

V445 Puppis – Dustier than a Thousand Novae

2 D. P. K. BANERJEE ¹, A. EVANS ², C. E. WOODWARD ³, S. STARRFIELD ⁴, K. Y. L. SU ⁵, N. M. ASHOK ¹ AND
3 R. M. WAGNER ^{6,7}

4 ¹Physical Research Laboratory, Navrangpura, Ahmedabad, Gujarat 380009, India

5 ²Astrophysics Group, Keele University, Keele, Staffordshire, ST5 5BG, UK

6 ³Minnesota Institute for Astrophysics, University of Minnesota, 116 Church Street SE, Minneapolis, MN 55455, USA

7 ⁴School of Earth & Space Exploration, Arizona State University, Box 871404, Tempe, AZ 85287-1404, USA

8 ⁵Steward Observatory, University of Arizona, 933 North Cherry Avenue, Tucson, AZ 85721, USA

9 ⁶Department of Astronomy, Ohio State University, 140 W. 18th Avenue, Columbus, OH 43210, USA

10 ⁷Large Binocular Telescope Observatory, 933 North Cherry Avenue, Tucson, AZ 85721, USA

11 (Received 2023 May 10; Revised 2023 May 30; Accepted 2023 June 16; Published To appear in the ApJL)

12 ABSTRACT

13 V445 Puppis, the only known Galactic helium nova, is a unique testbed to verify supernova (SN) theories in
14 the single degenerate channel that involve a white dwarf (WD) accreting matter from a helium-rich donor. An
15 estimate of the mass of the helium shell on the WD is crucial to deciding whether or not it will undergo a SN
16 detonation. In this context, this study estimates the dust and ejecta masses in the 2000 November eruption of
17 V445 Pup. Subsequent to its outburst, the star became cocooned in a dust envelope. An analysis of the spectral
18 energy distribution (SED) of the dust using infrared data shows that V445 Pup produced at least $10^{-3} M_{\odot}$ of
19 dust which is unprecedented for a classical or recurrent nova. The SED can be explained by a combination of a
20 cold dust component at 105 ± 10 K, mass $(1.9 \pm 0.8) \times 10^{-3} M_{\odot}$, and a warm dust component at 255 ± 10 K,
21 mass $(2.2 \pm 1.2) \times 10^{-5} M_{\odot}$. For a conservative choice of the gas-to-dust mass ratio in the range 10–100, the
22 mass of the ejecta is 0.01–0.1 M_{\odot} . Such a high mass range raises the question: why did V445 Pup not detonate
23 as a Type Ia SN as is predicted in certain double-detonation sub-Chandrasekhar supernovae formalisms? We
24 re-examine the nature of V445 Pup and discuss its role as a potential SN progenitor.

25 *Keywords:* Classical novae (251), Chemical abundances (224), Dust shells (414), Explosive Nucleosynthesis
26 (503), Type Ia supernovae (1728)

27 1. INTRODUCTION

28 V445 Pup erupted in 2000 reaching a peak V brightness
29 of 8.46 mag on 2000 Nov 29, and then slowly declined
30 with a $t_2 \gtrsim 100$ days (t_2 being the elapsed time to decline
31 2 mags from peak brightness). Although the outburst was
32 first reported on 2000 December 30 by Kanatsu (Kato et al.
33 2000) archival All Sky Automated Survey (ASAS, Pojman-
34 ski 1997) records demonstrated the outburst had begun ear-
35 lier (Goranskij et al. 2010). V445 Pup appeared to be a slow
36 nova except that the spectra, both in the optical and near-
37 infrared (NIR), recorded in the immediate and post-eruption
38 stages, were unique in not showing the hydrogen lines con-
39 ventionally seen in a nova outburst. Instead, there were many

40 lines of carbon, helium, and other metals; the C and He lines
41 were specially prominent in the NIR (Iijima & Nakanishi
42 2008; Woudt & Steeghs 2005; Lynch et al. 2004; Ashok &
43 Banerjee 2003; Lynch et al. 2001; Wagner et al. 2001a,b).
44 Based on its spectrum Ashok & Banerjee (2003) proposed
45 V445 Pup to be a helium nova that had undergone a ther-
46 monuclear runaway in helium-rich matter accreted onto a
47 white dwarf’s (WD) surface from a helium-rich donor (e.g.,
48 Kato & Hachisu 2003; Iben & Tutukov 1994).

49 On 2001 Jan 2, about 34 days after peak brightness, JHK
50 photometry showed that hot dust had begun to form (Ashok
51 & Banerjee 2003) and 3–14 μm spectroscopy obtained on
52 2001 Jan 31, confirmed the presence of significant amounts
53 of carbon dust (Lynch et al. 2001). The dust shell rapidly
54 thickened from 2001 June and by 2001 October V445 Pup
55 had faded below $V = 20$ mag (Goranskij et al. 2010).

56 A remarkable $\simeq 2''$ hourglass nebula, expanding with time,
 57 (detected around the object with adaptive optics K_s band im-
 58 agery) showed high velocity outflows (Woudt et al. 2009).
 59 The knots at the tips of the hourglass had velocities as large
 60 as $\sim 8500 \text{ km s}^{-1}$ (Woudt et al. 2009). Flaring radio syn-
 61 chrotron radiation was persistently observed from the object
 62 from the beginning and up to 7 years after the outburst (Nya-
 63 mai et al. 2021; Rupen et al. 2001a). It is now clear that
 64 this non-thermal synchrotron emission was produced from
 65 shocks caused by the interaction of ejected matter (or a wind)
 66 from the WD with a pre-existing equatorial density enhance-
 67 ment collimating the ejecta to create the hour glass nebula
 68 with its pinched waist (Nyamai et al. 2021).

