29 | 600 | $\mathrm{SI},-\mathrm{S} 2_{4} ; \mathrm{S2}_{3}-\mathrm{Sl} 1_{2}$ | 2 |

[^21]density syntheses, which delayed the identification of the light atoms. Interpretation of the Patterson function of the 1,2,4trithiacyclopentane from pinacolone reveals that Sl and S3 exhibit a similar specialisation in position. However, the effect in this structure will be much less severe as both Sl and S3 have a greater displacement from (020), ( $0.13 \AA$ and $0.26 \AA$ compared with $0.01 \AA$ in the desaurin), and in addition S2 will make a 'normal' contribution to all structure factors.

Structure factors were calculated using the sulphur positions determined from the Patterson function and assigning an overall temperature factor of $3.5 \mathrm{~A}^{\circ}$; reflection data classified as unobserved were excluded from the calculation which yielded $R=0.41$. The subsequent electron density synthesis ${ }^{\neq}$was devoid of false symnetry and the light atom locations were clearly revealed as the sixteen highest, (non-sulphur), peaks on the sections. A calculation of interatomic distances and angles showed the molecular dimensions to be reasonable for this stage of the structure analysis. Structure factors were re-calculated from the positions of all nineteen atoms and yielded a residual index of 0.25 . Inspection of the second electron density synthesis showed that all the light atoms included in the structure factor calculation had increased in peak density and that there were no spurious peaks on the sections. The atomic coordinates estimated from this second electron density synthesis are listed in Table 12.7. The atomic coordinates tabulated comprise an

* The electron density syntheses were computed for the quarter of the unit cell and with mesh intervals previously defined for the Patterson synthesis.
(asymmetric) molecule, (the sulphur locations listed are the invertion equivalents of the positions determined from the Patterson function). The atomic numbering system employed is shown in Figure 12.1

Refinement of the atomic parameters by the method of least squares was now initiated.

| Atom | $\mathrm{X} / \mathrm{a}$ | $\mathrm{Y} / \mathrm{b}$ | $\mathrm{Z} / \mathrm{c}$ |
| :--- | :---: | :---: | :---: |
| $\mathrm{S}(1)$ | 0.302 | -0.016 | $0.146_{5}$ |
| $\mathrm{~S}(2)$ | $0.202_{5}$ | -0.110 | $0.066_{5}$ |
| $\mathrm{~S}(3)$ | -0.230 | -0.021 | 0.087 |
| $\mathrm{O}(1)$ | 0.385 | 0.093 | 0.242 |
| $\mathrm{O}(2)$ | $0.067_{5}$ | -0.2355 | $-0.023_{5}$ |
| $\mathrm{C}(1)$ | 0.025 | 0.021 | 0.148 |
| $\mathrm{C}(2)$ | $-0.012_{5}$ | 0.076 | $0.199_{5}$ |
| $\mathrm{C}(3)$ | 0.190 | 0.113 | 0.2475 |
| $\mathrm{C}(4)$ | 0.150 | 0.165 | 0.300 |
| $\mathrm{Me}(5)$ | 0.395 | 0.200 | 0.348 |
| $\mathrm{Me}(6)$ | -0.020 | 0.257 | 0.278 |
| $\mathrm{Me}(7)$ | 0.044 | 0.092 | 0.338 |
| $\mathrm{C}(8)$ | -0.115 | -0.108 | 0.050 |
| $\mathrm{C}(9)$ | $-0.262_{5}$ | -0.168 | 0.006 |
| $\mathrm{C}(10)$ | -0.160 | -0.236 | $-0.028_{5}$ |
| $\mathrm{C}(11)$ | -0.310 | -0.318 | -0.071 |
| $\mathrm{Me}(12)$ | $-0.542_{5}$ | -0.263 | -0.118 |
| $\mathrm{Me}(13)$ | -0.410 | -0.400 | -0.0345 |
| $\mathrm{Me}(14)$ | -0.185 | -0.364 | -0.112 |

Atomic coordinates estimated from the second electron density synthesis.


The atomic numbering system used throughout
the structure analysis.

Figure 12.1

## CHAPTER 13.

## REFINETENT OF THE STRUCTURE.

The standard version of OR-FLS was employed for the initial stages of refinement, but was replaced, for cycles involving anisotropic temperature factors, by an extensively modified version able to accommodate simultaneous refinement of 172 parameters. Each observed structure amplitude was weighted using its standard deviation calculated in the data reduction procedure,

$$
w_{F_{0}}=\frac{1}{\left(\sigma F_{0}\right)^{2}}
$$

An overall scale factor was treated as a variable throughout the refinement; the function minimised was $\sum w\left(\left|F_{0}\right|-\left|F_{c}\right|\right)^{2}$. During the course of the refinement two residual indices are quoted; $R_{1}$ is the residual index obtained using all 2778 reflections for which intensity measurements were made, whilst $R_{2}$ is the index produced when reflections termed unobserved are excluded from the summation.

The initial positional parameters used in the refinement were those obtained from the second electron density synthesis, (Table 12.7), and each atom was allocated an isotropic temperature factor of $3.5 \AA^{2}$. Two cycles of refinement were computed with the thermal parameters of the atoms fixed, followed by four cycles in which all 76 positional and thermal parameters were refined, $R_{1}$ and $R_{2}$ improving to 0.188 and 0.168 respectively.

The parameter shifts calculated in the last cycle were less than one half of the corresponding standard deviation and convertion of the temperature factors into anisotropic mode was now considered in order. The atomic parameters at this stage are listed in Table. 13.1.

It was noted that the observed and calculated structure amplitudes for 011,023 and 024 were in extremely poor agreement for such low order reflections,

|  | $\left\|F_{0}\right\|$ | $\left\|F_{C}\right\|$ |
| :---: | :---: | :---: |
| 011 | 89.8 | 136.8 |
| 023 | 142.7 | 181.1 |
| 024 | 88.2 | 103.2 |

On re-checking the diffractometer output it was revealed that there had been insufficient attenuation available during the intensity data collection for these reflections; consequently they were excluded from all further refinement.

The modified version of OR-FLS was introduced and refinement continued employing anisotropic temperature factors. Two cycles were computed with the thermal parameters of the carbon atoms of the methyl groups fixed, followed by four cycles in which all 172 parameters were refined, $R_{1}$ and $R_{2}$ improving to 0.094 and 0.074 respectively. The maxinum shift calculated in the final cycle represented about $20 \%$ of the standard deviation of the parameter and refinement was considered complete. The final positional and thermal parameters and their estimated standard deviations, calculated from the least squares residuals,

| Atom | X/a | Y/b | 2/c | $B_{i s o}\left(\dot{A}^{2}\right)$. |
| :---: | :---: | :---: | :---: | :---: |
| S(1) | 0.2938 | -0.0140 | 0.1448 | 3.762 |
| S(2) | 0.1952 | -0.1089 | 0.0655 | 3.480 |
| S(3) | -0.2300 | -0.0192 | 0.0880 | 3.229 |
| O(1) | 0.3834 | 0.0925 | 0.2402 | 5.050 |
| O(2) | 0.0541 | -0.2321 | -0.0223 | 5.087 |
| c(1) | 0.0158 | 0.0220 | 0.1494 | 2.650 |
| c(2) | -0.0181 | 0.0774 | 0.1951 | 2.982 |
| c(3) | 0.1855 | 0.1113 | 0.2444 | 3.609 |
| C(4) | 0.1486 | 0.1680 | 0.3000 | 3.892 |
| $\mathrm{Me}(5)$ | 0.3885 | 0.1972 | 0.3447 | 6.874 |
| Me (6) | -0.0108 | 0.2610 | 0.2777 | 6.561 |
| Me (7) | 0.0282 | 0.0938 | 0.3328 | 7.678 |
| c(8) | -0.1110 | -0.1076 | 0.0477 | 2.553 |
| c(9) | -0.2571 | -0.1659 | 0.0048 | 2.638 |
| c(10) | -0.1593 | -0.2355 | -0.0298 | 3.114 |
| c(11) | -0.3150 | -0.3165 | -0.0732 | 3.752 |
| Me(12) | -0.5390 | -0.2663 | -0.1186 | 6.868 |
| Me (13) | -0.3897 | -0.3933 | -0.0309 | 9.014 |
| $\mathrm{Me}(14)$ | -0.1857 | -0.3670 | -0.1138 | 9.271 |

Atomic parameters at the conclusion of the refinement employing isotropic temperature factors.

Table 13.1
are listed in Tables 13.2 and 13.3. $B_{i j}$, in Table 13.3, are coefficients in the anisotropic temperature factor expression,

$$
\begin{aligned}
\exp -1 / 4 & \left(h^{2} a^{* 2} B_{11}+k^{2} b^{*} B_{B_{22}}+1^{2} c^{* 2} B_{33}\right. \\
& \left.+2 h k a^{*} b^{*} B_{12}+2 h l a^{*} c^{*} B_{13}+2 k I b^{*} c^{*} B_{23}\right)
\end{aligned}
$$

From Table 13.3 it can be seen that the themal vibrations of the atoms are 'non-isotropic' in nature. Although the values of $B_{i j}$ produced from a least squares refinement must be treated with caution, since they will tend to encompass systematic errors in the $\left|F_{0}\right|$, those of the carbon atoms of the six methyl groups, ('C(Me)'), appear especially notable. The thermal parameters of the $\mathrm{C}(\mathrm{Me})$ atoms show a high degree of anisotropy and the magnitudes of the motion of these atoms contrast with those of the fourth carbon atom of the t-butyl groups, (C(4) and $C(11))$. A difference synthesis computed at this point did not suggest any corrections to the $C(M e)$ parameters. Although peaks attributable to the hydrogen atoms bonded to $C(2)$ and $C(9)$ were apparent in the synthesis, the regions where 't-butyl hydrogens' are theoretically expected were characterised by smears of low electron density and consequently their locations could not be fixed.

The observed and calculated structure factors, (the latter from the atomic parameters listed in Tables 13.2 and 13.3), are tabulated in Appendix IV. Superimposed sections of a final electron density distribution, taken through the atomic centres parallel to (010), are shown in Figure 13.1, together with a drawing of the corresponding molecular unit.

| Atom | $\sigma(\mathrm{X} / \mathrm{a})$ |  | $\sigma(Y / b)$ |  | $\sigma(z / c)$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | X/a | $\times 10^{4}$ | Y/b | $\times 10^{5}$ | Z/c | $\times 10^{5}$ |
| S(1) | 0.2936 | 3 | -0.0142 | 13 | 0.1446 | 8 |
| S(2) | 0.1956 | 3 | -0.1087 | 13 | 0.0657 | 7 |
| S(3) | -0.2295 | 3 | -0.0192 | 12 | 0.0882 | 7 |
| O(1) | 0.3849 | 7 | 0.0944 | 37 | 0.2407 | 21 |
| O(2) | 0.0562 | 8 | -0.2329 | 35 | -0.0225 | 22 |
| c(1) | 0.0160 | 8 | 0.0203 | 40 | 0.1491 | 24 |
| C(2) | -0.0177 | 10 | 0.0752 | 41 | 0.1950 | 25 |
| c(3) | 0.1879 | 11 | 0.1109 | 45 | 0.2448 | 28 |
| c(4) | 0.1512 | 11 | 0.1664 | 50 | 0.3002 | 27 |
| Me (5) | 0.3902 | 13 | 0.1975 | 69 | 0.3450 | 36 |
| Me (6) | -0.0125 | 15 | 0.2614 | 60 | 0.2770 | 35 |
| Me (7) | 0.0291 | 15 | 0.0898 | 69 | 0.3328 | 36 |
| c(8) | -0.1099 | 9 | -0.1064 | 39 | 0.0481 | 23 |
| c(9) | -0.2566 | 9 | -0.1663 | 40 | 0.0047 | 24 |
| c(10) | -0.1583 | 10 | -0.2360 | 43 | -0.0298 | 26 |
| c(11) | -0.3129 | 11 | -0.3167 | 45 | -0.0725 | 27 |
| Me (12) | -0.5363 | 14 | -0.2656 | 63 | -0.1171 | 37 |
| Me (13) | -0.3882 | 19 | -0.3921 | 65 | -0.0301 | 40 |
| Me(14) | -0.1847 | 16 | -0.3674 | 80 | -0.1121 | 48 |

The final fractional coordinates and standard deviations of the non-hydrogen atoms.

Table 13.2

| Atom | $\mathrm{B}_{11}$ | $\sigma \mathrm{~B}_{11}$ | $\mathrm{~B}_{22}$ | $\sigma \mathrm{~B}_{22}$ | $\mathrm{~B}_{33}$ | $\sigma \mathrm{~B}_{33}$ | $\mathrm{~B}_{12}$ | $\sigma \mathrm{~B}_{12}$ | $\mathrm{~B}_{13}$ | $\sigma \mathrm{~B}_{13}$ | $\mathrm{~B}_{23}$ | $\sigma \mathrm{~B}_{23}$ |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\mathrm{~S}(1)$ | 2.63 | 0.06 | 5.16 | 0.08 | 5.44 | 0.09 | -0.77 | 0.06 | 1.51 | 0.06 | -2.34 | 0.07 |
| $\mathrm{~S}(2)$ | 2.77 | 0.07 | 4.60 | 0.08 | 4.80 | 0.08. | -0.20 | 0.06 | 1.50 | 0.06 | -1.68 | 0.07 |
| $\mathrm{~S}(3)$ | 2.82 | 0.06 | 3.71 | 0.07 | 4.16 | 0.08 | 0.57 | 0.06 | 0.70 | 0.05 | -1.44 | 0.06 |
| $\mathrm{O}(1)$ | 3.11 | 0.21 | 7.77 | 0.31 | 7.42 | 0.29 | -0.57 | 0.20 | 1.91 | 0.21 | -3.59 | 0.25 |
| $\mathrm{O}(2)$ | 3.99 | 0.23 | 5.89 | 0.27 | 7.99 | 0.31 | -0.52 | 0.20 | 2.21 | 0.21 | -3.56 | 0.24 |
| $\mathrm{C}(1)$ | 2.34 | 0.23 | 2.47 | 0.24 | 3.26 | 0.25 | -0.30 | 0.20 | 0.56 | 0.20 | -0.60 | 0.21 |
| $\mathrm{C}(2)$ | 3.77 | 0.28 | 2.61 | 0.26 | 3.33 | 0.28 | -0.42 | 0.22 | 0.61 | 0.23 | -1.13 | 0.23 |
| $\mathrm{C}(3)$ | 4.15 | 0.31 | 2.92 | 0.29 | 4.51 | 0.33 | -0.43 | 0.25 | 1.19 | 0.26 | -0.66 | 0.26 |
| $\mathrm{C}(4)$ | 4.77 | 0.33 | 4.63 | 0.35 | 3.34 | 0.30 | -0.23 | 0.28 | 1.18 | 0.26 | -1.29 | 0.27 |
| $\mathrm{Me}(5)$ | 4.90 | 0.39 | 11.94 | 0.64 | 6.65 | 0.46 | -1.61 | 0.42 | 0.30 | 0.35 | -5.78 | 0.47 |
| $\mathrm{Me}(6)$ | 10.32 | 0.58 | 6.00 | 0.46 | 7.18 | 0.49 | 3.73 | 0.43 | 2.41 | 0.44 | -1.25 | 0.39 |
| $\mathrm{Me}(7)$ | 8.39 | 0.53 | 10.37 | 0.60 | 6.73 | 0.47 | -3.32 | 0.46 | 4.05 | 0.42 | -0.08 | 0.46 |
| $\mathrm{C}(8)$ | 2.98 | 0.24 | 2.05 | 0.23 | 2.80 | 0.24 | 0.57 | 0.20 | 1.18 | 0.20 | -0.19 | 0.20 |
| $\mathrm{C}(9)$ | 3.52 | 0.27 | 2.28 | 0.25 | 3.00 | 0.26 | 0.33 | 0.21 | 1.19 | 0.22 | -0.66 | 0.21 |
| $\mathrm{C}(10)$ | 3.71 | 0.28 | 2.63 | 0.27 | 3.60 | 0.29 | -0.08 | 0.23 | 0.89 | 0.24 | -0.59 | 0.23 |
| $\mathrm{C}(11)$ | 4.17 | 0.31 | 3.06 | 0.29 | 4.17 | 0.33 | -0.01 | 0.24 | 1.07 | 0.26 | -1.29 | 0.25 |
| $\mathrm{Me}(12)$ | 5.90 | 0.42 | 7.50 | 0.50 | 8.12 | 0.52 | 0.55 | 0.40 | -1.21 | 0.38 | -2.16 | 0.44 |
| $\mathrm{Me}(13)$ | 15.30 | 0.83 | 6.43 | 0.51 | 8.02 | 0.55 | -5.68 | 0.55 | 1.35 | 0.54 | 1.22 | 0.45 |
| $\mathrm{Me}(14)$ | 7.88 | 0.55 | 15.38 | 0.83 | 14.93 | 0.79 | -2.56 | 0.53 | 5.12 | 0.55 | -12.88 | 0.70 |

The final anisotropic thermal parameters, ( $\AA^{2}$ ), and their standard deviations.