69 In this study, we re-analyze the SED of the dust from more
 70 recent archival data and conclude that V445 Pup has pro-
 71 duced an unprecedented amount of dust for a nova. The im-
 72 plications of the large dust mass on the role of V445 Pup as
 73 a SN Type 1a progenitor are discussed.

74 2. DISTANCE AND REDDENING

75 From the observed expansion parallax of the nebula,
 76 Woudt et al. (2009) derived a distance of $8.2 \pm 0.5 \text{ kpc}$. Iijima
 77 & Nakanishi (2008) using the radial velocities of the Na D1,
 78 D2 lines from high dispersion spectra, in conjunction with HI
 79 21cm radio data, estimated the reddening and distance to be
 80 $E_{B-V} = 0.51 \text{ mag}$ and $3.5 \lesssim d(\text{kpc}) \lesssim 6.5$. We point out that
 81 the equivalent width of 0.95 \AA for the Na DI line in the Iijima
 82 & Nakanishi (2008) data, calibrated using the Richmond et
 83 al. (1994) relations, yields $E_{B-V} = 0.45$, in reasonable agree-
 84 ment with the above. The $E_{B-V} = 0.51$ derived from the
 85 measurements of Iijima & Nakanishi (2008) is consistent
 86 with the estimate of Wagner et al. (2001b), $E_{B-V} \lesssim 0.8 \text{ mag}$.
 87 However, their distance estimate is based on low spatial res-
 88 olution extinction maps of Neckel & Klare (1980) whereas
 89 better data are now available. We use reddening data from
 90 Green et al. (2019) to estimate the distance (Fig. 1). The
 91 figure demonstrates, even up to the maximum distance of
 92 6.2 kpc (beyond which the data are stated to be unreliable)
 93 the E_{B-V} value has still not reached 0.51 (nor has it reached
 94 $E_{B-V} = 0.45$ obtained from the Na DI line). This suggests
 95 that the distance to V445 Pup is $\gtrsim 6.2 \text{ kpc}$. We adopt values
 96 of $d = 6.2 \text{ kpc}$ and $E_{B-V} = 0.51$ (hence $A_V = 1.6$ using
 97 $A_V = 3.1 \times E_{B-V}$).

98 3. ANALYSIS AND RESULTS

99 3.1. V445 Pup Dust mass

100 The SED of the dust was analyzed adopting the formal-
 101 ism of Sakon et al. (2016). The archival data used for the
 102 analysis are presented in Table 1. We assume a pure carbon
 103 composition for the dust given that the object showed a rich
 104 carbon spectrum in the optical and NIR (Iijima & Nakanishi
 105 2008). Further the $3\text{--}14 \mu\text{m}$ spectrum obtained and modeled

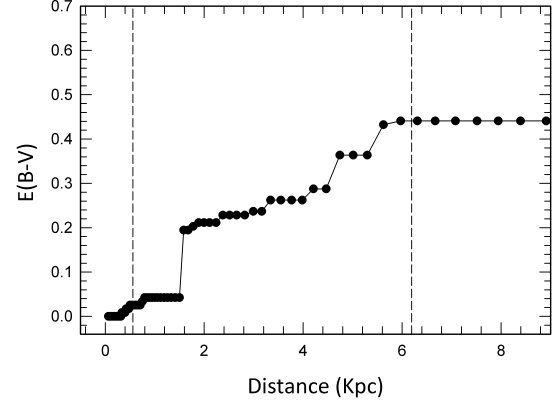


Figure 1. Reddening versus distance plot for V445 Pup from Green et al. (2019). The $E(g-r)$ values listed therein were converted to a mean E_{B-V} using the two relations given by the authors at <http://http://argonaut.skymaps.info/>. The vertical dotted lines give the range over which the method is stated to be reliable. See text for details.

106 by Lynch et al. (2001) showed a featureless continuum, with-
 107 out any silicate or unidentified infrared features (UIRs, see
 108 Evans et al. 2016), and thus strongly favors a carbon compo-
 109 sition.

110 Assuming an optically thin shell of spherical carbon dust
 111 grains with a uniform radius a , a total mass M_i with an equi-
 112 librium temperature $T_i(\text{K})$, located at a distance d from the
 113 observer, the observed flux density f ($\text{W m}^{-2} \mu\text{m}^{-1}$) is:

$$f_\nu^i = M_i \left(\frac{4\pi\rho_{(AC)}a^3}{3} \right)^{-1} \pi B_\nu(\lambda, T_i) Q_{abs}^{(AC)} \left(\frac{a}{d} \right)^2 \quad (1)$$

114 where $\rho_{(AC)}$ is the density of amorphous carbon (AC) dust
 115 (1.87 g cm^{-3}) and $Q_{abs}^{(AC)}$ is the absorption efficiency of
 116 amorphous carbon of radius $a(\mu\text{m})$ (BE sample; Zubko et al.
 117 1996). At IR wavelengths, $Q_{abs} \propto (8\pi a/\lambda)$ and so Eqn. (1)
 118 becomes independent of the dust grain size a (e.g., Kruegel
 119 2003; Bohren & Huffman 1983).

120 The foreground extinction of the emission by silicate
 121 grains in the interstellar medium (ISM) is taken into account
 122 by multiplying the equation above with the exponential term:

$$\exp\left\{-\tau_{9.7} \left(\frac{Q_{abs}^{Sil}(\lambda)}{Q_{abs}^{Sil}(9.7 \mu\text{m})} \right)\right\}. \quad (2)$$

123 We have normalized with the optical depth at $9.7 \mu\text{m}$ ($\tau_{9.7}$)
 124 determined from $(A_V/\tau_{9.7}) = 18.5 \pm 1.5$ (Roche & Aitken
 125 1984) with $A_V = 1.6$.

126 The SED can be understood as a combination of two com-
 127 ponents of amorphous carbon dust. The first component
 128 at $105 \pm 10 \text{ K}$ has a mass $(1.9 \pm 0.8) \times 10^{-3} M_\odot$. The
 129 second is a warm component, $255 \pm 10 \text{ K}$, with a mass
 130 $(2.2 \pm 1.2) \times 10^{-5} M_\odot$. The decomposition of the SED is

Table 1. V445 Puppis Archival IR/mm Photometry^{*}

Facility	λ (μm)	Brightness (Jy/mag)	Flux ($\text{W m}^2 \mu\text{m}^{-1}$)	Epoch
WISE	3.4	10.39 ± 0.23 (mag)	5.7080×10^{-15}	2010
WISE	4.6	7.297 ± 0.020 (mag)	2.9105×10^{-14}	2010
WISE [†]	12.0	0.554 ± 0.027 (mag)	3.9112×10^{-13}	2010
WISE [†]	22.0	-1.316 ± 0.010 (mag)	1.7106×10^{-13}	2010
AKARI-IRC	18.0	35.39 ± 1.82 (Jy)	3.2746×10^{-13}	2006–2007
AKARI-FIS2	65.0	12.55 ± 0.445 (Jy)	8.9051×10^{-15}	2006–2007
AKARI-FIS2	90.0	8.023 ± 0.236 (Jy)	2.9694×10^{-15}	2006–2007
AKARI-FIS2	140.0	3.430 ± 0.623 (Jy)	5.2460×10^{-16}	2006–2007
AKARI-FIS2	160.0	2.710 (Jy)	3.1736×10^{-16}	2006–2007
Spitzer-IRAC [‡]	3.6	0.11 ± 0.003 (Jy)	2.540×10^{-14}	2005
Spitzer-IRAC [‡]	4.5	0.32 ± 0.019 (Jy)	4.750×10^{-14}	2005
Spitzer-MIPS [‡]	70.0	7.53 ± 0.038 (Jy)	4.603×10^{-15}	2005
Herschel	70.0	7.582 ± 0.064 (Jy)	4.6388×10^{-15}	2012
Herschel	160.0	1.260 ± 0.044 (Jy)	1.3071×10^{-16}	2012
SEST	1200.0	0.0295 ± 0.0054 (Jy)	6.1400×10^{-20}	2003

NOTE— ^{*}Data retrieved from various mission archives hosted at the NASA/IPAC Infrared Science Archive including Spitzer (doi: 10.26131/IRSA433), WISE (doi: 10.26131/IRSA142), Herschel (doi: 10.26131/IRSA79), and AKARI (doi:10.26131/IRSA180, 10.26131/IRSA181). [†]The fractional pixel saturation in the 12 and 22 μm fluxes are 0.28 and 0.12 respectively but saturation effects are corrected by using profile fitted magnitudes (https://wise2.ipac.caltech.edu/docs/release/allsky/expsup/sec6_3d.html). [‡] Spitzer data reductions described in Su et al. (2020). The AKARI fluxes are averages over multiple detections made during the mission lifetime between 2006-2007.

131 shown in Fig. 2. To cross-check the mass estimates, the Q_{abs}
 132 of ACAR sample (Zubko et al. 1996) were also used, yielding
 133 very similar results as the BE sample. We have not consid-
 134 ered the 1.2 mm point in the fits as it is unclear whether the
 135 mm continuum flux is due to dust or from free-free emission
 136 from ionized gas (as discussed later).

137 Clearly there is a large mass of cool dust, based largely
 138 on the SED modeling of the long wavelength ($\lambda \gtrsim 10 \mu\text{m}$)
 139 photometry. Were the dust shell optically thick (e.g., $\tau \gtrsim 5$)
 140 at these wavelengths, rather than optically thin as assumed,
 141 would require an $A_V \simeq 50$ for extinction $\propto \lambda^{-1}$, compa-
 142 rable to that seen in the most opaque molecular clouds and
 143 likely is not reasonable. It is outside the scope of this study to
 144 compute the dust mass for more complex geometries (e.g., a
 145 bipolar morphology with ad-hoc assumptions on the amount
 146 of equatorial material enhancement).

147 Our SED fits involve modeling of data that is not contem-
 148 poraneous. However, the main result of the paper - the large
 149 dust mass - is largely based on fitting the longer wavelength
 150 ($\gtrsim 10 \mu\text{m}$) AKARI, Herschel, WISE and Spitzer data taken

151 between 2005-2012. These data are well-fit by our model
 152 suggesting that the dust temperature did not change signif-
 153 icantly between the different epochs, increasing our confi-
 154 dence that the data at There is some variability in the WISE
 155 and Spitzer 3.6 μm and 4.5 μm data, but the hotter dust com-
 156 ponent which is used to fit this data contributes only a small
 157 percent of the dust mass. So the total dust mass estimate
 158 should be reasonably reliable.