Table 13.3


Composite electron density normal to [010]. Contour intervals at carbon and oxygen atoms, $1.0 \mathrm{e}^{-3}$, at sulphur atoms, $2.0 \mathrm{e}^{-3}$, , (the zero contour is omitted).

Figure 13.1

## CHAPTER 14.

## DISCUSSION.

Structural data were computed using the 'X-ray 63 system' (72) and a general interatomic distances and angles program (73). The standard deviations in bond lengths and angles involving sulphur, carbon and oxygen atoms, calculated for each molecule, take into account both the standard deviation in atomic position, as estimated from the least squares residuals, and the standard deviations in the cell dimensions.

The desaurin from acetophenone, PDM and the $2,2,4$-trithiacyclopentane from pinacolone have a common chemical grouping of the $\underline{\beta}$-dithia- $\alpha, \beta$-unsaturated carbonyl system,


The three structure analyses thus afford a means of comparing molecular dimensions for this group, and perhaps ascertaining the effect of the different substituents on the system. Cruickshank (74) has emphasised the importance of applying proper significance tests before drawing any conclusions from a comparison of this type. If two independent experimental determinations

[^22]of a specific bond length yield values $L_{1}$ and $L_{2}$ with standard deviations $\sigma_{1}$ and $\sigma_{2}$ respectively, Cruickshank suggests the adoption of the following significance levels;

| $\mathrm{P}>5 \%$ | $\Delta \mathrm{~L}<1.65 \sigma_{A}$ | not significant |
| ---: | ---: | :--- |
| $5 \%>\mathrm{P}>1 \%$ | $1.65 \sigma_{A}<\Delta \mathrm{L}<2.33 \sigma_{A}$ | possibly significant |
| $1 \%>\mathrm{P}>0.1 \%$ | $2.33 \sigma_{A}<\Delta \mathrm{L}<3.09 \sigma_{A}$ | significant |
| $\mathrm{P}<0.1 \%$ | $\Delta \mathrm{~L}>3.09 \sigma_{A}$ | highly significant |

where, $\mathrm{P}=$ the probability that the difference between the two bond lengths is due to experimental error,

$$
\begin{aligned}
& \Delta L=\left(\left|L_{1}-L_{2}\right|\right), \\
& \sigma_{A}=\left(\sigma_{1}^{2}+\sigma_{2}^{2}\right)^{1 / 2}
\end{aligned}
$$

It is evident therefore that only differences of more than about three standard deviations may be taken as definitely significant.

To analyse the packing of the molecules in the unit cells all intermolecular interatomic distances of less than $4.0 \AA$ were calculated for each structure. In general the interatomic contact distances between molecules should not be less than the sum of the van der Waals radii, except in cases where hydrogen bonding is present in the crystal. In this work the van der Waals radii of oxygen, carbon and the methyl group were taken as $1.4 \AA, 2.7 \AA$ and $1.95 \AA$ respectively (75). The van der Waals radius of sulphur quoted by Pauling (75) is $1.85-1.95 \AA$; however there have been suggestions that this value is excessively
(75) L. Pauling, "The Nature of the Chemical Bord", 2nd Edit., Cornell University Press, Ithaca, N. Y., (1940), pp. 187-193.
high. The sulphur-sulphur contact distances in rhombic sulphur indicate a value of $1.65 \pm 0.01 \AA(76)$, whilst crystal structure analyses of other sulphur-containing organic compounds imply values ranging from $1.65-1.75 \AA(77,78,79)$. A constant value for the van der Waals radius cannot strictly be assumed for any atom, such factors as the state of hybridisation of the atom and the electronegativity of atoms attached to. it will affect the value found in any specific case. In this work, however, the minimum observed value of $1.65 \AA$ was assumed for the van der Waals radius of sulphur ${ }^{\neq}$.

Least squares planes were calculated for various atomic groups using the X-ray 63 system. All planes are formulated by an equation of type,

$$
P x+Q y+R z=S,
$$

where $x, y$ and $z$ are atomic fractional coordinates with respect to the crystallographic axes. In the tables of least squares plane data, the atomic deviations from the plane are shown as the distance $\Delta$, (in $\AA)$; an asterisk with this distance indicates that the atom so marked was not used to define the plane.
(76) B. E. Warren and J. T. Burwell, J. Chem. Phys., 3, 6 (1935).
(77) H. L. Yakel, Jn., and E. W. Hughes, Acta Cryst., 2,291 (1954).
(78) R. E. Marsh, Ibid, $\frac{8}{8}$, 91 (1955).
(79) J. Donohue, J. Am. Chem. Soc., 72, 2701 (1950).
$\neq$ The short intramolecular sulphur-oxygen distances found are therefore compared with a van der Waals contact distance of 3.05 A.
14.1 Details of the molecular and crystal structure of the desaurin from acetophenone.

Structural data were calculated from the final atomic coordinates, (and standard deviations), of the sulphur, oxygen and carbon atoms listed in Table 6.2, and the hydrogen atom locations obtained from the difference synthesis, (Table 6.4).

The analysis confirms the 1,3-dithiete structure proposed by Meyer (2) and supported by Yates and Moore (9). The question of the relative stereochemistry of the benzoyl substituents on the $C(2)$ and $C\left(2^{\prime}\right)$ carbons was resolved at an early stage in the diffraction study, since the space group and the presence of two molecules per unit cell require a centrosymmetrical molecule; the desaurin is therefore the trans isomer. A further implication of a molecular centre of symmetry is that the $1,3-$ dithiete ring is planar, in contrast to the situation found in the six-membered sulphur-containing heterocycles XVI, XVII and XVIII, in which the molecules are folded about an axis through


XVI


XVII


XVIII
the heteroatoms $(80,81,82)$. The structure analysis of the desaurin reveals that the $\alpha, \beta$-unsaturated carbonyl system adopts
(80) H. Lynton and E. G. Cox, J. Chem. Soc., 4886 (1956). (81) P. A. Howell, R. M. Curtis and W. N. Lipscomb, Acta Cryst., (82) S. Hosoya, ibid, , 20, 429 (1966).
an $\underline{\text { s-cis }}$ conformation, as illustrated in Figure 14.1.
Bond lengths and angles, (involving non-hydrogen atoms), are shown in Figure 14.11, and are listed, together with their estimated standard deviations, in Tables $14.1_{1}$ and $14.1_{2}$.

The two independent carbon-sulphur bond lengths in the desaurin are remarkably similar, ( $1.766 \AA$ and $1.764 \AA$ ), and are in the centre of the range of $c\left(s p^{2}\right)-S(I I)$ distances found in XVI, XVII and XVIII, ( $1.78 \AA, 1.76 \AA$ and $1.75 \AA$ respectively). As expected, in the dithiete ring, the C-S-C angles, ( $82.0^{\circ}$ ), are smaller than the S-C-S angles, (98.0年). In five- and sixmembered heterocycles containing divalent sulphur, the sulphur valence angles are usually in the range $90-105^{\circ}$ (83), so that the angle found in the l,3-dithiete ring is significantly reduced from its normal value. However, the C-S-C angle in the desaurin is not the smallest observed; in the three-membered ring of ethylene sulphide, the angle subtended at the sulphur atom is $48.3^{\circ}\left(84,85^{\ddagger}\right)$.

The bond distances and angles in the $\alpha, \beta$-unsaturated carbonyl system are discussed in section 14.4 , where a comparative study of the results of the three structure analyses is made. The C(3)-C(4) bond length, $1.495 \AA$, is in accord with distances
(83) S. C. Abrahams, Quart. Rev., X, 407 (1956). (84) H. Dreizler and H. D. Rudolf, Z. Natur., 17a, 712 (1962). (85) R. Desiderato and R. L. Sass, Acta Cryst., 23, 430 (1967).
. $\ddagger$ Desiderato and Sass draw attention to an error in the original paper ( 84 ) where the sulphur valence angie is reported as $58.8^{\circ}$.


Desaurin from acetophenone:

Bond lengths and angles.

Figure 14.11

|  |  | Distances ( ${ }^{\circ}$ ) | $\sigma(\AA)$ |
| :---: | :---: | :---: | :---: |
| s | -c(1) | 1.764 | 0.006 |
| st | - c(1) | 1.766 | 0.006 |
| c(1) | - c(2) | 1.324 | 0.008 |
| c(2) | $-c(3)$ | 1.462 | 0.009 |
| c(3) | - c(4) | 1.495 | 0.008 |
| c(3) | - 0 | 1.216 | 0.007 |
| C(4) | - c(5) | 1.404 | 0.009 |
| c(5) | - c(6) | 1.372 | 0.010 |
| c(6) | - c(7) | 1.385 | 0.011 |
| c(7) | - $c(8)$ | 1.399 | 0.011 |
| c(8) | - c(9) | 1.404 | 0.010 |
| c(9) | - c(4) | 1.383 | 0.009 |
| S | - 0 | 2.640 | 0.006 |
| s | - s' | 2.665 | 0.006 |
| c(1) | - c(1) | 2.315 | 0.009 |

Desaurin from acetophenone:
Intramolecular interatomic distances and their standard deviations.

Table 14.11

|  |  |  |  |  | Angle (0) | $\sigma(0)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| S | - | C(1) | - | S ${ }^{1}$ | 98.0 | 0.4 |
| c(1) | - | S | - | C(1) | 82.0 | 0.4 |
| S | - | C(1) | - | c(2) | 128.5 | 0.5 |
| S ${ }^{1}$ | - | C(1) | - | c(2) | 133.4 | 0.5 |
| c(1) | - | c (2) | - | c(3) | 117.9 | 0.5 |
| C(2) | - | C(3) | - | 0 | 120.2 | 0.6 |
| c(4) | - | c(3) | - | 0 | 119.8 | 0.5 |
| c(2) | - | C(3) | - | c(4) | 120.0 | 0.5 |
| C(3) | - | C(4) | - | c(5) | 117.4 | 0.5 |
| c(3) | - | C(4) | - | c(9) | 122.7 | 0.6 |
| c(9) | - | C(4) | - | c(5) | 119.9 | 0.6 |
| C(4) | - | C(5) | - | c(6) | 120.1 | 0.6 |
| c(5) | - | c(6) | - | c(7) | 120.6 | 0.7 |
| c(6) | - | c(7) | - | c(8) | 120.1 | 0.7 |
| c 7 ) | - | c (8) | - | C(9) | 119.4 | 0.6 |
| c (8) | - | c(9) | - | c(4) | 119.9 | 0.6 |

Desaurin from acetophenone:
Bond angles and their standard deviations.

$$
\text { Table } 24.1_{2}
$$

found in similarly conjugated systens (86). The bond distances and angles in the benzene ring, 1.372-1.404 $\AA$, (mean 1.391 $\AA$ ), and 119.4-120.60, (mean 120.0 ${ }^{\circ}$ ), are consistent with the accepted values. Deviations from the least squares plane for the benzene ring, (Table $14.1_{3}$, plane A), are small; the maximum deviation being $\pm 0.003 \AA$, ( $C(6)$ and $C(7))$. Bond distances and angles involving hydrogen atoms are listed in Table 14.14, (the numbering system employed for the hydrogen atoms is shown in Figure 5.1). The $\mathrm{C}-\mathrm{H}$ bond lengths range from 0.96 to $1.14 \AA$, (mean $1.055^{\circ}$ ), and the $\mathrm{C}-\mathrm{C}-\mathrm{H}$ angles from $100^{\circ}$ to $132^{\circ}$, (mean $120^{\circ}$ ). Little significance can be attached to individual deviations of these bond lengths and angles from the normally accepted values since the hydrogen atoms were not included in the least squares refinement.

The most striking feature of the molecular structure of the desaurin, (also of PDM and the 1,2,4-trithiacyclopentane from pinacolone vide supra), is the close planarity of the ( $\underline{s}$-cis) $\alpha, \beta$-unsaturated carbonyl system and the sulphur atoms, which incurs a correspondingly short intramolecular sulphuroxygen distance. The least squares plane calculated for the central ten atoms of the molecule,

summarised in Table 14.13 $_{3}$, (plane B), shows the maximum atomic
(86) "Tables of Interatomic Distances and Configurations in Molecules and Ions", Chem. Soc. Special Publication, Nos 11 and 18 , London, 1958 and 1965.

| Atom | $\Delta(i)$ |
| :---: | :---: |
| $C(4)$ | +0.0030 |
| $C(5)$ | +0.0001 |
| $C(6)$ | -0.0032 |
| $C(7)$ | +0.0032 |
| $C(8)$ | -0.0001 |
| $C(9)$ | -0.0029 |

## Plane A: $\quad 2.752 x-3.443 y+12.125 z=-0.1237$

| Atom | $\Delta(A)$ |
| :---: | :---: |
| $\left(S^{\prime}\right), S$ | $(-),+0.0060$ |
| $\left(0^{\prime}\right), 0$ | $(+1),-0.0012$ |
| $\left(C\left(1^{\prime}\right)\right), C(1)$ | $(-),+0.0168$ |
| $\left(C\left(2^{\prime}\right)\right), C(2)$ | $(+1,-0.0021$ |
| $\left(C\left(3^{\prime}\right)\right), C(3)$ | $(+),-0.0036$ |

Plane B: $\quad 1.796 x-3.621 y+15.131 z=0.0$

Desaurin from acetophenone;
Least squares plane data.

Table $14.1_{3}$

| Distance ( ${ }_{\text {a }}$ ) |  |  |  |  |  | Angle (0) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $c(2)$ - | H(1) | 1.05 | c(1) | - c(2) | - H(I) | 120. |
| c(5) - | H(2) | 0.96 | c(3) | - c(2) | - H(1) | 122. |
| c(6) - | H(3) | 1.14 | C(4) | - C(5) | - H(2) | 111. |
| $\mathrm{c}(7)$ - | H(4) | 1.01 | c(6) | - c(5) | - H(2) | 127. |
| c(8) - | H(5) | 2.01 | c(5) | - c(6) | - H(3) | 100. |
| c(9) - | H(6) | 1.11 | c(7) | - c(6) | - H(3) | 132. |
|  |  |  | c(6) | - c(7) | - H(4) | 123. |
|  |  |  | c(8) | - c(7) | - H(4) | 117. |
|  |  |  | c(7) | $-\mathrm{c}(8)$ | - H(5) | 122. |
|  | . |  | c(9) | - c(8) | - H(5) | 119. |
|  |  |  | C(8) | - C(9) | - H(6) | 111. |
|  |  |  | c(4) | - C(9) | $-\mathrm{H}(6)$ | 130. |

Desaurin from acetophenone:
Bond distánces and angles involving hydrogen atoms.