159 A similar but brief study, estimating the dust mass in
 160 V445 Pup was conducted by Shimamoto et al. (2017). How-
 161 ever, their modeling was limited to only the AKARI data with
 162 no data short-ward of 9 μm to constrain the Wien side of
 163 the SED. The source of the 9 μm AKARI flux used by Shi-
 164 mamoto et al. (2017) is also unclear. It is not listed in the
 165 master records in the AKARI database¹. The mm flux was
 166 not discussed and Shimamoto et al. (2017) invoked an un-
 167 realistic value of foreground extinction of $A_V = 12.5$ to fit

¹ https://darts.isas.jaxa.jp/astro/akari/data/AKARI-IRC_Catalogue_AllSky_PointSource_1.0.html

the SED. The present modeling is hence more comprehensive, improved, and realistic. Shimamoto et al. (2017) do conclude that large dust masses are extant, comprised of a combination of cold amorphous carbon (125 K) with a mass of $(0.45^{+0.66}_{-0.27}) \times 10^{-3} M_{\odot}$ and warm amorphous carbon (250 K) with a mass of $(1.8^{+1.0}_{-0.5}) \times 10^{-5} M_{\odot}$.

What becomes evident is that no nova, either recurrent nova (RN) or classical nova (CN), has produced as much dust ($\simeq 10^{-3} M_{\odot}$) as V445 Pup (adopting distances of 3.5 kpc or 8.2 kpc as extrema do not radically change the mass estimate as the dust mass scales as d^2). The typical mass of the dust produced in a dust producing nova is 10^{-6} to $10^{-9} M_{\odot}$ (Evans & Gehrz 2022). If a canonical value of gas-to-dust mass of 100 is assumed, the mass of the ejecta could be as large as $0.1 M_{\odot}$. This is unprecedented.

If the Swedish-ESO Submillimeter Telescope (SEST, Booth et al. 1989) 1.2 mm data point is included in the dust SED fitting (however, see §3.2), then additional modeling suggests that, apart from the two components used in the present analysis, an additional cooler component at $\simeq 30$ to 50 K and with a mass of $10^{-2} M_{\odot}$ is required to fit the composite SED. It thus appears certain that the mass ejected by V445 Pup was very large and that the accreted shell on the WD shell at the time of outburst was massive, at least $10^{-2} M_{\odot}$, for a most conservative choice of 10 for the gas-to-dust mass ratio. V445 Pup could have potentially erupted as a SN 1a, which curiously it did not. However, we first discuss the possible origin of the mm flux.

3.2. Millimeter/sub-mm studies of novae

Millimeter/sub-mm studies of novae appear to be few. The two possible origins for mm continuum fluxes are free-free emission (e.g., V1974 Cygni, Ivison et al. 1992) and/or the Rayleigh-Jeans tail of dust emission (e.g., V4743 Sgr, Schmidtobreick et al. 2005; Nielbock & Schmidtobreick 2003). Reasonably, the assertion that the detection of 1.2 mm emission from the nebula around V445 Pup is from free-free emission has a basis in two arguments. First, the 1.2 mm detection in V445 Pup was made in 2003 May about ~ 885 d after the outburst, at which stage the ionized nebula around V445 Pup was expected to be $\lesssim 1''$ in diameter (based on image sizes in Woudt et al. 2009). Second, the optically thin free-free flux $F_{ff}(\lambda)$ ($\text{W cm}^{-2} \mu\text{m}^{-1}$) at wavelength $\lambda(\mu\text{m})$ from an ionized gas with electron density $n_e(\text{cm}^{-3})$, assumed equal to n_i , the ion density, at temperature $T(K)$ and occupying a volume $V(\text{cm}^3)$ is (Banerjee et al. 2001):

$$F_{ff}(\lambda) = \frac{2.05 \times 10^{-30} \lambda^2 z^2 g T^{-0.5} n_e n_i V}{4\pi d^2} \quad (3)$$

where g is the Gaunt factor (between 0.3 and 0.5), z is the charge (2 for an ionized Helium gas) and d (in cm) is the distance.

As an illustrative example, the observed mm flux $F(1.2 \text{ mm}) = 6.14 \times 10^{-24} \text{ W cm}^{-2} \mu\text{m}^{-1}$ can be reproduced by considering a singly ionized He nebula of 0.6 arc-second extent at $T = 10,000 \text{ K}$ and $d = 6.2 \text{ kpc}$, with a $n_e = 2 \times 10^5 \text{ cm}^{-3}$ (since [O III] 5007Å was seen at around this time, Woudt & Steeghs 2005). We approximate the density to be less than the critical density of that line which is $\simeq 6 \times 10^5 \text{ cm}^{-3}$. These assumptions are reasonable and serve to show that the mm flux is almost certainly due to free-free emission.