Table $14.1_{4}$

Atom ' $A$ ' Atom ' $B$ ' | Unit translation |
| :--- |
| on Atom ' $B$ '. |

| $s$ | $S_{2}$ | $(-1,0,0)$ | 3.42 |
| :--- | :--- | :--- | :--- |
| $s$ | 0, | $(0,-1,0)$ | 3.85 |
| $s$ | $c(3)_{1}$ | $(0,-1,0)$ | 3.63 |
| $s$ | $c(4)_{1}$ | $(0,-1,0)$ | 3.89 |
| 0 | $c(2)_{1}$ | $(-1,0,0)$ | 3.61 |
| 0 | $c(4)_{1}$ | $(0,-1,0)$ | 3.88 |
| 0 | $c(8)_{1}$ | $(-1,-1,0)$ | 3.45 |
| 0 | $c(9)_{1}$ | $(-1,0,0)$ | 3.57 |
| $c(1)$ | $c(1)_{2}$ | $(0,1,0)$ | 3.70 |
| $c(1)$ | $c(2)_{2}$ | $(0,1,0)$ | 3.63 |
| $c(1)$ | $c(9)_{1}$ | $(0,-1,0)$ | 3.59 |
| $c(2)$ | $c(8)_{1}$ | $(0,-1,0)$ | 3.59 |
| $c(2)$ | $c(9)_{1}$ | $(0,-1,0)$ | 3.42 |
| $c(3)$ | $c(8)_{1}$ | $(0,-1,0)$ | 3.50 |
| $c(3)$ | $c(9)_{1}$ | $(0,-1,0)$ | 3.64 |
| $c(4)$ | $c(7)_{1}$ | $(0,-1,0)$ | 3.65 |
| $c(4)$ | $c(8)_{1}$ | $(0,-1,0)$ | 3.63 |
| $c(5)$ | $c(7)_{1}$ | $(0,-1,0)$ | 3.63 |
| $c(5)$ | $c(8)_{1}$ | $(-1,0,0)$ | 3.46 |
| $c(5)$ | $c(9)_{1}$ | $(-1,0,0)$ | 3.74 |

Atom 'B' is generated from the coordinates listed in Table 6.2 using the symmetry operation ()$_{i}$ and unit translations shown.
$l_{1}: \quad x, y, z$
$l_{2}: \quad-x,-y,-z$

Desaurin from acetophenone:

Closest contacts of neighbouring molecules.

Table 14.15


Desaurin from acetophenone:
Projection of the structure along [100], illustrating the packing of the molecules.

Figure 14.12.
deviation from the plane is $\mp 0.017 \AA$, ( $C(1)$ and $C\left(I^{\prime}\right)$ ), with a mean atomic deviation of $0.006 \AA$. However, the molecule as a whole is not planar; the plane of the benzene ring is inclined at an angle of $11.4^{\circ}$ to that of the central ten atom system. The intramolecular sulphur-oxygen distance is $2.640 \AA$, being $0.41 \AA$ less than the sum of the van der Waals radii of sulphur and oxygen. The approximate planarity of the central portion of the molecule and the short intramolecular sulphur-oxygen distance are further discussed in section 14.4.

Intermolecular interatomic distances of less than $3.9 \AA$ are listed in Table $14.1_{5}$, and the molecular packing is illustrated in Figure 14.12. The contact distances between molecules all appear reasonable; the only interatomic separation of note is the sulphur-sulphur contact of $3.42 \AA$, which provides additional support that the van der Waals radius of sulphur is somewhat less than the value suggested by Pauling.

### 14.2 Details of the molecular and crystal structure of 3-phenyd-

 1-propene-1,3-dione 1-(dimethyl mercaptole) ${ }^{*}$.Structural data were calculated from the final atomic coordinates, (and standard deviations), of the sulphur, oxygen and carbon atoms, listed in Table 10.2, and the hydrogen locations obtained from the difference synthesis, (Table 10.4).

The structure analysis reveals that the $\alpha, \beta$-unsaturated

[^23]carbonyl system adopts an s-cis conformation, analogous to the situation found in the desaurin. Bond distances and angles, (involving non-hydrogen atoms), are shown in Figure 14.21, and are listed, together with their estimated standard deviations, in Tables 14.2 , and 14.22 .

The $C(1)-S$ bond distances in PDM, $1.745 \AA$ and $1.748 \AA$, are identical within experimental error, but are shorter than the corresponding lengths in the desaurin, (1.766 $\AA$ and $1.764 \AA$ ); however the differences in these bond lengths between the two molecules are not significant, representing about $1.4 \sigma_{A}$. The major differences between the molecular structures of PDM and the desaurin lie in the valence angles at the sulphur atoms and in the $S-C(1)-S$ angle;

|  | PDM | Desaurin |
| :---: | :---: | :---: |
| S-C(1)-S | 115.10 | 98.00, |
| $\widehat{\mathbf{S}}$ | $-1 \begin{aligned} & 104.4^{\circ} \\ & 204.30 \\ & \end{aligned}$ | $82.0^{\circ}$. |

These angles in the 1,3 -dithiete ring of the desaurin are, as expected, significantly smaller than the corresponding angles in PDA. Consequently the intramolecular sulphur-sulphur distance in PDM, ( $2.948^{\circ} \mathrm{A}$ ), is appreciably longer than in the desaurin, (2.665 $\AA$ ). The $\mathrm{S}-\mathrm{C}\left(\mathrm{sp}^{3}\right)$ bond lengths, $1.823 \AA$ and $1.843 \AA$, are not significantly different and are consistent with the accepted carbon-sulphur single bond length of $1.82 \AA$ ( 83 ).

Bond distances and angles in the $\alpha, \beta$-unsaturated carbonyl system show little difference from the values found in similarly


PDM:
Bond lengths and angles.
Figure 14.21

|  |  |  | Distance ( ${ }_{\text {A }}$ ) | $\sigma(A)$ |
| :---: | :---: | :---: | :---: | :---: |
| S(1) | - | C(1) | 1.745 | 0.011 |
| S(1) | - | $\mathrm{Me}(10)$ | 1.843 | 0.011 |
| S(2) | - | c(1) | 1.748 | 0.010 |
| S(2) | - | Me (11) | 1.823 | 0.011 |
| 0 | - | C(3) | 1.246 | 0.013 |
| C(1) | - | c(2) | 1.378 | 0.013 |
| c(2) | - | c(3) | 1.440 | 0.014 |
| C(3) | - | c(4) | 1.511 | 0.013 |
| C(4) | - | c (5) | 1.386 | 0.016 |
| c(4) | - | c(9) | 1.414 | 0.016 |
| c(5) | - | c(6) | 1.408 | 0.014 |
| $c(6)$ | - | c (7) | 1.407 | 0.018 |
| c(7) | - | c (8) | 1.411 | 0.018 |
| c(8) | - | c(9) | 1.388 | 0.014 |
| S(1) | - | S(2) | 2.948 | 0.007 |
| S(1) | - | 0 | 2.727 | 0.010 |

PDM:
Intramolecular interatomic distances and their standard deviations.

Table 14.21

| C(1) | - | $S(1)$ | - | Angle (0) |  | $\sigma(0)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | Me(10) | 104.3 | 0.5 |
| C(1) | - | $S(2)$ | - | Me(11) | 104.4 | 0.5 |
| $S(1)$ | - | C(1) | - | $S(2)$ | 215.1 | 0.5 |
| S(1) | - | C(1) | - | C(2) | 123.0 | 0.8 |
| S (2) | - | C(1) | - | C(2) | 121.9 | 0.8 |
| C(1) | - | C(2) | - | C(3) | 122.5 | 0.9 |
| 0 | - | C(3) | - | c(2) | 122.5 | 0.8 |
| 0 | - | C(3) | - | C(4) | 119.9 | 0.9 |
| C(2) | - | C(3) | - | C(4) | 117.6 | 0.9 |
| c(3) | - | C(4) | - | C(5) | 122.4 | 2.0 |
| c(3) | - | C(4) | - | C(9) | 116.6 | 1.0 |
| c(5) | - | C (4) | - | C(9) | 120.9 | 0.9 |
| C(4) | - | C(5) | - | c (6) | 119.7 | 1.1 |
| C(5) | - | c(6) | - | C(7) | 119.6 | 1.1 |
| c(6) | - | C (7) | - | c(8) | 120.3 | 1.0 |
| C(7) | - | C(8) | - | c(9) | 119.7 | 1.1 |
| C(4) | - | C(9) | - | C(8) | 119.8 | 1.1 |

PDM:

Bond angles and their standard deviations.
$\xlongequal{\text { Table } 14.22}$
conjugated atomic groups; the dimensions of this system are discussed in detail, and compared with the analogous dimensions in the desaurin and trithiacyclopentane from pinacolone, in section 14.4 . The $C(3)-C(4)$ bond length, $1.511 \AA^{\circ}$, is comparable with the analogous distance in the desaurin of $1.495 \AA$, and the bond distances and angles in the benzene ring, 1.386-1.414 ${ }^{\circ}$, (mean 1.402 $\AA$ ), and 219.6-120.90, (mean $120.0^{\circ}$ ) are consistent with the accepted values. Deviations from the least squares plane of the benzene ring, summarised in Table $14.2_{3}$, (plane A), are small; the maximum atomic displacement from the plane is $0.019 \AA,[\mathrm{C}(4)]$, and the mean atomic deviation is $0.010 \AA$.

Bond lengths and angles involving hydrogen atoms are listed in Table $14.2_{4}$, (the numbering system used for the hydrogen atoms is shown in Figure 10.1). As in the desaurin, little significance can be attached to any individual deviation of these distances and angles from the accepted values since the hydrogen atoms were not included in the least squares refinement. The mean values of the bond lengths and angles involving hydrogen atoms are, $1.08 \AA$ for $C\left(\underline{s p}^{2}\right)-H ; 1.06 \AA$ for $C\left(\underline{s p}^{3}\right)-H ;$ $119^{\circ}$ for $\mathrm{C}-\mathrm{C}\left(\underline{s p}^{2}\right)-\mathrm{H} ; 105^{\circ}$ for $\mathrm{S}-\mathrm{C}\left(\underline{s p}^{3}\right)-\mathrm{H}$ and $114^{\circ}$ for $\mathrm{H}-\mathrm{C}\left(\right.$ sp $\left.^{3}\right)-\mathrm{H}$.

The coplanarity of the $\alpha, \beta$-unsaturated carbonyl system and the sulphur atoms is still much in evidence in PDM, though to a lesser degree than found in the desaurin structure. Various least squares planes were calculated, and are summarised in Table $14.2_{3}$, (planes $B, C$ and $\underline{D}$ ). Least squares plane $B$, calculated for the $S(1), S(2), C(1)$ and $C(2)$ system, illustrates
$\frac{\text { Plane } \underline{A}}{\Delta(\AA)} \quad \frac{\text { Plane B }}{\Delta(\AA)} \quad \frac{\text { Plane C }}{\Delta(\AA)} \quad \frac{\text { Plane } \underline{D}}{\Delta(\AA)}$

| $S(1)$ | 0.0004 | 0.0136 | 0.0149 |
| :--- | :---: | :---: | :---: |
| $S(2)$ | 0.0004 | $-0.0127^{*}$ | -0.0043 |
| 0 |  | $-0.0308^{*}$ | -0.0275 |
| $c(1)$ | -0.0013 | -0.0076 | -0.0305 |
| $c(2)$ |  | 0.0005 | -0.0152 |
| $c(3)$ | $0.0468^{*}$ | 0.0366 | -0.0123 |
| $c(4)$ | -0.0194 |  |  |
| $c(5)$ | 0.0095 |  |  |
| $c(6)$ | 0.0050 |  |  |

$c(7) \quad-0.0094$
$C(8) \quad-0.0006$
c(9) . 0.0148

| $\mathrm{Me}(10)$ |  | $0.2125 *$ | $0.2325 *$ | $0.2366 *$ |
| :--- | :---: | :---: | :---: | :---: |
| $\mathrm{Me}(11)$ |  | $0.1635 *$ | $0.1302 *$ | $0.1407 *$ |
| P | -2.6616 | -3.5265 | -3.5295 | -3.5396 |
| Q | 6.9544 | 3.9428 | 4.0432 | 4.0231 |
| R | 13.836 | 16.955 | 16.854 | 16.858 |
| S | 6.5785 | 6.9277 | 6.9067 | 6.9016 |

PDM;
Least squares plane data.

$$
\text { Table } 14.23
$$

Distance $(\dot{A})$

| $\mathrm{C}(2)-\mathrm{H}(1)$ | 1.04 | $\mathrm{C}(1)-\mathrm{C}(2)-\mathrm{H}(1)$ | 127. | $\mathrm{S}(1)-\mathrm{Me}(10)-\mathrm{H}(7)$ | 93. |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\mathrm{C}(5)-\mathrm{H}(2)$ | 1.15 | $\mathrm{C}(3)-\mathrm{C}(2)-\mathrm{H}(1)$ | 109. | $\mathrm{S}(1)-\mathrm{Me}(10)-\mathrm{H}(8)$ | 111. |
| $\mathrm{C}(6)-\mathrm{H}(3)$ | 0.96 | $\mathrm{C}(4)-\mathrm{C}(5)-\mathrm{H}(2)$ | 127. | $\mathrm{~S}(1)-\mathrm{Me}(10)-\mathrm{H}(9)$ | 113. |
| $\mathrm{C}(7)-\mathrm{H}(4)$ | 1.04 | $\mathrm{C}(6)-\mathrm{C}(5)-\mathrm{H}(2)$ | 111. | $\mathrm{H}(7)-\mathrm{Me}(10)-\mathrm{H}(8)$ | 87. |
| $\mathrm{C}(8)-\mathrm{H}(5)$ | 1.12 | $\mathrm{C}(5)-\mathrm{C}(6)-\mathrm{H}(3)$ | 147. | $\mathrm{H}(7)-\mathrm{Me}(10)-\mathrm{H}(9)$ | 138. |
| $\mathrm{C}(9)-\mathrm{H}(6)$ | 1.19 | $\mathrm{C}(7)-\mathrm{C}(6)-\mathrm{H}(3)$ | 93. | $\mathrm{H}(8)-\mathrm{Me}(10)-\mathrm{H}(9)$ | 113. |
| $\mathrm{Me}(10)-\mathrm{H}(7)$ | 0.92 | $\mathrm{C}(6)-\mathrm{C}(7)-\mathrm{H}(4)$ | 112. | $\mathrm{~S}(2)-\mathrm{Me}(11)-\mathrm{H}(10)$ | 102. |
| $\mathrm{Me}(10)-\mathrm{H}(8)$ | 1.27 | $\mathrm{C}(8)-\mathrm{C}(7)-\mathrm{H}(4)$ | 127. | $\mathrm{~S}(2)-\mathrm{Me}(11)-\mathrm{H}(11)$ | 94. |
| $\mathrm{Me}(10)-\mathrm{H}(9)$ | 0.90 | $\mathrm{C}(7)-\mathrm{C}(8)-\mathrm{H}(5)$ | 126. | $\mathrm{~S}(2)-\mathrm{Me}(11)-\mathrm{H}(12)$ | 114. |
| $\mathrm{Me}(11)-\mathrm{H}(10)$ | 1.07 | $\mathrm{C}(9)-\mathrm{C}(8)-\mathrm{H}(5)$ | 114. | $\mathrm{H}(10)-\mathrm{Me}(11)-\mathrm{H}(11)$ | 112. |
| $\mathrm{Me}(11)-\mathrm{H}(11)$ | 1.12 | $\mathrm{C}(8)-\mathrm{C}(9)-\mathrm{H}(6)$ | 125. | $\mathrm{H}(10)-\mathrm{Me}(11)-\mathrm{H}(12)$ | 103. |
| $\mathrm{Me}(11)-\mathrm{H}(12)$ | 1.09 | $\mathrm{C}(4)-\mathrm{C}(9)-\mathrm{H}(6)$ | 112. | $\mathrm{H}(11)-\mathrm{Me}(11)-\mathrm{H}(12)$ | 130. |