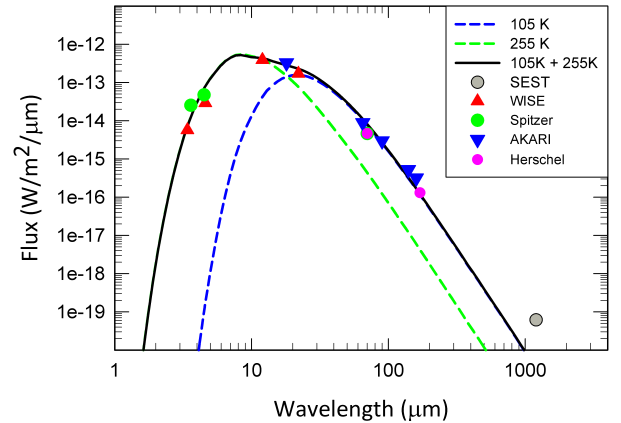


Figure 2. Model fit to the SED of V445 Pup using amorphous carbon grains (BE sample; Zubko et al. 1996). The black bold line is the composite of a 105 K plus a 255 K component. The 1.2mm point was not considered in the fitting.

3.3. Pathways for SN 1a explosions and novae

One of the pathways for SN1a explosions, within the single-degenerate channel, is the double-detonation sub-Chandrasekhar mechanism wherein the WD accretes from a helium donor (Maoz et al. 2014). A detonation in the accreted helium shell can cause a secondary detonation of the carbon/oxygen white dwarf (C/O WD) core (Starrfield et al. 2021; Fink et al. 2010; Shen & Bildsten 2009; Livne & Arnett 1995; Woosley & Weaver 1994). This happens at a total mass below the Chandrasekhar limit and depends on two factors – first the formation of a detonation in the helium shell and second whether a successful detonation of the helium shell can detonate the core. V445 Pup, by virtue of having had a shell detonation, is thus the only object which can be used as a testbed for the above theory.

A secondary detonation can be triggered in two different ways: either directly when the helium detonation shock hits the core/shell interface (“edge-lit”), or with some delay, after the shock has converged near the center (Fink et al. 2010). Fink et al. (2010, 2007) examined the delayed mechanism

since it could lead to a core detonation even for shocks too weak for the edge-lit case. They find that secondary core detonations are triggered for all their simulated models, ranging in core mass from $0.810 M_{\odot}$ up to $1.385 M_{\odot}$ with corresponding shell masses from $0.126 M_{\odot}$ down to $0.0035 M_{\odot}$. For convenience, we reproduce from Table 1 of Fink et al. (2010), the core mass of the WD and shell mass in the format $(M_{\text{core}}, M_{\text{shell}})$ for all their models that undergo double detonation: $(0.81, 0.126)$, $(0.92, 0.084)$, $(1.025, 0.055)$, $(1.125, 0.039)$, $(1.28, 0.013)$, $(1.385, 3.5 \times 10^{-3})$. The end result of their modeling is that as soon as a detonation occurs in a helium shell covering a carbon/oxygen WD a subsequent core detonation is virtually inevitable (Fink et al. 2010).

The WD mass is unknown but we speculate it is low based on the low amplitude outburst, the extremely long time to decline, the formation of dust, the amount of mass ejected, the low excitation spectrum at outburst (Ashok & Banerjee 2003), and the lack of coronal line emission even 3 yrs after outburst (Lynch et al. 2004). Piersanti et al. (2014) suggest that the initial mass of the WD in V445 Pup could have been close to $0.8 M_{\odot}$. A low mass could explain why a SN Ia explosion was averted by the double-detonation channel.

The 2000 eruption of V445 Pup was not a SN explosion. With $m_V^{\text{max}} = 8.46$ (Goranskij et al. 2010), $A_V = 1.6$, d in the range 3.5 to 8.2 kpc, the absolute magnitude M_V is in the range -6 to -7.7 , typical of a slow nova. Thus, its peak brightness falls greatly short of even the weakest SN “impostor” explosions which have $M_V \sim -13$ to -14 (Kasliwal 2012; Smith et al. 2009).

3.4. V445 Pup – Its nature and connection with SNe

Useful insight about the system can be obtained from Fig. 3 which presents a plot of the extinction corrected pre-outburst SED of the progenitor from 2MASS (Skrutskie et al. 2006) and DENIS (Epchtein 1994) magnitudes, both of which were recorded almost simultaneously on 1999 February 02 and 11 respectively (2MASS: $J = 12.271 \pm 0.026$, $H = 11.94 \pm 0.024$, $K_s = 11.52 \pm 0.025$; DENIS: $I = 12.85 \pm 0.02$, $J = 12.24 \pm 0.06$, $K_s = 11.36 \pm 0.09$). We also use $B = 14.3 \pm 0.3$ from Goranskij et al. (2010). The SED is well fit by a 10^4 K star, (cf., Goranskij et al. 2010) but with a discernible IR excess. This excess may be explained by free-free emission from ionized circumbinary gas that existed before the eruption (Fig. 3).

Several factors point towards the existence of material around V445 Pup prior to the eruption. First, the radio synchrotron emission requires shocks formed by nova ejecta plowing into pre-existing material. Second, the bipolar morphology of the nebula necessitates an equatorial constriction for shaping (the radio images confirm the constriction; Nyamai et al. 2021). Third, the putative presence of free-free emission demands ionized gas in the vicinity of the central

star. Fourth, the presence of a large number of narrow absorption features (e.g., Ba II 493.408 nm, Sc II 552.679 nm and several lines of Ti II, Cr II and Fe II) seen in the rich, high dispersion ($R = 8000$) spectra (Iijima & Nakanishi 2008) likely are Transient Heavy Element Absorption (THEA) systems (Williams et al. 2008). The latter authors propose that gas causing the THEA absorption systems in novae must be circumbinary, exists before the outburst, with a likely origin arising from mass ejection from the secondary star.