PDM:

Angle ( 0 )
Angle ( 0 )
93.
11.
7.
38.
13.
4.


that $C(1)$ shows no significant deviation from the formal sp $^{2}$ hybrid state and acts as a base-line for the other calculated planes; atoms C(3) and $O$ are displaced on either side of plane $B$ by $0.047 \AA$ and $0.031 \AA$ respectively. The X-ray diffraction study of 2-desylidene-1,3-dithiolane, XIV, by Tulinsky (15,16),


XIV*
showed that the $S 1-C=C-C=O$ system is approximately coplanar (the mean atomic deviation being $0.034 \AA$ ), but that $S 2$, which is not involved in any sulphur-oxygen interaction, is sienificantly displaced from the plane by $0.24 \AA$. Consequently least squares plane $C$ was calculated for the analogous $S(1)-C(1)=C(2)-C(3)=0$ system of PDM. The out-of-plane distances for plane $\mathbf{C}$ show a direct contrast to the findings of Tulinsky; $S(2)$ is displaced $0.013 \AA$ from the plane and is, in fact, closer to the plane than four of the defining atoms. Least squares plane $\underline{D}$, calculated for the $S(1), S(2), C(1), C(2), C(3)$ and 0 systen, is the closest comparison available to the ten atom plane of the desaurin. The maximum atomic deviation from plane $\underline{D}$ and the mean atomic deviation are $0.036 \AA,[C(3)]$, and $0.017 \AA$ respectively, (compared with $0.017 \AA$ and $0.006 \AA$ respectively for plane $B$ of the desaurin). The intramolecular S(1)-0 distance in PDM, $2.727 \AA^{\circ}$, is greater

[^24]| Atom 'A' | Atom 'B' | Unit translation on Atom 'B'. | Distance ( $\AA$ ) . |
| :---: | :---: | :---: | :---: |
| S(1) | $\mathrm{S}(\mathrm{I})_{2}$ | $(1,0,1)$ | 3.41 |
| S(1) | $\mathrm{Me}(10)_{2}$ | ( 1, 0, 1) | 3.89 |
| S(2) | $\mathrm{S}(2)_{2}$ | ( 2, 1, 1) | 3.57 |
| S(2) | $\mathrm{C}(2)_{2}$ | ( 1, 1, 1) | 3.88 |
| S(2) | C(3) 1 | ( $1,0,0$ ) | 3.69 |
| S(2) | C(4), | $(1,0,0)$ | 3.77 |
| $s(2)$ | c(9) ${ }_{1}$ | ( $1,0,0$ ) | 3.79 |
| S(2) | $\mathrm{Me}(11){ }_{2}$ | $(2,1,1)$ | 3.80 |
| 0 | $\mathrm{C}(5)_{4}$ | $(0,-1,0)$ | 3.88 |
| 0 | $\mathrm{C}(6){ }_{4}$ | $(0,-1,0)$ | 3.41 |
| 0 | $\mathrm{Me}(10)_{1}$ | $(-1,0,0)$ | 3.53 |
| 0 | $\mathrm{Me}(10)_{2}$ | ( $1,0,1$ ) | 3.47 |
| c(1). | C(9), | ( $1,0,0)$ | 3.46 |
| c(1) | $\mathrm{Me}(11){ }_{2}$ | ( $1,1,1$ ) | 3.71 |
| $c(2)$ | $\mathrm{c}(8){ }_{1}$ | $(1,0,0)$ | 3.56 |
| $c(2)$ | c(9), | ( $1,0,0)$ | 3.56 |
| c(3) | $\mathrm{c}(8) 1$ | ( $1,0,0$ ) | 3.88 |
| C(4) | $\mathrm{Me}(11)$, | $(-1,0,0)$ | 3.80 |
| c(6) | $\mathrm{C}(9) 4$ | $(-1,0,0)$ | 3.74 |
| c(6) | $\mathrm{Me}(10)_{3}$ | $(-1,0,-1)$ | 3.68 |
| c(7) | $\mathrm{C}(9)_{4}$ | $(-1,0,0)$ | 3.89 |
| c(7) | $\mathrm{Me}(11){ }_{4}$ | $(0,-1,0)$ | 3.74 |
| C(8) | $\mathrm{Me}(11) 4$ | $(1,-1,0)$ | 3.45 |

Atom ' $B$ ' is generated from the coordinates listed in Table 10.2 using the symmetry operation ()$_{i}$ and unit translation shown. $\begin{array}{llrl}l_{1}: x, y, & z . & l_{3}: x, 1_{2}-y, l_{2}+z . \\ l_{2}:-x,-y,-z . & l_{4}:-x, 1_{2}+y, l_{2}-z .\end{array}$

PDM:

> Closest contacts of neighbouring molecules.

## Table 14.25


than the corresponding length in the desaurin, ( $2.640 \AA$ ), but is still appreciably less than the sum of the van der Waals radii of sulphur and oxygen. The approximate coplanarity of the $\alpha, \beta-$ unsaturated carbonyl system and the sulphur atoms, and the short sulphur-oxygen distance, are further discussed in section 14.4. As in the desaurin, the plane of the benzene ring is inclined to that of the sub-system, and in this molecule is rotated $19.3^{\circ}$ from plane $\underline{D}$.

Intermolecular interatomic distances of less than $3.9 \AA^{\circ}$ are listed in Table 14.25 , and the molecular packing is illustrated in Figure 14.22. The contact distances between molecules range upwards from $3.41 \AA$, ( $C(6)-0 ; S(1)-S(1))$, and appear reasonable, with the exception of the $C(8)-\mathrm{Me}(11)$ separation of $3.45 \AA$. This distance is less than the theoretical carbon-methyl group contact distance of $3.65 \AA$, although it is greater than twice the van der Waals radius of a carbon atom. Examination of a threedimensional difference synthesis in the region of $S(2)$ revealed no density which might suggest an alternative position for Me(11). The sulphur-sulphur separation, $3.41 \AA$, across a space group inversion centre, is very similar to the situation found in the desaurin.

### 14.3 The molecular and crystal structure of the $1,2,4$-trithiacyclopentane from pinacolone.

Structural data were calculated from the final atomic coordinates, (and standard deviations), listed in Table 13.2. Bond lengths and angles are shown in Figure 14.31, and are


Trithiacyclopentane from pinacolone:
Bond lengths and angles.
Figure 14.3 ,

Iisted, together with their estimated standard deviations, in Tables 14.3 , and 14.32 . The structure analysis confirms the 1,2,4-trithiacyclopentane structure forwarded by Gompper (7) and supported by Lynch (10), and reveals that the molecule is the cis isomer VII $_{a}$, with both $\alpha, \beta$-unsaturated carbonyl systems adopting an s-cis conformation, as illustrated in Figure 14.3..

Hordvik ( 87 ) conducted an extensive survey of structural data for molecules containing an S(II)-S(II) bond which showed that the dihedral angle ${ }^{\neq}$between the valences of the sulphur atoms is correlated to the disulphide bond length. The variation in disulphide bond length with dihedral angle was assumed to be in part due to lone pair repulsion, which is most pronounced when the dihedral angle is $0^{\circ}$, (cis), and in part due to pi bonding, which is most pronounced at dihedral angle $90^{\circ}$. A plot of the dihedral angle against the S(II)-S(II) bond length, for a number of linear and unsaturated heterocyclic disulphides, resulted in a smooth curve from which Hordvik estimates the length of a single bond between two divalent sulphur atoms at dihedral angle zero is $2.10 \AA$.

Least squares plane data for various groups of atoms in the trithiacyclopentane from pinacolone molecule are listed in Table 14.33. Plane A, calculated for the $C(1)-S(1)-S(2)-C(8)$ system, shows this grouping is approximately planar, (the mean atomic deviation is $0.013{ }^{\circ}$ ), but that $S(3)$ is significantly
(87) A. Hordvik, Acta Chem. Scand. 20 , 1885 (1966).
$\ddagger$ In the disulphide system $-C_{1}-S_{1}-S_{2}-C_{2}-$, the angle between the $C_{1} S_{1} S_{2}$ and $S_{1} S_{2} C_{2}$ planes is termed the "dihedral angle".

|  |  |  | Distance ( ${ }_{\text {A }}$ ) | $\sigma(\AA)$ |
| :---: | :---: | :---: | :---: | :---: |
| S(I) | - | S(2) | 2.103 | 0.003 |
| S(1) | - | $C(1)$ | 1.741 | 0.006 |
| S(2) | - | C(8) | 1.732 | 0.006 |
| S(3) | - | C(1) | 1.748 | 0.005 |
| S(3) | - | c(8) | 1.739 | 0.006 |
| O(1) | - | c(3) | 1.221 | 0.008 |
| O(2) | - | $c(10)$ | 1.233 | 0.008 |
| C(1) | - | c(2) | 1.330 | 0.008 |
| C(2) | - | C(3) | 1.460 | 0.007 |
| C(3) | - | C(4) | 1.525 | 0.010 |
| C(8) | - | c(9) | 1.345 | 0.007 |
| C(9) | - | c(10) | 1.436 | 0.009 |
| C(10) | - | C(11) | 1.529 | 0.008 |
| C(4) | - | $\mathrm{Me}(5)$ | 1.523 | 0.009 |
| C(4) | - | $\mathrm{Me}(6)$ | 1.561 | 0.010 |
| C(4) | - | $\mathrm{Me}(7)$ | 1.550 | 0.012 |
| C(11) | - | $\mathrm{Me}(12)$ | 1.548 | 0.009 |
| c(11) | - | $\mathrm{Me}(13)$ | 1.536 | 0.012 |
| C(11) | - | $\mathrm{Me}(14)$ | 1.500 | 0.014 |
| $S(1)$ | - | O(1) | 2.517 | 0.005 |
| S(2) | - | O(2) | 2.509 | 0.005 |

Trithiacyclopentane from pinacolone:
Intramolecular interatomic distances and their standard deviations.

|  |  |  |  | Angle (0) | $\sigma(0)$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| S(2) | - | S(1) - | C(1) | 100.6 | 0.2 |
| S(1) | - | S(2) - | C(8) | 100.4 | 0.2 |
| c(1) | - | S(3) | C(8) | 103.4 | 0.3 |
| S(1) | - | c(1) - | S(3) | 116.7 | 0.3 |
| S(1) | - | c(1) - | c(2) | 124.0 | 0.4 |
| S(3) | - | c(l) - | C(2) | 119.3 | 0.4 |
| c(1) | - | $c(2)$ - | C(3) | 119.1 | 0.6 |
| O(1) | - | c(3) | C(2) | 118.4 | 0.6 |
| O(1) | - | C(3) | C(4) | 122.1 | 0.5 |
| c(2) | - | c(3) - | C(4) | 119.5 | 0.6 |
| c(3) | - | c(4) | Mie (5) | 109.7 | 0.6 |
| c(3) | - | C(4) | Me (6) | 109.3 | 0.5 |
| c(3) | - | C(4) | Me (7) | 107.3 | 0.6 |
| $\mathrm{Me}(5)$ | - | c(4) - | Me (6) | 111.7 | 0.6 |
| Me (5) | - | c(4) - | Me (7) | 109.2 | 0.6 |
| $\mathrm{Me}(6)$ | - | C(4) | $\mathrm{Me}(7)$ | 109.7 | 0.7 |
| S(2) | - | c(8) | S(3) | 117.3 | 0.3 |
| S(2) | - | c(8) - | C(9) | 123.6 | 0.5 |
| S(3) | - | C(8) | C(9) | 119.1 | 0.4 |
| c(8) | - | $c(9)-$ | C(10) | 119.1 | 0.5 |
| O(2) | - | c(10) - | c(9) | 118.9 | 0.5 |
| O(2) | - | c(10) - | C(11) | 120.0 | 0.6 |
| c(9) | - | c(10) - | C(11) | 121.0 | 0.5 |
| c(10) | - | c(11) - | $\mathrm{Me}(12)$ | 110.2 | 0.5 |
| c(10) | - | c(11) - | $\mathrm{Me}(13)$ | 106.0 | 0.5 |
| $\mathrm{C}(10)$ | - | c(11) - | $\mathrm{Me}(14)$ | 111.6 | 0.6 |
| $\mathrm{Me}(12)$ | - | c(11) - | Me (13) | 109.2 | 0.6 |
| Me (12) | - | c(11) - | Me(14) | 106.7 | 0.6 |
| Me (13) | - | c(11) - | Mie (14) | 113.3 | 0.7 |

Trithiacyclopentane from pinacolone:
Bond angles and their standard deviations.

Table 14.32

| Plane $\underline{A}$ | Plane B | Plane C | Plane D |
| :---: | :---: | :---: | :---: |
| $\Delta(\AA)$ | $\Delta(\AA)$ | $\Delta(\AA)$ | $\Delta(\AA)$ |


| $S(1)$ | -0.0145 | $0.3170^{*}$ | -0.0218 | $-0.2672^{*}$ |
| :--- | :---: | :--- | :---: | :---: |
| $S(2)$ | 0.0145 | $0.3611^{*}$ | $-0.1474^{*}$ | -0.0268 |
| $S(3)$ | $0.2145^{*}$ | 0.0 | $0.1180^{*}$ | $0.1395^{*}$ |
| $O(1)$ | $-0.1438^{*}$ |  | 0.0329 |  |
| $O(2)$ | $-0.1729^{*}$ |  |  | 0.0417 |
| $C(1)$ | 0.0111 | 0.0 | 0.0206 |  |
| $C(2)$ | $-0.0938^{*}$ |  | 0.0031 |  |
| $C(3)$ | $-0.2204^{*}$ |  | -0.0349 |  |
| $C(8)$ | -0.0112 | 0.0 |  | 0.0238 |
| $C(9)$ | $-0.1661^{*}$ |  |  | 0.0070 |
| $C(10)$ | $-0.3098^{*}$ |  |  | -0.0457 |
|  |  |  |  |  |
| $P$ | -0.9624 | -2.1053 | -0.9892 | -0.7946 |
| $Q$ | -10.299 | -9.6764 | -10.834 | -9.5099 |
| $R$ | 13.898 | 14.763 | 12.613 | 15.424 |
| $S$ | 1.8589 | 1.9710 | 1.6655 | 1.8649 |

Trithiacyclopentane from pinacolone:
Least squares plane data.

Table 14.33
displaced from the plane by $0.215 \AA$. Consequently the heterocycle is most easily defined by plane $A$ and a second plane through the $C(1)-S(3)-C(8)$ system, (plane $B$ ), which intersect at an angle of $11.5^{\circ}$. The sulphur-sulphur bond length found in this cis, effectivly-planar ${ }^{*}$ disulphide system is 2.103 ${ }^{\circ}$, which according to Hordvik's work (87) must be considered as an example of a $S(I I)-S(I I)$ single bond with no pi character. It is evident, however, that if the short intramolecular sulphuroxygen distances found in this molecule, PDM and the desaurin from acetophenone are interpreted as partial bonding between these atoms, the net observed effect in this molecule could be a lengthening of the disulphide bond. This point is further discussed in section 14.4 .