Hence, preexisting (equatorial) material, if dusty, would contribute to the hotter (less massive) dust component. Our analysis does not rule out this possibility. Colder dust, at the temperature of the cooler component, may also have pre-existed but in quantities below that required to allow a pre-outburst detection by IRAS (Neugebauer et al. 1984) which had an average 10σ sensitivity of 0.85 Jy at $60 \mu\text{m}$ and 3 Jy at $100 \mu\text{m}$.

These factors, when viewed collectively, would favor a pre-eruption configuration for V445 Pup that consists of a binary system, with remnants of the common envelope (CE) phase forming a torus in the equatorial plane. Radiation from a hot source, likely the WD or a hot accretion disc (or a combination of both) ionized parts of the CE remnant material leading to the observed free-free emission. The secondary was likely a star with $T_{\text{eff}} = 10^4$ K, that was a periodic variable star with a probable orbital period of 0.650654 ± 0.000011 d (Goranskij et al. 2010).

These parameters (spectral type, orbital period of secondary, mass of the ejecta) are essential inputs while modeling the thermal response of a degenerate C/O WD accreting helium from a helium-star donor (e.g., Brooks et al. 2016; Piersanti et al. 2014). Both these cited studies show that various outcomes are possible from the accretion process de-

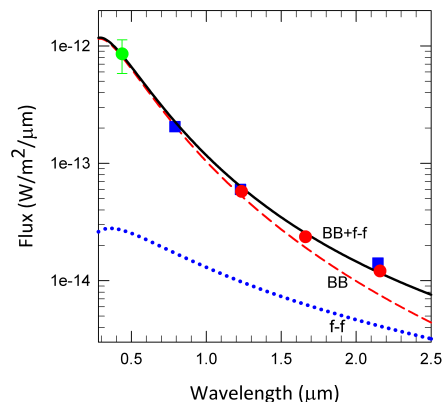


Figure 3. The pre-eruption SED of the progenitor fit by the sum of a blackbody (BB) at 10000 K (dashed line) and free-free (f-f, dotted line) emission at 10000 K. The sum of the BB and f-f emission is shown by the bold line.

330 pending on the accretion rate and WD mass, whether there is
 331 steady He burning on the surface, mild shell flashes, strong
 332 shell flashes, or quiet accumulation of matter up to the final
 333 SN 1a explosion when the mass crosses the Chandrasekhar
 334 limit. If the accretion rate is low ($\dot{M} \lesssim 10^{-6} M_{\odot} \text{ yr}^{-1}$) he-
 335 lium flashes result (Brooks et al. 2016; Piersanti et al. 2014)
 336 yielding a helium nova (Jacobson-Galán et al. 2019).

337 Furthermore, helium novae could be related to Type 1ax
 338 SNe because the helium emission in some Type 1ax SNe ap-
 339 pears to arise from the circumstellar environment rather than
 340 from the supernova ejecta itself (Jacobson-Galán et al. 2019).
 341 The argument is that the ejecta of a SN 1ax detonation, fol-
 342 lowing an earlier helium nova eruption on the same star, en-
 343 trains the helium injected by the latter’s eruption into the cir-
 344 cumstellar environment. Type 1ax supernovae share similar
 345 characteristics as SN 1a but exhibit lower peak luminosities
 346 and ejecta velocities. Jacobson-Galán et al. (2019) point out
 347 that the helium emission in two SN 1ax, SNe 2004cs and
 348 2007J, is consistent with coming from the ejecta of a rela-
 349 tively recent helium nova, and note in particular, that the ve-
 350 locity of the material in these two SNe is similar to that of the
 351 galactic helium nova V445 Pup. Recently, Kool et al. (2023)
 352 discuss the strong likelihood of a V445 Pup type object being
 353 the progenitor of the first radio-detected Type 1a SN 2020eyj
 354 which has a helium-rich circumstellar medium. Explaining
 355 the radio light curve and the bolometric light-curve tail of
 356 SN2020eyj requires a circumstellar medium mass between
 357 $0.3\text{--}1.0 M_{\odot}$. A mass of $0.67 M_{\odot}$, in good agreement with
 358 that posited for SN 2020eyj, can be provided by the he-
 359 lium donor in V445 Pup if we adopt a distance of 8.2 kpc
 360 (Woudt et al. 2009) instead of 6.2 kpc, a plausible gas-to-
 361 mass ratio of 200, in tandem with the cold component mass
 362 of $(1.9 \pm 0.8) \times 10^{-3} M_{\odot}$ that we derive. V445 Pup is thus
 363 a unique test platform for testing single-degenerate channel
 364 theories that involve a helium-rich donor.