The four independent $C\left(\underline{s p}^{2}\right)-S$ bond lengths in the molecule, (1.741 $1.1 .748 \AA, 1.739 \AA$ and $1.732 \AA$ ), are not significantly different and the mean value of $1.740 \AA$ is consistent with the corresponding mean distance in PDM of $1.747 \AA$. The valence angles at $S(1)$ and $S(2), 100.6^{\circ}$ and $100.4^{\circ}$, are identical within experimental error, but are smaller than the angle subtended at $S(3),\left(103.4^{\circ}\right)$. The endocyclic angles at $C(1)$ and $C(8)$, $116.7^{\circ}$ and $117.3^{\circ}$, are comparable with the analogous S-C(1)-S bond angle of $115.1^{\circ}$ in PDM.

The molecular dimensions of the two $\alpha, \beta$-unsaturated carbonyl systems in the molecule are not significantly different, (Figure 14.31 ). The maximum deviation between 'equivalent' bond lengths and angles in the two systems is between the

[^25] $S(1) S(2) C(8)$ planes is 2.00 .
$C(2)-C(3)$ and $C(9)-C(10)$ distances, $(1.460 \AA$ and $1.436 \AA$ respectively), representing about $2 \sigma_{A}$. The dimensions of the conjugated carbonyl systems are discussed in section 14.4. The C(3)-C(4) and $C(10)-C(11)$ bond lengths, $1.525 \AA$ and $1.529 \AA$ respectively, are not significantly different and are consistent with the distance normally associated with a $C\left(\right.$ spe $\left.^{3}\right)-\mathrm{CO}$ bond (86). Fortunately the dimensions of the t-butyl groups were of minimal interest in this study; the high thermal motion of the methyl groups is reflected in the large individual deviation in bond lengths and angles from the accepted values. However the mean $\mathrm{C}-\mathrm{Me}$ distance and $\mathrm{Me}-\mathrm{C}-\mathrm{Me}$ angle, $1.545 \AA$ and $109.5^{\circ}$ at $\mathrm{C}(4)$, $1.528 \AA$ and $109.5^{\circ}$ at C(11), are reasonably consistent with the normal values.

Least squares planes were calculated for,

which showed, (Table 14.3 ), that each atomic grouping is approximately coplanar. The maximum atomic displacement from the plane and the mean atomic deviation are $0.035 \AA, 0.023 \AA$ for plane C; $0.046 \AA, 0.029 \AA$ for plane D. However the central portion of the molecule is strictly non-planar; planes $\underline{C}$ and $\underline{D}$ intersect at an angle of $10.1^{\circ}$, and $\mathrm{S}(3)$ is significantly displaced from each plane, ( $0.118 \AA \AA$ from plane $\underline{C}$ and $0.140 \AA$ from plane $\underline{D}$ ). The

| Atom 'A' | Atom 'B' | Unit translation on Atom 'B' | Distance ( ${ }^{\circ}$ ) |
| :---: | :---: | :---: | :---: |
| S(1) | $S(3){ }_{1}$ | $(1,0,0)$ | 3.45 |
| S(2) | $\mathrm{S}(3)$, | $(1,0,0)$ | 3.49 |
| S(2) | $\mathrm{S}(3)_{2}$ | $(0,0,0)$ | 3.93 |
| $s(3)$ | $\mathrm{O}(2)_{2}$ | $(0,0,0)$ | 3.88 |
| S(2) | $\mathrm{Me}(5)_{4}$ | $(1,-1,0)$ | 3.67 |
| S(2) | $\mathrm{C}(8)_{2}$ | $(0,0,0)$ | 3.74 |
| S(2) | C(9) ${ }_{1}$ | $(1,0,0)$ | 3.98 |
| S(2) | $\mathrm{C}(9){ }_{2}$ | $(0,0,0)$ | 3.98 |
| S(3) | $\mathrm{C}(9){ }_{2}$ | $(-1,0,0)$ | 3.95 |
| O(1) | C(2), | ( $1,0,0)$ | 3.99 |
| O(1) | Me (7), | ( $1,0,0)$ | 3.75 |
| O(1) | $\mathrm{Me}(12)_{2}$ | $(0,0,0)$ | 3.90 |
| O(2) | $\mathrm{C}(1)_{2}$ | $(0,0,0)$ | 3.93 |
| $0(2)$ | $\mathrm{Me}(12)$, | $(1,0,0)$ | 3.72 |
| O(2) | $\mathrm{Me}(13)$, | $(1,0,0)$ | 3.94 |
| c(1) | $\mathrm{Me}(6)_{4}$ | $(0,-1,0)$ | 3,77 |
| Me (6) | $\mathrm{Me}(12)_{2}$ | $(-1,0,0)$ | 3.79 |
| Me (7) | $\mathrm{Me}(14)_{3}$ | $(0,-1,0)$ | 3.54 |
| c(8) | $\mathrm{C}(8){ }_{2}$ | $(0,0,0)$ | 3.99 |

Atom ' B ' is generated from the coordinates listed in Table 13.2 using the symmetry operation ( $)_{i}$, and unit translations shown.
$l_{1}: x, \quad y, \quad z$
$l_{2}:-x,-y, \quad-z$
$j_{3}: x, 1 / 2-y, 1 / 2+z$
$)_{4}:-x, 1 / 2+y, 1 / 2-z$

Trithiacyclopentane from pinacolone:

Closest contact of neighbouring molecules.

Table 14.34

intramolecular sulphur-oxygen distances of $2.509 \AA$ and $2.517 \AA$ are the smallest observed in the three structure analyses; the mean S--0 distance of $2.513 \AA$ being $0.54 \AA$ less than the sum of the van der Waals radii of sulphur and oxygen. The approximate coplanarity of the $\alpha, \beta$-unsaturated carbonyl and sulphur subsystem, and the short sulphur-oxygen distances, are discussed in the following section.

Intermolecular interatomic distances of less than $4.0 \AA$ are listed in Table $14.3_{4}$, and the molecular packing is illustrated in Figure 14.32. Two unusual contact distances are the $\mathrm{Me}(7)-\mathrm{Me}(14)$ and $\mathrm{Me}(6)-\mathrm{Me}(12)$ separations of $3.54 \AA$ and $3.79 \AA$ respectively. Each of these distances is shorter than the separation of $3.90 \AA$ often found between methyl groups (75), but they are greater than twice the van der Waals radius of a carbon atom. Similar observations of unusually short contact distances between two methyl groups have been reported in other crystal structures $(88,89)$ and a careful study did not suggest any alternative positions for the (methyl) atoms in this analysis.

### 14.4 A comparison of molecular structures.

The dimensions of the $\underline{\beta}$-dithia- $\alpha, \beta$-unsaturated carbonyl systems in the three molecules are shown collectively in Figure 14.41. Since the dimensions of the two $\alpha, \beta$-unsaturated carbonyl systems of the l,2,4-trithiacyclopentane from pinacolone are not significantly different, (section 14.3 ), the values shown

[^26]

Desaurin from acetophenone.


3-phenyl-l-propene-1,3-dione 1-(dimethyl mercaptole). PDM


1,2,4-trithiacyclopentane from pinacolone, (mean dimensions).
in Figure 14.41 are the mean dimensions of the two systems. The $C\left(\underline{p p}^{2}\right)-S(I I)$ bond lengths in the three molecules are,

Desaurin
$=1.766 \AA, 1.764 \AA( \pm 0.006 \AA)$; mean $1.765 \AA^{\circ}$,
PDM
: $1.745 \AA \AA^{\circ}, 1.748 \AA \AA^{\circ}( \pm 0.011 \AA)$; mean $1.747 \AA$,
 from pinacolone $\quad 1.748 \AA, 1.739 \AA$

Data so far accumulated indicate that the length of a carbonsulphur single bond is $1.82 \AA$, whilst bond-order, bond-length relationships suggest a value of $2.61 \AA$ for a pure double bond $(83,90)$. The carbon-sulphur distances found are significantly less than the single bond length, indicating that there is some interaction between the electron in the p-pi orbital of $C(1)$ and the lone-pair electrons of the sulphur atom. The mean $C(1)-S$ bond length in the desaurin is greater than the corresponding mean lengths in PDM and the trithiacyclopentane from pinacolone. Although the difference is not highly significant, $\left(\sim 2 \sigma_{A}\right)$, it appears possible that the $C(1)-S$ bonds in the desaurin are lengthened to relieve some steric strain in the 1,3-dithiete ring in an analogous manner to the lengthening of carbon-carbon bonds in cyclobutane systems (91,92).

The bond lengths of the conjugated carbonyl systems in the three molecules are:

[^27]Desaurin

| $c(1)-c(2)$ | $c(2)-c(3)$ | $c(3)-0$ |  |
| :---: | :---: | :---: | :---: |
| 1.324 | $0.008 \AA$ | 1.462 | $0.009 \AA$ |
| 1.378 | $0.013 \AA$ | 1.216 | $0.007 \AA$ |
| 1.330 | $0.008 \AA$ | $0.014 \AA$ | 1.246 |
| $1.0 .013 \AA$ |  |  |  |
| 1.345 | $0.007 \AA$ | 1.460 | $0.007 \AA$ |
| 1.436 | $0.009 \AA$ | 1.221 | 1.233 |
|  | $0.008 \AA$ |  |  |

Application of Cruickshank's significance tests to each 'set' of bond lengths for the system, reveals that the only possibly significant difference is in the $C(1)-C(2)$ distance. This bond in PDM is $0.054 \AA$ longer than in the desaurin, $\left(\sim 3.7 \sigma_{A}\right)$, and $0.040 \AA$ longer than the mean in the trithiacyclopentane from pinacolone, ( $\sim 2.7 \sigma_{A}$ ). Simple carbon-carbon double bond lengths are normally in the range $1.31-1.34 \AA$, being lengthened in a $C=C-C=0$ system by conjugation to around $1.36 \AA(86)$. The $C(1)-C(2)$ bond lengths found therefore suggest that the electrons are more localised in this bond of the desaurin and 1,2,4-trithiacyclopentane from pinacolone molecules. The $C(2)-C(3)$ bond lengths in the three molecules are not significantly different and are consistent with the distance of $1.44 \AA$ normally associated with this conjugated system (86). Values ranging from $1.17 \AA$ to $1.20 \AA$ are suggested for the length of a carbon-oxygen double bond (83), whilst in ketones and conjugated systems the carbonyl bond length normally lies in the range $1.22-1.24 \AA(86)$. The C(3)-0 distances found in the three structure analyses are not significantly different, the maximum difference is $0.030 \AA$ between
$\neq$ Bond lengths in the analogous $C(8)=C(9)-C(10)=0(2)$ system.
the values in PDM and the desaurin, $\left(\sim 2 \sigma_{A}\right)$. When the three bond lengths are considered overall, the distances found in PDM suggest.a greater delocalisation of electrons in the system than in the desaurin and the 1,2,4-trithiacyclopentane from pinacolone.

The valence angles at $C(2)$ and $C(3)$ in the three molecules, (Figure 14.41 ), are reasonably consistent with the value of $120^{\circ}$ normally associated with an $\underline{s p}^{2}$-hybridised carbon atom. The magnitudes of the exocyclic angles ${ }^{\neq}$at $C(1)$ are, of course, governed by the S-C(1)-S angle. In the desaurin this latter angle is considerably reduced from the normal value to $98.0^{\circ}$ in the l,3-dithiete ring and consequently the exocyclic angles in this molecule are significantly greater than those in PIM and the trithiacyclopentane from pinacolone.

The 2-methylene-1,3-dithiacyclobutane derivative, XIII (14), and 2-desylidene-1,3-dithiolane, XIV $(15,16)$, have some similarity to the structures reported in this thesis. Dimensional details of the comparative portions of XIII and XIV are shown in Figure 14.42.

The 2-methylene-1,3-dithiacyclobutane system in XIII is reported to be approximately planar and is therefore comparable with the planar l,3-dithiete ring of the desaurin. Two notable features of XIII are the 'non-equivalence' of the C(spis)-S bond lengths, $1.80 \AA$ and $1.72 \AA$, (contrasting with $1.766 \AA$ and $1.764 \AA$ in the desaurin), and the extremely long $\mathrm{S} 2-\mathrm{C}\left(\mathrm{sp}^{3}\right)$ length of

[^28]

XIII


## XIV

Dimensional details of the comparative portions of the 2 -methylene-1,3-dithiacyclobutane derivative XIII, and 2-desylidene-1,3dithiolane XIV.

Figure 14.42
1.90A. The endocyclic valence angles in the desaurin, $98.0^{\circ}$ at $C\left(\underline{s p}^{2}\right)$ and $82.0^{\circ}$ at $S(I I)$, are not significantly different from the analogous angles in XIII, $100.6^{\circ}$ and the mean value of $83.5^{\circ}$ respectively; the slightly larger values of these angles in XIII are coupled with the small valence angle of 92.30 at the $\underline{s p}^{3}$-hybridised carbon atom $C(4)$.

The $C\left(\right.$ sp $\left.^{2}\right)-S$ bond lengths in XIV, $1.72 \AA$ and $1.76 \AA$, are consistent with the corresponding distances in PDM, (mean 1.747 ${ }^{\circ}$ ), and the trithiacyclopentane from pinacolone, (mean $1.740 \AA$ ). The S(1)-C(1)-S(2) angles in XIV and PDM, (115.50 and $115.1^{\circ}$ respectively), are comparable, although the valence angles at the sulphur atoms in XIV, $94.6^{\circ}$ and $96.7^{\circ}$, reflect the severe puckering in the l,3-dithiolane system and are significantly smaller than the analogous angles in PDM of $104.3^{\circ}$ and $104.4^{\circ}$. The $C\left(\underline{s p}^{3}\right)$-S bond lengths in XIV, $1.83 \AA$ and $1.80 \AA$, and PDM, $1.843 \AA$ and $1.823 \AA$, are in quite good agreement.

Bond lengths in the $C=C-C=O$ systems of XIII and XIV indicate pronounced conjugation and overall appear more consistent with the distances found in PDM than those in the desaurin and trithiacyclopentane from pinacolone molecules.

Attention has already been drawn in the foregoing sections to the approximate planarity of the (s-cis) $\alpha, \beta$-unsaturated carbonyl and sulphur atom sub-systems, in the desaurin, PDM and the 1,2,4-trithiacyclopentane from pinacolone molecules, which incur short intramolecular sulphur-oxygen distances. Similar observations in other molecules have been described in Chapter 1 ,
for which partial sulphur-oxygen bonding has been proposed (11, 12,13).

In the 2-methylene-1,3-dithiacyclobutane derivative XIII, (Figure 14.42 ), it is reported that the four-membered sulphurcontaining ring and the $\alpha, \beta$-unsaturated carbonyl system are approximately coplanar, (with a mean deviation of $0.04 \AA^{\circ}$ ), as is found for the central ten atom system of the desaurin from acetophenone, (mean deviation $0.006 \AA$ ). It is interesting to note that the distortions of the exocyclic angles at $C(1)$ found in the desaurin, (Figure 14.41 ), are apparent to some extent in XIII, (Figure 14.42). The intramolecular suiphur-oxygen distances, $2.640 \AA$ in the desaurin and $2.63 \AA$ in XIII, are some $0.41 \AA$ less than the sum of the van der Waals radii of sulphur and oxygen. Comparison of the dimensions of the $\underline{\beta}$-dithia- $\alpha, \beta-$ unsaturated carbonyl systems in the desaurin and XIII, suggest that the extremely close similarity in sulphur-oxygen distance is probably fortuitous.

Least squares plane data for the $\alpha, \beta$-unsaturated carbonyl and sulphur sub-system in PDM,

has been given in section 14.2 , (Table 14.23, plane D), and showed that this group of six atoms is approximately coplanar, (with a mean atomic deviation of $0.017 \AA$ ). As previously stated,
in 2-desylidene-1,3-dithiolane XIV (16), (Figure 14.42), it is reported that although the analogous $S(1)-C(1)=C(2)-C(3)=0$ grouping is approximately coplanar, (mean deviation $0.034 \AA$ ), $S(2)$ is significantly out-of-plane by $0.24 \AA$. The findings in the PDM molecule are in direct contrast to those in XIV on this point; $S(2)$ in PDM is closely coplanar with the five atom system, (Table $14.2_{3}$, plane C). The distortions in exocyclic angles at $C(1)$, found in the desaurin, are not apparent in the PDM molecule, (Figure 14.41); however in XIV a severe 'reverse' distortion is evident, (Figure 14.42). The intramolecular sulphur-oxygen distances in PDM, (2.727 ${ }^{\circ}$ ), and XIV, (2.70 $)$, are closely similar, but it is unwise to take a distance of around $2.7 \AA$ as characteristic for the grouping on the basis of two determinations.

In the l,2,4-trithiacyclopentane from pinacolone molecule, least squares planes were calculated for,

which showed, (Table $14.3_{3}$ ), that each atomic grouping is approximately coplanar. Although the mean atomic deviations, ( $0.023 \AA$ for plane $\underline{C}$ and $0.029 \AA$ for plane $D$ ), are greater than those found in the desaurin and PDPI, the intramolecular sulphuroxygen distances found are the shortest in the three analyses.

The unusually short S--O distances found in this molecule, which involve the sulphur atoms of a five-membered heterocyclic disulphide system, are similar to those found in 2,5-dimethyldithiofurophthene, X (11), and 3,5-diacetamido-1,2-dithiolium bromide, XIX (93).

$x^{\neq}$


XIX

The molecular skeleton of 2,5-dimethyl-dithiofurophthene is reported to be essentially planar, (with a maximum atomic deviation of $0.04 \AA$ ), incuring an extremely short intramolecular sulphur-oxygen distance of $2.44 \AA$, with a disulphide bond length of 2.12 $\AA$. Mammi et. al. (11), after comparing the molecular dimensions of $X$ with those of the previously reported isologue 2,5-dimethyl-thiothiophthene XI (12), interpreted the short S-O distance as evidence of partial bonding between these atoms.

In 3,5-diacetamido-1,2-dithiolium bromide, the 1,2dithiolium ring is found to be planar within experimental error. Although the nitrogen atoms lie close to this plane, the exocyclic carbon atoms, $C_{4}$ and $C_{6}$, and the carbonyl oxyeen atoms, $\mathrm{O}_{1}$ and $\mathrm{O}_{2}$, are out-of-plane by $+0.091 \mathrm{~A},+0.094 \mathrm{~A},+0.048 \mathrm{~A}$ and +O.130A respectively. Least squares planes through the $S_{1}-C_{1}$ -$\mathrm{N}_{1}-\mathrm{C}_{4}-\mathrm{O}_{1}$ and $\mathrm{S}_{2}-\mathrm{C}_{2}-\mathrm{N}_{2}-\mathrm{C}_{6}-\mathrm{O}_{2}$ systems of XIX would therefore appear to be akin to the 'butterfly-wing' disposition of planes

C and $\underline{D}$ about the heterocycle in the 1,2,4-trithiacyclopentane from pinacolone molecule. The intramolecular sulphur-oxygen distances in XIX, $2.571 \AA$ for $S_{1}-0_{1}$, and $2.515 \AA$ for $S_{2}-O_{2}$, are comparable with those found in trithiacyclopentane from pinacolone, $(2.517 \AA$ and $2.509 \AA$ ). Hordvik (93) suggests that the slight asymmetry of the sulphur-oxygen distances in XIX is due to the oxygen atom 0 , entering into a weaker contact, ( $3.113 \AA$ ), with a second $S_{1}$ atom of a symmetry related ion, and concludes that the short $\mathrm{S}-\mathrm{O}$ distances are indicative of partial bonding between these atoms. By taking the sum of the covalent radii of sulphur and oxygen as $1.71 \AA$, the sum of the van der Waals radii as $3.25 \AA^{\circ}$, and assuming a linear bondorder, bond-length relationship, the sulphur-oxygen contacts of $2.571 \AA$ and $2.515 \AA$ were taken as corresponding to sigma bond-orders of 0.47 and 0.44 respectively. By comparing the disulphide bond length in XIX, (2.080 ${ }^{\circ}$ ), with those found in thiuret hydrobromide, $2.081 \AA$ (94), and thiuret hydrochloride hemihydrate, $2.071 \AA(95,96)$, Hordvik concludes that the partial bonding between the sulphur and oxygen atoms does not appreciably affect the disulphide bond length. Whether the close S--O contacts in the trithiacyclopentane from pinacolone effectively lengthen the disulphide bond is not at all easy to detect without further structural data on 1,2,4-trithiacyclopentane derivatives. In this context a structure analysis

[^29]of the second cis isomer type,

would unequivocally test this point.
14.5 Conclusion.

The structures of the desaurin and $1,2,4$-trithiacyclopentane are confirmed and dimensional details of these novel heterocyclic systems are now established. The structure analyses reveal that in each molecule the $\alpha, \beta$-unsaturated carbonyl systems adopt an s-cis conformation and the $\mathrm{S}-\mathrm{C}=\mathrm{C}-\mathrm{C}=0$ systems are approximately planar, incuring intramolecular sulphur-oxygen distances which are considerably less than the sum of the van der Waals radii of sulphur and oxygen. In the light of similar observations in related systems, and the conclusions drawn therefrom, the short sulphur-óxygen distances found in this work could perhaps be considered as evidence for partial bonding between the sulphur and oxygen atoms, or a strong charge interaction.

In 2,5-dimethyl-thiothiophthene XI, Giacometti and
Rigatti (97) considered that the sulphur atoms use only their 3p orbitals for bonding, and viewed the sulphur-sulphur bonds as having both sigma and pi contributions. Mammi et al. (1l) considered the partial bonding between the sulphur and oxygen atoms in the related 2,5-dimethyl-dithiofurophthene $X$ to be of
the same general nature, as did Hordvik in 3,5-diacetamido-1,2dithiolium bromide, XIX (93). A reasonable mechanism by which partial bonding between sulphur and oxygen could occur in the three molecules studied in this work, is through overlap of a 2 p orbital on oxygen with a suitably orientated 3d orbital on sulphur. If the $2 p_{x}$ and $2 p_{z}$ orbitals of oxygen are involved in bonding to carbon, the $2 p_{y}$ orbital remains for bonding to the sulphur. This orbital lies in the plane of the $\alpha, \beta$-unsaturated carbonyl system and is directed towards the sulphur atom. Consequently overlap between this orbital and a 3d orbital on the sulphur may give rise to a sigma type bond. An alternative formulation involves p-pi-d-pi bonding between the oxygen $2 p_{z}$ orbital and a sulphur 3d orbital.

## APPENDIX I

The desaurin from acetophenone.

Observed and calculated structure factors.

Unobserved reflections are marked *


M 1 10\% 10F


APPENDIX II.

3-Phenyl-1-propene-1,3-dione 1-(dimethyl mercaptole).

A summary of the symbolic addition procedure.

* denotes reflections for which the phase generated had both absolute and symbolic contributors, indicating $A=$ - .
*** denotes reflections which were not used as input for cycle 13.
*** denotes the initial symbolic assignment $A$.

For cycles 1 through 12, $|E| \geqslant 2.0$.
For cycle $13,2.0\rangle|E| \geqslant 1.5$.
cycle $h \quad k \quad l|E| \quad s \quad \mid$ cycle $h \quad k \quad l|E| \quad s$ no.

|  | 2 | 2 | -9 | 2.63 | + |
| :--- | ---: | ---: | ---: | ---: | ---: |
|  | 5 | 1 | -2 | 2.83 | + |
|  | 2 | 3 | 1 | 2.90 | + |
| 1. | 4 | 1 | -8 | 2.38 | - |
|  | 0 | 1 | 10 | 2.37 | + |


cyole $h \quad k \quad l \quad|E| \quad s$ no.
9. $1 \begin{array}{lllll}1 & 4 & -7 & 2.15 & -\mathrm{A}\end{array}$ cont. $128 \quad 1 \quad 2.04+A$

|  | 1 | 8 | 11 | 2.10 | $-A$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| 10. | 2 | 7 | 6 | 2.04 | $-A$ |


|  | 0 | 2 | 15 | 2.56 | - |
| :---: | :---: | :---: | :---: | :---: | :--- |
| 0 | 8 | 5 | 2.32 | +A |  |
|  | 1 | 3 | 10 | 2.41 | - |
|  | 3 | 2 | 11 | 2.26 | + |
| 11. | 2 | 1 | 16 | 2.11 | + |
|  | 2 | 2 | 3 | 2.25 | - |
|  | 3 | 2 | 9 | 2.12 | + |
|  | 0 | 5 | 15 | 2.28 | -A |


| $*$ | 4 | 3 | 6 | 2.12 | + |
| :--- | ---: | ---: | ---: | ---: | ---: |
| $*$ | 0 | 9 | 6 | 2.09 | - |
| $*$ | 0 | 4 | 14 | 3.04 | + |

** $0 \quad 12 \quad 0 \quad 2.40$ +
** $0 \begin{array}{lllll}\text { * } & 10 & 8 & 2.26 & +A\end{array}$ * 1 1 82.18 *. $3 \quad 9 \quad 2 \quad 2.01+$
13. * $2.5 .14 \quad 1.94$ -

0 3-10 $1.94+$
$2 \quad 9-10 \quad 1.90+\mathrm{A}$
$25 \quad 6 \quad 1.89+A$
0341.88 +

4 9. 21.85 tA
$\begin{array}{lllll}2 & 9 & -8 & 1.84 & +A\end{array}$
$\begin{array}{lllll}2 & 3 & -8 & 1.84 & -A\end{array}$
2. 7 - $\quad-6$ 1.81 +A

6 1 2. 1.78 -
2. $5-10 \quad 1.75-\mathrm{A}$ * $\begin{array}{lllll} & 3 & 8 & 1.74\end{array}+$
h k $\quad|E| \quad s$


* 21151.59 +
$\begin{array}{lllll}2 & 7 & 9 & 1.58 & +A\end{array}$
* 293 1.51 +
$2 \quad 4 \quad 12 \quad 1.97$ -
210 . $0 \quad 1.84+A$
* 4 4-12 1.80 +
$2014 \quad 1.75$ +
$\begin{array}{lllll}4 & 6 & 8 & 1.74 & -A\end{array}$
2 2 0. 1.67 +A
$4 \quad 4 \quad 0 \quad 1.60$ -
$0.8111 .99+A$
$\begin{array}{lllll}4 & 8 & -5 & 1.97 & -A\end{array}$
$22111.93-$
$\begin{array}{lllll}0 & 8 & 7 & 1.86 & +A\end{array}$
$6 \quad 2 \quad-9 \quad 1.83 \quad$ +
$04111.73-A$



#### Abstract

APPENDIX III.

3-Phenyl-1-propene-1,3-dione 1-(dimethyl mercaptole).


 Observed and calculated structure factors.Reflections classified as unobserved are marked * .



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## APPENDIX IV.

The 1,2,4-trithiacyclopentane from pinacolone.

## Observed and calculated structure factors.

Reflections classified as unobserved are marked *




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## APPENDIX V.

## Crystallographic programs.

(a) Interpolation program for scattering factor
curves.
(b) Lorentz-polarisation correction program for
equi-inclination Weissenberg photographs. xix.
(c) A program to generate normalised structure
factor magnitudes.

INTERPOLATION PROGRAM - INPUT IS F VALUES AT SIN THETA/LAM. 0.0 TO 1.10 IN STEPS OF 0.05 ( 23 OF 1. ..OUTPUT IS F VALUES FOR SIN THETA FROM 0.02 TO 1.00 IN STEPS /F 0.02 ( 50 OF)..
DATA IS 1. SINTHETA/LAM. FVALUE.. 23 OF IN FORMAT F6.4.F5.2..
SINTHETA/LAM RUNS FROM 0.00 TO 1.10..
DATA2. MM(FORMATI3). MM IS THE NO. OF VALUES OF SINTHETA
TO BE INTERPOLATED..
DATA 3.. MM VALUES OF SINTHETA IN FORMAT F6.4...
103 FORMAT ( $1 \times, 28 H$ SINT RHO SCATT)
102 FORMAT(13)
112 FORMAT(F6.4)
111 FORMAT(1X,F7.4,3X,F7.4,3X,F8.4)
100 FORMAT(F6.4,F5.2)
DIMENSION SN(23),F(23),B(4),C(3),D(2)
DO $11=1,23$
$1 \operatorname{READ}(5,100)$ SN(I),F(I)
READ (5,102) MM
WRITE(6,103)
$K K=0$
$420 K K=K K+1$
IF(KK-MM)421,421,422
$421 \operatorname{READ}(5,112) \mathrm{X}$
$S X=(X / 1.5418)$
$150 A=0.0500$
$200 A=A+0.0500$
IF(A-SX)200,201,202
$202 \mathrm{BI}=(\mathrm{A} / 0.050 \mathrm{C})+1.0000$
$I=B I$
$J=1-3$
DO $300 K=1,4$
$J=J+1$
$300 B(K)=(F(J) *(S N(J+1)-S X)-F(J+1) *(S N(J)-S X)) /(S N(J+1)-S N(J))$
$J=I-3$
$K=0$
DO $301 \mathrm{~L}=1.3$
$J=J+1$
$K=K+1$
$301 C(L)=(B(K) *(S N(J+2)-S X)-B(K+1) *(S N(J)-S X)) /(S N(J+2)-S N(J))$
$J=1-3$
$L=0$
DO $302 \mathrm{M}=1,2$
$L=L+1$
$J=J+1$
$302 D(M)=(C(L) *(S N(J+3)-S X)-C(L+1) *(S N(J)-S X)) /(S N(J+3)-S N(J))$ $J=1-2$
303 FCAL=(D(1)*(SN(J+4)-SX)-D(2)*(SN(J)-SX))/(SN(J+4)-SN(J))
304 WRITE $(6,111) X, S X, F C A L$
GO TO 420
$201 \mathrm{BI}=(\mathrm{A} / 0.0500)+1.0000$
$1=B I$
WRITE(6,111) X,SX,FII)
GO TO 420
. 422 RETURN
END
L.P. CORRECTION FOR EQUI-INCLINATION WEISSENBERGS