365 Although the ejected mass in V445 Pup is unusually high
 366 for a nova, a helium nova outburst is still the most favorable
 367 interpretation for the 2000 eruption. However, V445 Pup
 368 also shares certain similarities with CK Vul, the latter pro-
 369 posed to belong to the class of objects known as intermediate-
 370 luminosity red transients (ILRTs) or interchangeably Lumi-
 371 nous red novae (Banerjee et al. 2020). CK Vul also has a
 372 hourglass morphology, a similar dust mass of $4.3 \times 10^{-3} M_{\odot}$
 373 in the inner nebula, and peak expansion velocities of \simeq
 374 2000 km s^{-1} (Banerjee et al. 2020; Eyres et al. 2018). How-
 375 ever, CK Vul was not hydrogen deficient and was much more
 376 luminous at the peak of its outburst ($M_V \sim -12.4$) compared
 377 to V445 Pup ($M_V \sim -7$).

378 4. SUMMARY

379 The principal result is that V445 Pup, at the time of its out-
 380 burst, had a shell mass as massive as $0.01 M_{\odot}$ or more and

381 thus should have potentially undergone a SN 1a detonation
 382 by the double-detonation sub-Chandrasekhar pathway. That
 383 it did not suggests that the WD is not massive.

ACKNOWLEDGMENTS

The authors thank the anonymous referee’s critique of the manuscript that improved the scientific narrative. This work is based in part on observations made with the Spitzer Space Telescope, obtained from the NASA/ IPAC Infrared Science Archive (doi: 10.26131/IRSA433), both of which are operated by the Jet Propulsion Laboratory, California Institute of Technology under a contract with the National Aeronautics and Space Administration. This publication makes use of data products from the Wide-field Infrared Survey Explorer (doi: 10.26131/IRSA142), which is a joint project of the University of California, Los Angeles, and the Jet Propulsion Laboratory/California Institute of Technology, funded by the National Aeronautics and Space Administration. Herschel is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA (doi: 10.26131/IRSA79). This research is based on observations with AKARI (doi:10.26131/IRSA180, 10.26131/IRSA181), a JAXA project with the participation of ESA. This publication makes use of data products from the Two Micron All Sky Survey (doi: 10.26131/IRSA2), which is a joint project of the University of Massachusetts and the Infrared Processing and Analysis Center/California Institute of Technology, funded by the National Aeronautics and Space Administration and the National Science Foundation and the Deep Near-Infrared Survey of the Southern Sky (DENIS) Catalog (doi: 10.26131/IRSA478).

384 *Facilities:* AKARI, 2MASS, DENIS, WISE, Spitzer,
 385 SEST, Herschel

386 *Software:* IRAF, Astropy (Astropy Collaboration et al.
 387 2018)