IF ITEST IS 0 PRINT, IF $=1$ THEN PUNCH, IF $=2$ THEN PRINT AND PUNC ********************
** VERSION 103 **

1 FORMAT (3F7.3,4F7.4)
2 FORMAT(A5,4F8.4)
3 FORMAT (313,F9.2,F5.3,13)
4 FORMAT (13)
5 FORMAT(1H\$,3I9,F9.2,9X,F9.0,15)
6 FORMAT $1,1 \times, 6 H$
,F7.4,F7.2,314,F9.2,F8.3,F9.21
7 FORMAT $11 H^{\circ}, 58 \mathrm{H}$ SINT FOBS $H \quad K \quad L \quad O B S I \quad$ DLP 1 1*D !
8 FORMAT $11 \mathrm{HI}, 15 \mathrm{H}$ LAYER CODE IS A5,7H,MU IS F8.4.8H DEGREES,17H,SCALE
1 FACTOR IS F8.4)
9 FORMAT (IX.31H MINIMUM INTENSITY MEASURED IS F8.4,9H,ZETA IS F8.4)
10 FORMAT $11 X, 6 H^{* * * *, F 7.4, F 7.2,314, F 9.2, F 8.3, F 9.21) ~}$
11 FORMAT (1HO,12H REFLECTION 3I3.32H HAS SIN(THETA) GREATER THAN 1.0) DIMENSION RCP(6),PRCP(6)
READ (5,4)ITEST
READ (5,1). A, B, C, CAA,CBB,CCC, WAVE
$D=S Q R T(1 \cdot+2 \bullet * C A A * C B B * C C C-C A A * C A A-C B S * C B B-C C C * C C C)$
$D=1 . /(D * A * B * C)$
SAA $=$ SQRT(1.-CAA*CAA)
$S B B=S Q R T(1 .-C B B * C B B)$
SCC=SQRT(1.-CCC*CCC)
$R C P(1)=D * B * C * S A A$
$R C P(2)=D * A * C * S B B$
$R C P(3)=D * A * B * S C C$
$R C P(4)=(C B B * C C C-C A A) /(S B B * S C C)$
$R C P(5)=(C A A * C C C-C B S) /(S A A * S C C)$
$R C P(6)=(C A A * C S B-C C C) /(S A A * S B B)$
$\operatorname{PRCP}(1)=R C P(1) * R C P(1) * 0.25$
$\operatorname{PRCP}(2)=R C P(2) * R C P(2) * 0.25$
$\operatorname{PRCP}(3)=R C P(3) * R C P(3) * 0.25$
$\operatorname{PRCP}(4)=R C P(2) * R C P(2) * R C P(6) * 0.25$
$\operatorname{PRCP}(5)=R C P(1) * R C P(3) * R C P(5) * 0.25$
$\operatorname{PRCP}(6)=R C P(2) * R C P(3) * R C P(4) * 0.25$
100 CONTINUE
READ $(5,2)$ ICODE, AMU,SCALE,ZMIN],ETA
WRITE $(6,8)$ ICODE,AMU,SCALE
WRITE 6,9$)$ ZMINI,ETA
IF(ITEST-1) 120,121,120
120 CONTINUE
WRITE (6,7)
121 CONTINUE
112 CONTINUE
READ (5,3) IH,IK,IL,OBSI,Q,ID
IF(IH-999) 102,103,102
103 IF(IK-999) $100,104,100$
$102 \mathrm{FH}=\mathrm{IH}$
$F K=I K$
FL=IL

```
    RHO=FH*FH*PRCP(1)+FK*FK*PRCP(2)+FL*FL*PRCP(3)+2.*FH*FK*PRCP(4)+2.*
        1FK*FL*PRCP(6)+2.*FL*FH*PRCP(5)
        SQSINT = RHO*NAVE*WAVE
        SINT=SQRT(SQSINT)
        IF(SINT-1.) 400,400,401
    401 CONTINUE
        WRITE(6,11) IH,IK,IL
        GO TO 112
    400 CONTINUE
    CENT=SQRT(1.-(SINT*SINT))
    XISQ=(4.*SOSINT)-(ETA*ETA)
    XI=SQRT(XISQ)
    USING THE EQUATION GIVEN IN COGHRAN.
    D=(XI *CENT)/(1+COSSQ(TWOTHETA))
    CTWTH=(1.-(SINT*SINT*2.))
    CSQTW=CTWTH* CTWTH
    DLP=(XI *CENT)/(1.+CSQTW)
    IF(OBSI) 200.201.200
    201 CALCID=ZMINI*DLP*SCALE
    1D=1
    GO TO 202
    200 CALCID=OBSI*DLP*SCALE
    202 IH=FH
    IK=FK
    IL=FL
    FOBS=SQRT(CALCID)
    402 CONTINUE
        IF(ITEST-1) 206,207,206
    PRINT OUT
    206 IF(ID) 209,208,209
    208 WRITE(6,6) SINT,FOBS,IH,IK,IL,OBSI,DLP,CALCID
    GO TO 250
    209 WRITE(6,10).SINT,FOBS,IH,IK,IL,OBSI,DLP,CALCID
    250 IF(ITEST-1) 112,112,207
    207 IF(ID) 220,221,220
    220 FOBS=0.
    221 CONTINUE
        WRITE(6,5) IH,IK,IL,FOBS,Q,ID
        GO TO 112
    104 RETURN
    END
SDATA
```


# A PROGRAM TO GERERATE NORMALISED STRUCTURE FACTORS SUBROUTINE SYSTAB SETS UP THE CORRECT EPSILON VALUES 

1 FORMAT (3F7.3,4F7.4)
2 FORMAT (7F10.4)
3 FORMAT (1HO, $32 H$ RECIPROCAL CELL PARAMETERS ARE )
4 FORMAT (1X,6F10.5)
5 FORMAT (13)
6 FORMAT ( F9.6)
7 FORMAT(319,2F9.2,59.0.15)
8 FORMAT(1X.3I5,F9.4,F9.3,5F7.3.2F9.3,F9.4)
9 FORMAT(1HS,314,F9.4)
10 FORMAT (1HO,95H H K L SINT OBSI F1 F2 F3
1 F4 F5 SIG2 EP*SIG2 E**2 )
12 FORMAT(1HO,17H SCALE FACTOR $=$ F9.4///)
13 FORMAT(IX,6H Z IS F9.4,53H, NO OF ATOMS OF THIS TYPE IN THE IUNIT CELL IS 13//)
14 FORMAT (1HO,60H $A^{*} B^{*} C^{*} \operatorname{COS}(A L *) \operatorname{COS}(B E *)$ 1COS(GA*))
15 FORMAT(1013)
16 FORMAT ( $1 \mathrm{HO}, 25 \mathrm{X}, 23 \mathrm{H}$ NUMBER OF F CURVES IS 15 )
17 FORMAT ( $1 \mathrm{HO}, 25 \mathrm{x}, 27 \mathrm{H}$ DATA SET UP IS ON LOGICAL I5)
18 FORMAT( $1 \mathrm{HO}, 25 \mathrm{X}, 34 \mathrm{H}$ NO INITIAL PRINT OUT IS REQUESTED)
19 FORMAT(1HO,25X,34H INITIAL PRINT OUT IS REQUESTED)
220 FORMAT (1HO,25X,29H A K GENERATION IS REQUIRED)
221 FORMAT(1HO,25X,29H AN E GENERATION IS REQUIRED)
DIMENSION RCP(6), PRCP(6),AFM(11,51),FM(15),NOAT(10),CONT(10)
tAPE ALLOCATIONS.
NN IS BINY TAPE FOR EXPANDED DATA AND MM IS TAPE USED IN KGEN FOR STORAGE AND LSQFIT PICK-UP.

```
NFOUR=2
NN=1
MM=4
READ(5,1) A,B,C,CAA,CBB,CCC,WAVE
D=SQRT(1.+2.*CAA*CBB*CCC-CAA*CAA-CBB*CBB-CCC*CCC)
D=1./(D*A*B*C)
SAA=SQRT(1.-(AA*CAA)
SBB=SQRT(1.-CBB*CBS)
SCC=SQRT(1.-CCC*CCC)
RCP(1)=D*B*C*SAA
RCP(2)=D*A*C*SBB
RCP(3)=D*A*B*SCC
RCP(4)=(CBB*CCC-CAA)/(SBB*SCC)
RCP(5)=(CAA*CCC-CBB)/(SAA*SCC)
RCP(6)=(CAA*CBS-CCC)/(SAA*SBS)
PRCP(1)=RCP(1)*RCP(1)*0.25
PRCP(2)=RCP(2)*RCP(2)*0.25
PRCP(3)=RCP(3)*RCP(3)*0.25
PRCP(4)=RCP(1)*RCP(2)*RCP(6)*0.25
PRCP(5)=RCP(1)*RCP(3)*RCP(5)*0.25
PRCP(6)=RCP(2)*RCP(3)*RCP(4)*0.25
WRITE(6,3)
WRITE(6,14)
```



```
            IF(BRAG)28,28,27
            27 CONTINUE
            28 I=I-1
            FRAC=(BRAG+0.02)*50.
            DO 29 K=1,10
        29 FM(K)=0.
            DO 30 K=1,NATSC
            30 FM(K)=AFM(K,I)+(AFM(K,I+1)-AFM(K,I))*FRAC
            IF(ID.EQ.O) GO TO 999
            OBSI = 0.56
999 CONTINUE
    TOT=O.
    DO 104 L=1,NF
    FNOAT=NOAT(L)
        CONT=FNOAT*FM(L)*FM(L)
    TOT=TOT+CONT
104 CONTINUE
    SIG2 = TOT
    CALL SYSTAB (IH,IK,IL,EPSI)
    ESIG2 = SIG2 * EPSI
    ESQ = (OBSI/ESIG2)
    WRITE(NN) IH,IK,IL,ESIG2,ESQ,OSSI,FOBS,SINT,ID,EPSI
    IF(NOWR.EQ.O) GO TO 444
    WRITE(6,8) IH,IK,IL,SINT,OBSI,(FM(I),I=1,5),SIG2,ESIG2,ESQ
    444 CONTINUE
    GO TO 102
    50 IH = 999
    WRITE(NN) IH,IK,IL,ESIG2,ESQ,OSSI,FOBS,SINT,ID,EPSI
    END FILE NN
    REWIND NN
    IF (KCURV - 1) 777.666,666
666 CALL. KGEN (NN,MM)
    CONTINUE
    GO TO 995
777 CALL EGEN (NN,NFOUR)
    CONTINUE
995 RETURN
    END
```

SIBFTC KKK
SUBROUTINE KGEN（NN，MM）
THIS IS USED TO CALCULATE SUMS FOR GENERATING K CURVE，SO THAT
AN EXPONENTIAL FIT OF THE FORM K＝A＊EXP（B＊SINT＊＊C）
CAN be obtained and this used to place
THE INTENSITIES ON THE TRUE SCALE
DIMENSION S（30），TSIG（30），TIN（30），ZN（30），SM（30），CCK（30）
1 FORMAT（3I3，F9．6）
2 FORMAT（2F9．6）
3 FORMATIIHO，66H RANGE TOT（SIG2）TOTIINT）
IIN（TH）M NO／／I
4 FORMAT（IH ，15，3F15．4，F10．6，F6．1）

5 FORMAT(1X,3I5,2F15.4,15,F9.6,F6.1)
6 FORMAT ( $1 \mathrm{HI}, 50 \mathrm{H} \quad \mathrm{H} \quad \mathrm{L}$ E*SIG2 OBSI RANGE)
7 FORMAT(18H1KCURVE GENERATION//)
8 FORMAT(6HO )
9 FORMAT(2OH K CURVE SET NUMBER 13)
10 FORMAT(12H REFLECTION $313,36 H$ IS OUTSIDE THE MAXIMUM SIN(THETA), 115 H AND IS IGNOREDI
REWIND MM
11 CONTINUE
REWIND NN
READ(5,1) N,III,NSET,SINMAX
WRITE(MM) NONSET
IF(III.EQ.O) GO TO 100
WRITE $(6,6)$
100 CONTINUE
DO $1011=1, N$
TSIG(I) $=0.0$
TIN(I) $=0.0$
$Z N(1)=0.0$
CCK(I) $=0.0$
101 CONTINUE
CALL RANSET(N,SINMAX,S ,SM )
102 CONTINUE
READ (NN) IH,IK,IL,ESIG2,ESQ,OBSI,FOBS,SINT,ID,EPSI
IF(IH.EQ.999) GO TO 999
IF(SINT.LE.S(N)) GO TO 103
C WRITE(6,10) IH,IK,IL
GO TO 102
103 CONTINUE
DO $104 \mathrm{I}=1, \mathrm{~N}$
IF(SINT.GE.S(I)) GO TO 104
TSIG(I) $=$ TSIG(I) + ESIG2
TIN(I) $=$ TIN(I) + OBSI
$Z N(I)=2 N(I)+1.0$
IF. (III.EG.O) GO TO 150
WRITE(6,5) IH,IK,IL,ESIG2,OBSI,I,SINT,EPSI
150 CONTINUE
GO TO 102
104 CONTINUE
$C$
$C$
REACHES 999 AT THE END OF ALL DATA
999 CONTINUE
REWIND NN

PUT OUT THE REQUIRED SUMS
WRITE $(6,3)$
DO 110 I $=1, N$
CCK(I) $=(T S I G(I) / T I N(I))$
WRITE( 6,4 ) I,TSIG(I),TIN(I), CCK(I),SM(I), ZN(I)
WRITE(MM) I,TSIG(I),TIN(I),CCK(I),SM(I),ZN(I),S(I)
110 CONTINUE
IF(NSET.NE.O) GO TO 11

```
FINAL SET OF K CURVE
```

REWIND MM
WRITE(6,7)
WRITE 6,3$)$
$J K=0$
12 CONTINUE
$J K=J K+1$
WRITE(6,9) JK
WRITE 6,8 )
READ(MM) N,NSET
DO $14 \mathrm{I}=1, \mathrm{~N}$
READ (MM I I,TSIG(I),TIN(I),CCK(I),SM(I),ZN(I),S(I)
WRITE(6,4) I,TSIG(I),TIN(I),CCK(I),SM(I),ZN(1)
14 CONTINUE
IF(NSET.NE.O) GO TO 12
CALL LSQFIT(MM)
RETURN
END

SIBFTC ECALC
SUBROUTINE EGEN (NN,NFOUR)
DIMENSION FM(20)
THIS SUB READS BACK FROM BINY TAPE NN, AND SETS UP AVERAGES REQUIRED THE SIX QUANTITES REQUIRED ARE

| E DESIG | NUMBER TOTAL |
| :--- | :--- |
| TOTE | NUM |
| TESQM1 | NUM |
| TOTESQ | NUM |
|  |  |
|  | N3 |
|  | N2 |
|  | N1 |
|  |  |

THIS SUB ROUTINE ALSO OUTPUTS CARDS WITH THE REQUIRED.
FUNCTIONS ON. N.B. DATA IS SCALED USING THE EXP. FUNCTION
GENERATED FROM KGEN
1 FORMAT (1HI,17HTOTAL/E/ $=F 12.4 .26 H$ NUMBER OF OBSERVATIONS $=$
115.23H ***AVERAGE /E/ =F12.4)

2 FORMATI1HO,17HTOTAL /E**2-1/ = F12.4.31X.23H ***AVERAGE/E**2-1/= 1 F12.4)
3 FORMATI1HO,17HTOTAL /E**2/ = F12.4.31X,23H***AVERAGE/E**2 / = 1 F12.4)
4 FORMATIIHO,6H
5 FORMATILHO,26H NUMBER WITH /E/GT $3.0=15,18 \mathrm{H}$ PERCENT GT $3.0=F 8$ 1.4)