REFERENCES

- 388 Ashok, N. M. & Banerjee, D. P. K. 2003, *A&A*, 409, 1007.
389 doi:10.1051/0004-6361:20031160
- 390 Astropy Collaboration, Price-Whelan, A. M., Sipőcz, B. M., et al.
391 2018, *AJ*, 156, 123. doi:10.3847/1538-3881/aabc4f
- 392 Banerjee, D. P. K., Janardhan, P., & Ashok, N. M. 2001, *A&A*,
393 380, L13. doi:10.1051/0004-6361:20011508
- 394 Banerjee, D. P. K., Geballe, T. R., Evans, A., et al. 2020, *ApJL*,
395 904, L23. doi:10.3847/2041-8213/abc885
- 396 Bohren, C. F. & Huffman, D. R. 1983, *Light Scattering of Small*
397 *Particles*, New York: Wiley, 1983
- 398 Booth, R. S., Delgado, G., Hagstrom, M., et al. 1989, *A&A*, 216,
399 315
- 400 Brooks, J., Bildsten, L., Schwab, J., et al. 2016, *ApJ*, 821, 28.
401 doi:10.3847/0004-637X/821/1/28
- 402 Epchtein, N. 1994, *Experimental Astronomy*, 3, 73.
403 doi:10.1007/BF00430120
- 404 Evans, A. & Gehrz, R. D. 2022, arXiv:2211.12410.
405 doi:10.48550/arXiv.2211.12410
- 406 Evans, A., Gehrz, R. D., Woodward, C. E., et al. 2016, *MNRAS*,
407 457, 2871. doi:10.1093/mnras/stw352
- 408 Eyres, S. P. S., Evans, A., Zijlstra, A., et al. 2018, *MNRAS*, 481,
409 4931. doi:10.1093/mnras/sty2554
- 410 Fink, M., Hillebrandt, W., & Röpke, F. K. 2007, *A&A*, 476, 1133.
411 doi:10.1051/0004-6361:20078438
- 412 Fink, M., Röpke, F. K., Hillebrandt, W., et al. 2010, *A&A*, 514,
413 A53. doi:10.1051/0004-6361/200913892
- 414 Goranskij, V., Shugarov, S., Zharova, A., et al. 2010, *Peremennye*
415 *Zvezdy*, 30, 4. doi:10.48550/arXiv.1011.6063
- 416 Green, G. M., Schlafly, E., Zucker, C., et al. 2019, *ApJ*, 887, 93.
417 doi:10.3847/1538-4357/ab5362
- 418 Iben, I. & Tutukov, A. V. 1994, *ApJ*, 431, 264. doi:10.1086/174484
- 419 Iijima, T. & Nakanishi, H. 2008, *A&A*, 482, 865.
420 doi:10.1051/0004-6361:20077502
- 421 Ivison, R. J., Bang, M. K., & Bode, M. F. 1992, *IAUC*, 5516
- 422 Jacobson-Galán, W. V., Foley, R. J., Schwab, J., et al. 2019,
423 *MNRAS*, 487, 2538. doi:10.1093/mnras/stz1305
- 424 Kasliwal, M. M. 2012, *PASA*, 29, 482. doi:10.1071/AS11061
- 425 Kato, M. & Hachisu, I. 2003, *ApJL*, 598, L107.
426 doi:10.1086/380597
- 427 Kato, T., Kanatsu, K., Takamizawa, K., et al. 2000, *IAUC*, 7552
- 428 Kool, E. C., Johansson, J., Sollerman, J., et al. 2023, *Nature*, 617,
429 477. doi:10.1038/s41586-023-05916-w
- 430 Kruegel, E. 2003, *The physics of interstellar dust*, by Endrik
431 Kruegel. *IoP Series in astronomy and astrophysics*, ISBN
432 0750308613. Bristol, UK: The Institute of Physics, 2003.
- 433 Livne, E. & Arnett, D. 1995, *ApJ*, 452, 62. doi:10.1086/176279
- 434 Lynch, D. K., Rudy, R. J., Mazuk, S., et al. 2004, *AJ*, 128, 2962.
435 doi:10.1086/425889
- 436 Lynch, D. K., Russell, R. W., & Sitko, M. L. 2001, *AJ*, 122, 3313.
437 doi:10.1086/324465
- 438 Maoz, D., Mannucci, F., & Nelemans, G. 2014, *ARA&A*, 52, 107.
439 doi:10.1146/annurev-astro-082812-141031
- 440 Neckel, T. & Klare, G. 1980, *A&AS*, 42, 251
- 441 Neugebauer, G., Habing, H. J., van Duinen, R., et al. 1984, *ApJL*,
442 278, L1. doi:10.1086/184209
- 443 Nielbock, M. & Schmidtobreick, L. 2003, *A&A*, 400, L5.
444 doi:10.1051/0004-6361:20030107
- 445 Nyamai, M. M., Chomiuk, L., Ribeiro, V. A. R. M., et al. 2021,
446 *MNRAS*, 501, 1394. doi:10.1093/mnras/staa3712
- 447 Piersanti, L., Tornambé, A., & Yungelson, L. R. 2014, *MNRAS*,
448 445, 3239. doi:10.1093/mnras/stu1885
- 449 Pojmanski, G. 1997, *AcA*, 47, 467.
450 doi:10.48550/arXiv.astro-ph/9712146
- 451 Richmond, M. W., Treffers, R. R., Filippenko, A. V., et al. 1994,
452 *AJ*, 107, 1022. doi:10.1086/116915
- 453 Roche, P. F. & Aitken, D. K. 1984, *MNRAS*, 208, 481.
454 doi:10.1093/mnras/208.3.481
- 455 Rupen, M. P., Mioduszewski, A. J., & Dhawan, V. 2001a, *IAUC*,
456 7728
- 457 Sakon, I., Sako, S., Onaka, T., et al. 2016, *ApJ*, 817, 145.
458 doi:10.3847/0004-637X/817/2/145
- 459 Schmidtobreick, L., Nielbock, M., & Manthey, E. 2005, *The*
460 *Astrophysics of Cataclysmic Variables and Related Objects*,
461 330, 483
- 462 Shen, K. J. & Bildsten, L. 2009, *ApJ*, 699, 1365.
463 doi:10.1088/0004-637X/699/2/1365
- 464 Shimamoto, S., Sakon, I., Onaka, T., et al. 2017, *Publication of*
465 *Korean Astronomical Society*, 32, 109.
466 doi:10.5303/PKAS.2017.32.1.109
- 467 Skrutskie, M. F., Cutri, R. M., Stiening, R., et al. 2006, *AJ*, 131,
468 1163. doi:10.1086/498708
- 469 Smith, N., Ganeshalingam, M., Chornock, R., et al. 2009, *ApJL*,
470 697, L49. doi:10.1088/0004-637X/697/1/L49
- 471 Starrfield, S., Bose, M., Iliadis, C., et al. 2021, *The Golden Age of*
472 *Cataclysmic Variables and Related Objects V*, 2-7, 30.
473 doi:10.22323/1.368.0030
- 474 Su, K. Y. L., Rieke, G. H., Melis, C., et al. 2020, *ApJ*, 898, 21.
475 doi:10.3847/1538-4357/ab9c9b
- 476 Wagner, R. M., Schwarz, G., Starrfield, S. G., et al. 2001a, *IAUC*,
477 7717
- 478 Wagner, R. M., Schwarz, G., & Starrfield, S. G. 2001b, *IAUC*,
479 7571
- 480 Williams, R., Mason, E., Della Valle, M., et al. 2008, *ApJ*, 685,
481 451. doi:10.1086/590056
- 482 Woosley, S. E. & Weaver, T. A. 1994, *ApJ*, 423, 371.
483 doi:10.1086/173813

484 Woudt, P. A., Steeghs, D., Karovska, M., et al. 2009, *ApJ*, 706,
485 738. doi:10.1088/0004-637X/706/1/738

486 Woudt, P. A. & Steeghs, D. 2005, *Interacting Binaries: Accretion,*
487 *Evolution, and Outcomes*, 797, 647. doi:10.1063/1.2130306
488 Zubko, V. G., Mennella, V., Colangeli, L., et al. 1996, *MNRAS*,
489 282, 1321. doi:10.1093/mnras/282.4.1321