6 FORMATIIHO,26H NUMBER WITH /E/GT $2.0=15,18 H$ PERCENT GT $2.0=F 8$ 1.4)

7 FORMATI 1 HO,26H NUMBER WITH /E/GT $1.0=15,18 \mathrm{H}$ PERCENT GT $1.0=\mathrm{F} 8$ 1.41

8 FÓRMAT (1X,315,3F9.3)
9 FORMAT(1H1,42H $\quad \mathrm{H} \quad \mathrm{L} \quad \mathrm{E}$ E*2 (E**2-1))
10 FORMAT (3F10.5,110)
11 FORMAT(1H\$,314,5F9.4)
12 FORMAT ( $1 \mathrm{x}, 314,6 \mathrm{F9} .4$ )
13 FORMAT(1HI,57H H K L E E*2 E**2-1 SINT K ICALC )
14 FORMAT (1HO,5OH THE AVERAGE E**2 IS OUTSIDE THE PERMISSABLE RANGE// $1,35 \mathrm{H}$ AND CONSEQUENTLY THE JOB IS HALTED//,47H IF THE A COEFFICIENT 2 IS REPLACED BY THE VALUE F10.4,28H AN AVERAGE E**2 WILL RESULT//, 326 H OR RE-ANALYSE THE K CURVEI
15 FORMAT(1HO,51H THE K CURVE USED IS OF THE FORM K = $A^{*}(B(S I N T) * * C) /$ 1///,32H THE COEFFICIENTS USED HERE ARE ////. 5H A $=$ Flo.4//.5H B = 2F10.4//.5HC = F10.4//1
16 FORMAT (I3)
17 FORMAT (IHO,45H ONLY THE AVERAGES ARE REQUIRED NO INDIVIOUAL///. $140 H$ LISTING FOR EACH PLANE WILL BE GIVEN ,
$C$ READ IN PARAMETERS OF THE EXPONENTIAL FUNCTION
778 READ (5,10) A,B,C,NAVER
IF(A.GT.100.) GO TO 800
IF(NAVER.EQ.) GO TO 240
WRITE(6,17)
WRITE (6,13)
240 CONTINUE
c
C SET SCK AND NPUN TO 1 FOR INITIAL RUN
SCK $=1.0$
NPUN $=1$
IEND $=0$
$C$
400 CONTINUE
$\mathrm{N} 1=0$
$N 2=0$
N3 $=0$
NUM $=0$
TOTE $=0.0$
TOTESQ $=0.0$
TESQMI $=0.0$
REWIND NN
201 CONTINUE
READ (NN) IH,IK,IL,ESIG2,ESQ,OBSI,FOBS,SINT,ID,EPSI
IF(IH.EQ.999) GO TO 200
C IF(SINT.GT.0.75) GO TO 201
CONK $=A * E X P(B * S I N T * * C)$
$E 2=E S Q$ * CONK * SCK
$E 2 M 1=E 2-1.0$
E = SQRT(E2)
NUM $=$ NUM + 1
TOTE = TOTE + E
IF(E2M1.GE.0.0) GO TO 301
EEMI $=-E 2$ M1
GO TO 300
301 EEM1 = E2M1

```
    300 TESQM1 = TESQM1 + EEM1
        TOTESQ = TOTESQ + E2
C
    PRCENTAGE TOTALS
        IF(E.LT.3.0) GO TO 220.
        N1 = N1 + 1
        N2 =N2 + 1
        N3 = N3 + 1
        GO TO 555
    220 IF(E.LT.2.0) GO TO 210
        N1 = N1 + 1
        N2 = N2 + 1
        GO TO 555
    210 IF(E.LT.1.O) GO TO 555
    Nl = Nl + 1
    555 CONTINUE
        IF(NPUN.NE.O) GO TO 302
        IF(NAVER.NE.O) GO TO 302
        CALL CARD (IH,IK,IL,E,E2,E2MI,SINT,CONK)
        CONTINUE
    302. CONTINUE
    IF (IEND.NE.1) GO TO 201
C
    CALL FOUR2 (IH,IK,IL,E,E2MI,SINT,ID,FOBS,NFOUR,EPSI)
    IF(NAVER.NE.C) GO TO 201
    WRITE(6,12) IH,IK,IL,E,E2,E2MI,SINT,CONK,ESQ
    GO TO 201
    200 CONTINUE
    ZNUM = NUM
    2NI = NI
    ZN2 = N2
    ZN3 = N3
    AVE =TOTE/ZNUM
    AVESMI = TESQM1 / ZNUM
    AVESQ = TOTESQ/ZNUM
    PERE3 = (ZN3/ZNUM) * 100.0
    PERE2 = (ZN2/ZNUM) * 100.0
    PEREI = (ZN1/ZNUM) * 100.0
    WRITE(6,1) TOTE,NUM,AVE
    WRITE(6,2) TESQM1, AVESMI
    WRITE(6,3) TOTESQ,AVESQ
    WRITE(6,4)
    WRITE(6,5) N3, PERE3
    WRITE(6,6) N2, PERE2
    WRITE(6,7) N1, PEREI
    WRITE(6,15) A,B,C
    IF(IEND.EQ.1) GO TO 311
    IF(AVESQ.LT.0.98) GO TO 500
    IF(AVESQ.GT.1.02) GO TO 500
    IEND = 1
    NPUN = 0
    GO TO 400
C
    500 IEND = 0
```

SCK $=(1.0 / A V E S Q)$
$D=S C K * A$
WRITE $(6,14) \quad D$
311 CONTINUE
CALL FOUR2 (IH,IK,IL,E,E2MI,SINT,ID,FOBS,NFOUR)
REWIND NN
GO TO 778
800 RETURN
END

SIBFTC KFIT
SUBROUTINE LSQFIT (MM)
DIMENSION Y(1C0), X(100,2),A(2),W(100),Z(100),ZZ(100)
DIMENSION TSIG(100),TIN(100),CCK(100),SM(100),ZN(100),S(100)
REWIND MM
C READS BACK SINTHETA (SM) AND K (CCK) FROM TAPE MM READY TO GO TO THE LIBRARY SUBROUTINE FOR THE LEAST SQUARES FIT.
J1 $=1$
$J_{2}=0$
101 CONTINUE
READ (MM) N, NSET
$\mathrm{J} 2=\mathrm{J} 2+\mathrm{N}$
DO 104 I = Jl, J2
READ(MM) I,TSIG(I),TIN(I),CCK(I),SM(I),ZN(I),S(I)
104 CONTINUE
$J 1=J 1+N$
IF(NSET.NE.O) GO TO 101
NTOT = $\mathrm{J2}$
WRITE(6,111) NTOT
111 FORMAT (1H1,43H LEAST SQUARES FITTING SUSROUTINE IS CALLED,//36H TH IE TOTAL NUMBER OF POINTS. USED IS I5)

ENTER INTO THE L.S. SUBROUTINE.
THIS PROGRAM USES LEAST SQUARES LIBRARY SUBROUTINE
TO FIT NTOT POINTS $(W, Z)$ TO A CURVE
DO 106 I $=1, N T O T$
ZZ(I) = SM(I)
$W(I)=C C K(I)$
106 CONTINUE
ON = 1 .
100 DO $11=1, N T O T$
$Y(I)=A L O G(W(I))$
$Z(I)=Z Z(I) * * X N$
1 CONTINUE
DO $21=1$, NTOT
DO $2 \mathrm{~J}=1,2$
$2 X(I, J)=2(1) * *(J-1)$
CALL LSTSQ $(Y, X, 100, N T O T, 2, A, S)$
WRITE (6,12) XN
12 FORMAT (1X, 5HXN $=, F 10.5 / / 1)$
WRITE $(6,20) \mathrm{S},(J, A(J), J=1,2)$
20 FORMAT $11 \mathrm{X}, 17 \mathrm{HSUM}$ OF SQUARĢ $=, E 13.6 / 1$ (1X,I3, E20.6)!
$A A=\operatorname{ExP}(A(1))$
$S S=E X P(S)$
$B=A(2)$
WRITE 16,44$)$ AA,B,SS
44 FORMAT $(4 H O A=, F 10.5 ; 5 X, 4 H 3=, F 10.5,17 H S U M$ OF SQUARES $=, E 13.6 / 1)$
DO $4 K=1,5$
4 WRITE 16.301
30 FORMAT (4H)
$X N=X N+0.5$
IF (XN.GT.10.0) GO TO 110
GO TO 100
110 RETURN
END
\$IBFTC PUNCH
SUBROUTINE (ARD (IH,IK,IL,E,E2,E2MI,SINT,CONK)
IF (E.LT.1.4) GO TO 10
W-ITE (6,11) IH,IK,IL,E,E2,E2MI,SINT,CONK
11. FORMAT $11 \mathrm{HS}, 314,5 \mathrm{Fg} .41$

10 RETURN
END

SUBROUTINE RANSET (N,SINMAX,S ,SM 1
DIMENSION $S(30), S M(30), 53(30), S 3 M E A N(30)$

3 FORMAT(1H . $110,4(5 \mathrm{X}, \mathrm{F} 10.51)$
4 FORMATIIHI:48H RANSET HAS BEEN CALLED FOR GENERATION OF RANGES)
WRITE $(6,4)$
$Z N=N$
S3MAX $=(S I N M A X) * * 3$
S3IN1 $=(S 3 M A X) / Z N$
S3IN2 $=$ S3IN1/2.0
SMIN $=0.0$
$001001=1, N$
$Z I=1$
S3(1) $=$ SMIN + (S3IN1*21)
$S(I)=(S 3(1)) * * 0.333333$
S3MEAN(I) $=53(1)-531 N 2$
SM(I) $=($ S3MEAN(I) ) $* * 0.333333$
100 CONTINUE
WRITE(6,1) N,SINMAX
WRITE(6,2)
DO $101 \mathrm{I}=1, \mathrm{~N}$
WRITE(6,3) 1,S3(I),S3MEAN(I),S(11),SM(1)
101 CONTINUE RETURN END

SISFTC SETUP I
SUBROUTINE FOUR1 (NFOUR,RCP)
C THIS SETS UP REC. DATA ON NFOUR -- REQUIRED FOR SHARPENED PATTERSON
DIMENSION RCP(6)
REWIND NFOUR
WRITE (NFOUR) (RCP(I),I=1,6)
WRITE(6,1) NFOUR
1 FORMAT 11 HO, $47 \mathrm{H}^{*} \mathrm{H}^{*}$ RECIPROCAL CELL DATA IS WRITTEN ON LOGICAL I5) RETURN
END

SIBFTC SETUP2
SUBROUTINE FOUR2 (IH,IK,IL,E,E2MI,SINT,ID,FOBS,NFOUR,EPSI)
DIMENSION FM(20)
THIS SET UP IS THE MAIN ONE FOR DATA FOR A SUBSIQUENT SHARPENED AND ORIGIN REMOVED SHARPENED PATTERSON
$C$
$F H=I H$
$F K=I K$
$F L=I L$
$Y C=0.0$
$A C=0.0$
$B O=0.0$
$B C=0.0$
SIGYO=0.0
U = 1.U
RHO $=$ U. U
DO $100 \quad I=1,10$
FM(1) $=0.0$
100 CONTINUE
QE2 $=E P S 1 * E * E$
$0=S Q R T($ OE2)
Q2M1 = QE2 - 1.0
W-ITE (NFOUR) FH,FK,FL,U,YC, UZMI,AC,BU,BC,SIGYO,O,ID,RHO,
1 SINT, (FM(I),I=1,IU)
IF(IH.NE.999) GO TO 101
WRITE 6,102 )
102 FORMAT (1HO, $49 \mathrm{HIH}=999$ IS REACHED AND TKANSHEK CUMPLETE TO FOUR2 END FILE NFOUR
REWIND NFOUK
101 RETUKN
END

## SIBFTC SYST

SUBROUTINE SYSTAB (IH,IK,IL,EPSI)
$C$
$C$
$C$
P21/C

|  | IF (IK.EQ.O) |
| :--- | :--- |
|  | IF IK.NE.U) |
|  | IF IK.NE.O) |
| 101 | EPSI. $=2.0$ |
|  | GOTO 100 |
| 102 | EPSI $=1.0$ |
| 100 | RETURN |
|  | END |

560*CARDS


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[^1]:    (1) H. Bergreen, Ber., 21, 337 (18888).
    (2) V. Meyer, ibid., 23,1571 (1890).
    (3) Idem. 1 bi d. 21,353 (1888).
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[^2]:    (7) R. Gompper and W. Töpfl, ibid., 95 , 2871 (1962). (8) W. Wachter, ibid., 25, 1727 (1892).
    (9) P. Yates and D. R. Moore, J. Am. Chem. Soc., 80, 5577 (1958).

[^3]:    (10) T. R. Lynch, Ph.D. Thesis, University of Toronto, 1966.

[^4]:    (11) M. Mammi, R. Bardic, G. Traverso and S. Bezzi, Nature, 192, 1282 (1961).

[^5]:    (14) J. A. Kapecki, J. E. Baldwin and I. C. Paul, Tetrahedron Letters, 5307 (1967).
    (15) A. Tulinsky, private communication.
    (16) W. H. Schmidt and A. Tulinsky, Tetrahedron Letters, 5311 (1967).

[^6]:    (17) A. L. Patterson, Phys. Rev. $\frac{46}{} 372$ (1934).
    (18) Idem., Z. Krist., 20, 517 (1935).
    \# Where ( $x, y, z$ ) are fractional coordinates.

[^7]:    (24) G. A. Sim, ibid., 10, 536 (1957).
    (25) C. A. Beevers and J.H. Robertson, ibid., 2, 164 (1950). (26) M. J. Buerger, ibid., 4, 531 (195i).

[^8]:    (28) J. Gillis, ibid. $, \frac{1}{2}, 174$ (1948)
    (29) J. S. Kasper, C. M. Lucht and D. Harker, ibid., 2, 436 (1950).
    (30) J. Karle and H. Hauptman, ibid., 2, 181 (1950).
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[^9]:    (32) W. H. Zachariasen, ibid., 5, 68 (1952).
    (33) H. Hauptman and J. Karle, "Solution of the Phase Problem, I, The Centrosymetric Crystal". A. C. A. Monograph No. 3, The Letter Shop, Wilmington, (1953).

[^10]:    (34) M. M. Woolfson, Acta Cryst., 7, 61 (1954).
    (35) W. Cochran and M. M. Woolfson, ibid., $\underset{\text { g , i }}{ }$ (1955).

[^11]:    \# A listing of the Lp program is given in Appendix $V$.

[^12]:    (56) B. B. Cetlin and S. C. Abrahams, Acta Cryst., 16, 943 (1963).
    (57) S. C. Abrahams, ibid., 17,1327 (1964).
    (58) S. C. Abrahams and J. L. Bernstein, ibid., 18 926, (1965).

[^13]:    \# Wilson's method for obtaining the absolute scaling factor and overall temperature factor is discussed in Chapter 4.

[^14]:    (63) The program used to compute the $\Sigma_{2}$ listing and combinational count was written by Mr. F. T. C. Leung, Chemistry Department, University of Toronto.
    (64) H. Hauptman and J. Karle, Acta Cryst., 12, 93 (1959).
    (65) K. Lonsdale and H. J. Grenville-Wells, ibid., 7, 490 (1954).

[^15]:    (67)K. S. Dragonette and I. L. Karle, Acta Cryst., 22, 978 (1965).

[^16]:    \# The electron density synthesis was computed for the quarter of the unit cell and with mesh intervals previously defined for the E-map synthesis.

[^17]:    \# The errors quoted with the cell dimensions represent the the maximum deviation of the calculated quantities from the mean value.
    $\neq \neq \mu$ was calculated in the manner described in Chapter 4.

[^18]:    \# The Picker diffractometer and its mode of operation for intensity data collection are described in Chapter 7, p. 47.

[^19]:    \# Details of the data reduction are given in Chapter 7, p. 51.

[^20]:    \# Normalised structure factor magnitudes were calculated using equation (8.4), p. 58 .

[^21]:    The highest peaks in the Patterson function, and the S-S vectors to which they correspond.

[^22]:    (72) J. M. Stewart, University of Washington, Seattle (1965).
    (73) R. Shiono and S. S.C. Chu, X-Ray Crystallography Laboratory, University of Pittsburg, Pittsburg (1964).
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[^23]:    * The abbreviation PDM for this molecule is retained throughout this chapter.

[^24]:    ₹ Reproduced from chapter 1.

[^25]:    \# The strict dihedral angle between the C(1)S(1)S(2) and

[^26]:    (88) A. C. Macdonald and J. Trotter, J. Chem. Soc. ${ }^{\text {B }}$, 1966, 929. (89) H. Luth and J. Trotter, Acta Cryst., 12, 614 (1965).

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[^28]:    \# Exocyclic is used loosely to indicate the $\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{C}(2)$ and $S(2)-C(1)-C(2)$ angles in PDM.

[^29]:    $\neq$ Hordvik assumes the van der Waals radius of sulphur is $1.85 \AA^{\circ}$.
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    (95) A. Hordvik and J. Sletten, ibid., 12, 753 (1965).
    (96) Idem., ibid., 20 , 1907 ( 1966 ).

